

Supernova Nucleosynthesis

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November 25, 2009

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Overview of Supernova Nucleosynthesis

- changes the content of the ISM
- produces heavy elements
- nuclear decay of ^{56}Co , produced during supernova nucleosynthesis, powers the light from supernova remnants
- observations of gamma rays emitted by supernova nucleosynthesis can provide important constraints on theoretical models
 - currently the deepest probe of supernova ejecta (gravitational waves would be deeper)

Radioactive Isotopes Synthesized in Supernovae

Isotope	Decay chain	Lifetime	Line energy (keV)
^{56}Ni	$^{56}\text{Ni} \rightarrow ^{56}\text{Co}$	8.8d	158, 812, 750, 480
^{56}Co	$^{56}\text{Co} \rightarrow ^{56}\text{Fe}$	111d	847, 1238
^{57}Ni	$^{57}\text{Ni} \rightarrow ^{57}\text{Co} \rightarrow ^{57}\text{Fe}$	(52h) 390d	122, 136
^{44}Ti	$^{44}\text{Ti} \rightarrow ^{44}\text{Sc} \rightarrow ^{44}\text{Ca}$	89yr (5.4h)	78, 68, 1157
^{26}Al	$^{26}\text{Al} \rightarrow ^{26}\text{Mg}$	1.0×10^6 yr	1809
^{60}Fe	$^{60}\text{Fe} \rightarrow ^{60}\text{Co} \rightarrow ^{60}\text{Ni}$	2.0×10^6 yr (7.6yr)	1173, 1332

Core-Collapse Supernovae

- nucleosynthesis in core-collapse supernovae occurs in the form of explosive burning when the shock wave passes through layers of unburned material, igniting and accelerating it
- although reaction systems are complicated, yields can be predicted with use of some simplifying approximations:
 - **Nuclear Statistical Equilibrium:** for sufficiently high temperature and density, strong and electromagnetic thermonuclear reaction rates are rapid enough to achieve equilibrium within the timescale of the supernova
 - **Quasi-Equilibrium:** for temperatures in the range of 3-6 GK, although there is no global equilibrium, many nuclei are in local equilibrium with their neighbors

Carbon Burning

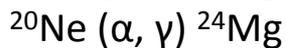
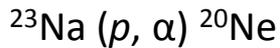
- calculations of explosive carbon burning for constant temperature done in Arnett and Truran 1969
 - took constant density of $5 \cdot 10^5 \text{ g cm}^{-3}$
 - nucleosynthesis during carbon burning does not sensitively depend on the density
 - initial composition of equal parts ^{12}C and ^{16}O

- reactions in carbon burning that produce light particles:

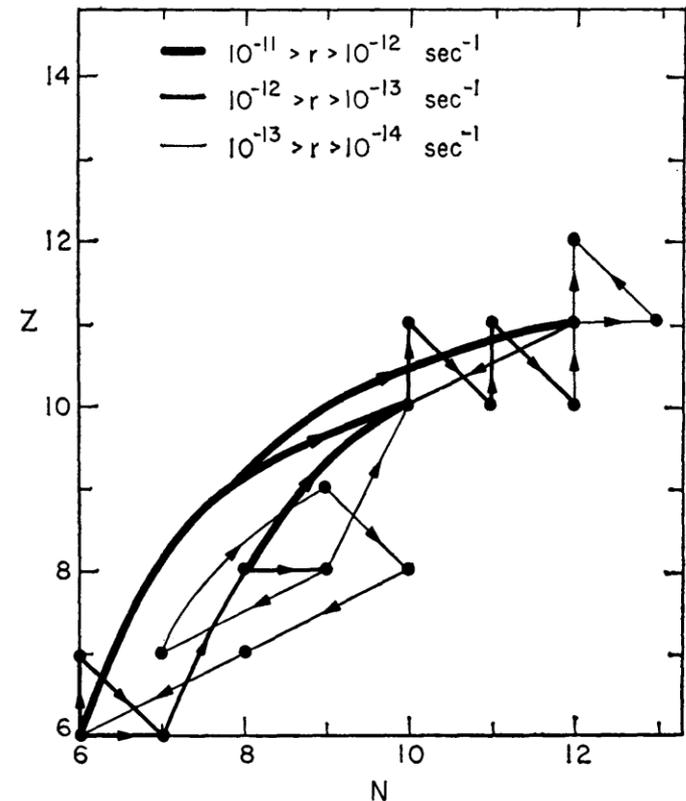


- these light particles react much faster than the original $^{12}\text{C} + ^{12}\text{C}$ reaction

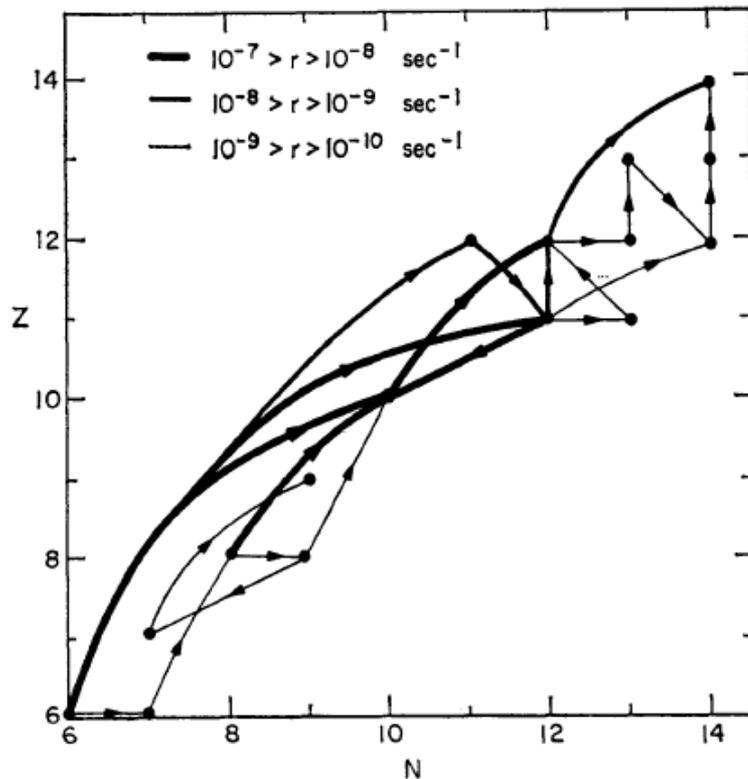
- most significant reactions:



Nuclear reaction flow diagram for carbon burning at $T = 1.2 \cdot 10^9 \text{ K}$ with little carbon depletion



Carbon Burning



Nuclear reaction flow diagram for carbon burning at $T = 1.2 \cdot 10^9$ K at half carbon depletion

○ neutron excess = (total neutrons – total protons)/total nucleons

○ the number of free neutrons, protons, and alpha particles is small

○ an increase in proton absorbers, particularly ^{23}Na , results in a decrease in protons and a constant number of alpha particles through the $^{23}\text{Na} (p, \alpha) ^{20}\text{Ne}$ reaction

○ neutron number is kept high by $^{12}\text{C} (^{12}\text{C}, n) ^{23}\text{Mg}$

○ when ^{12}C depletion becomes significant, the number of neutrons and alpha particles also begins to decrease

Carbon Burning

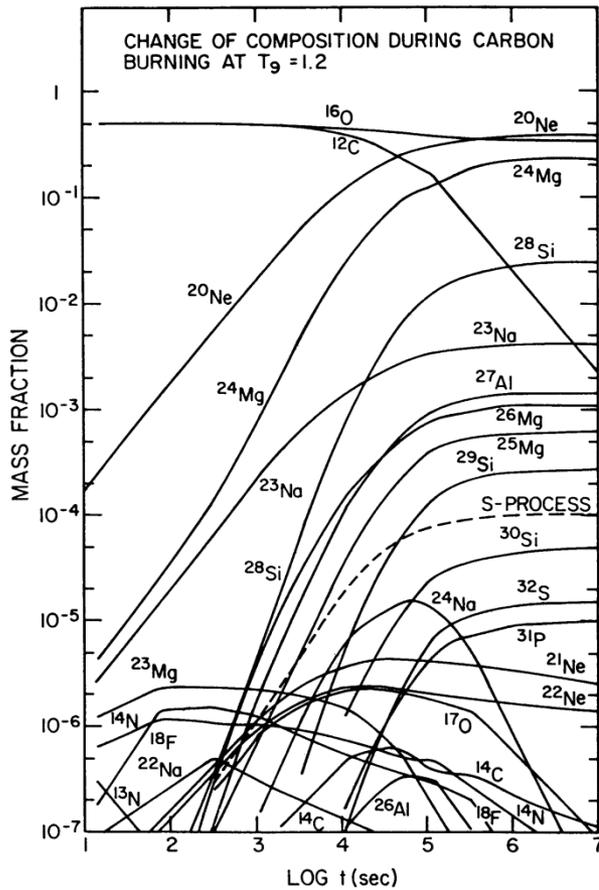


FIG. 4

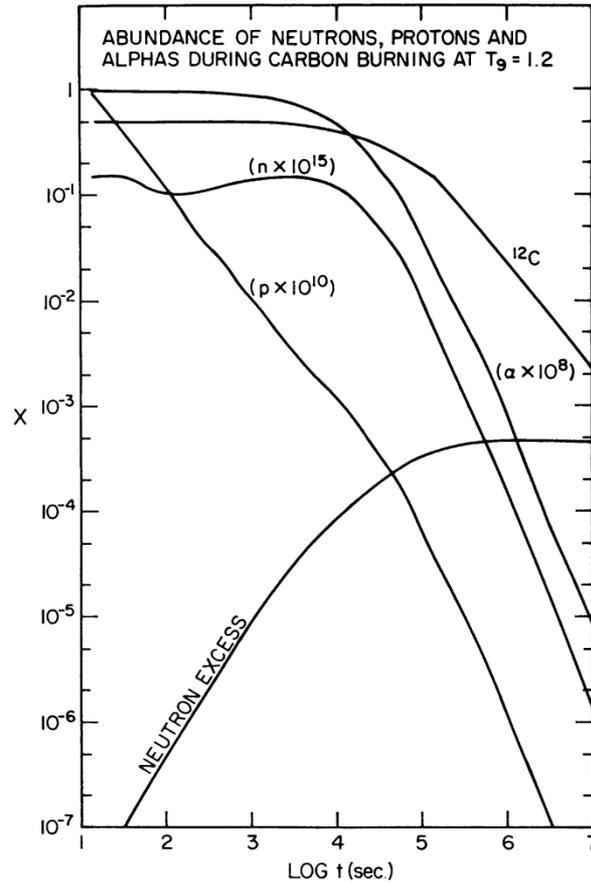


FIG. 5

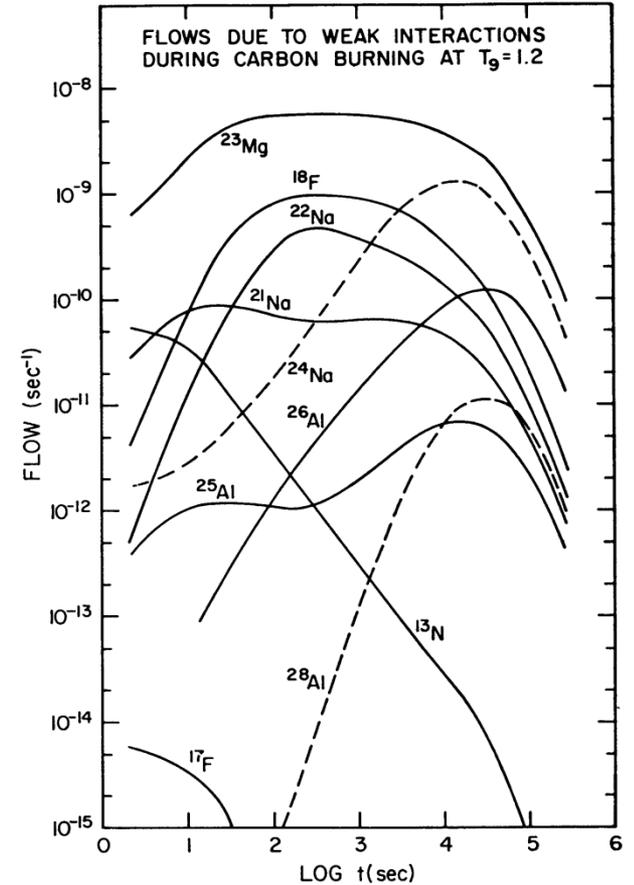


FIG. 6

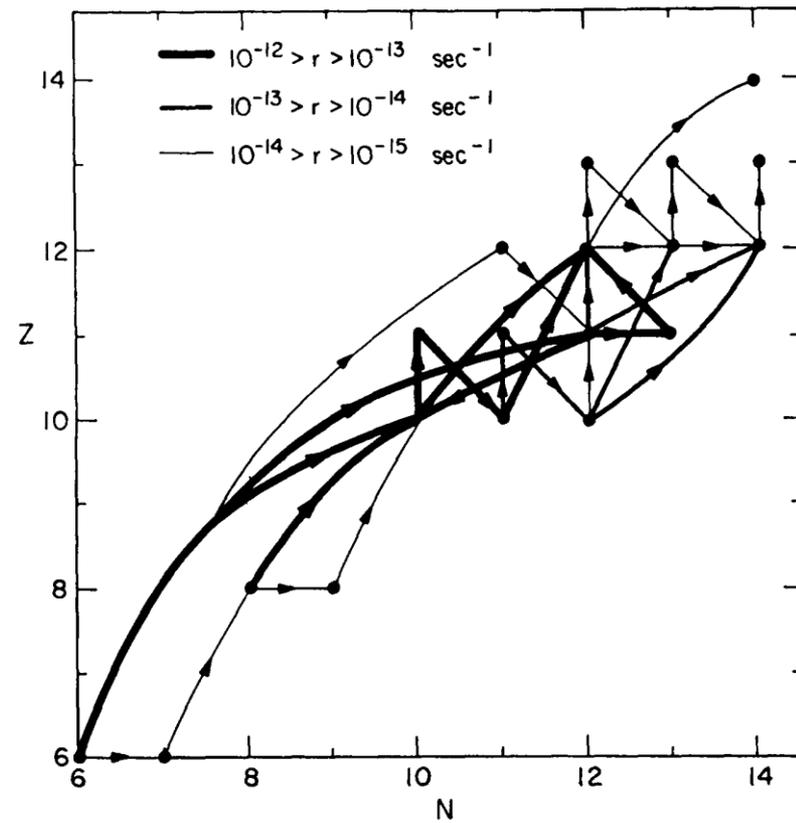
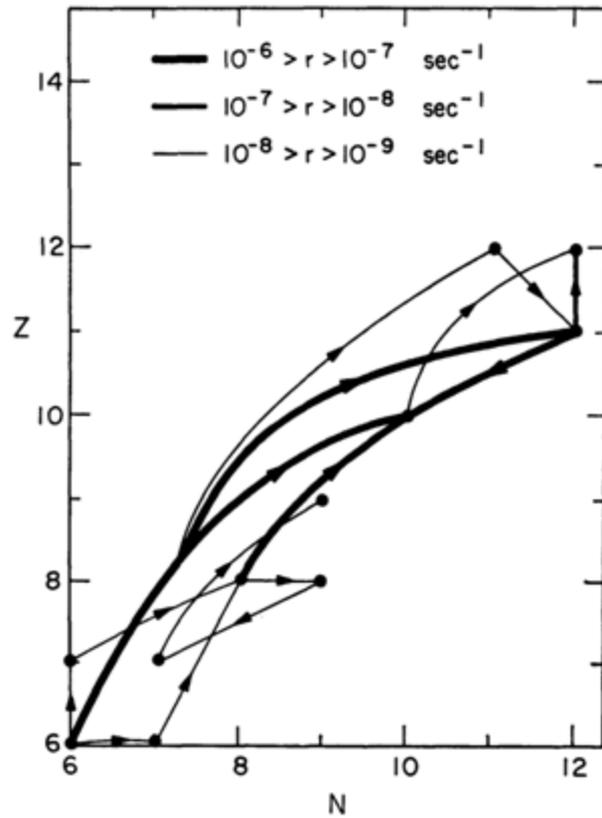
FIG. 4.—Variation of isotopic abundance by mass with time during carbon burning at $T_9 = 1.2$.

FIG. 5.—Variation of abundance by mass of neutrons, protons, and α -particles with time during carbon burning at $T_9 = 1.2$.

FIG. 6.—Variation of flows due to the weak interaction with time during carbon burning at $T_9 = 1.2$. *Solid lines, β^+ decays; dashed lines, β^- decays.*

Carbon Burning

○ at $T = 0.8 \cdot 10^9$ K, the gross behavior is similar to that at $T = 1.2 \cdot 10^9$ K



Carbon Burning

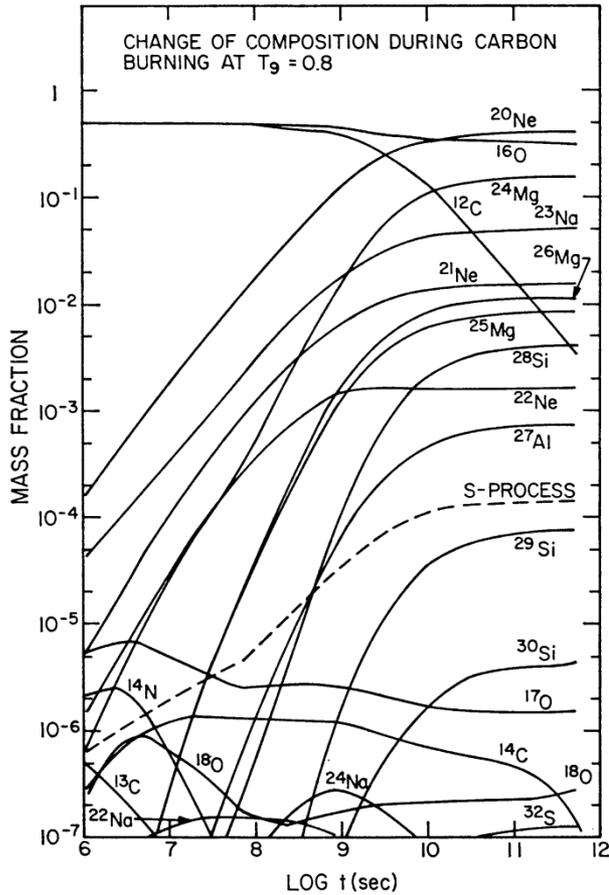


FIG. 9

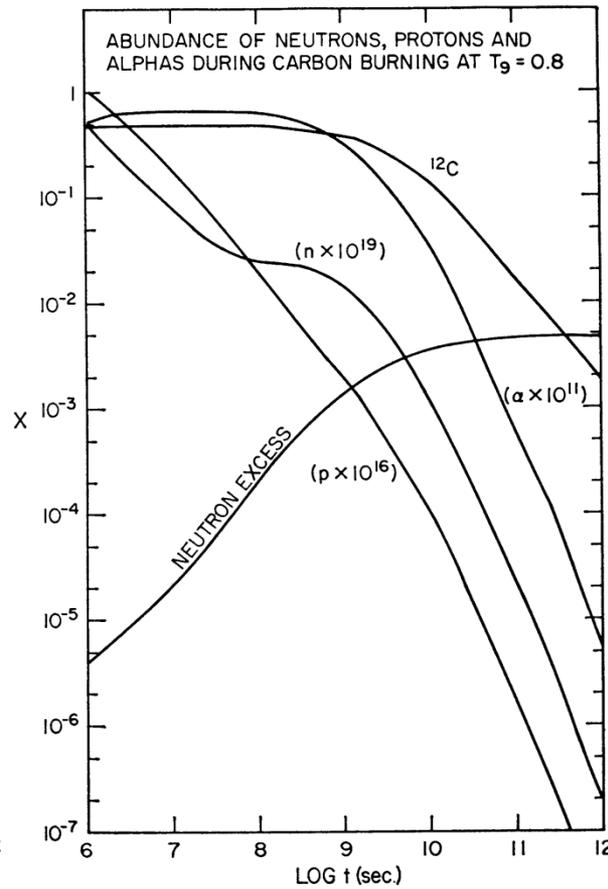


FIG. 10

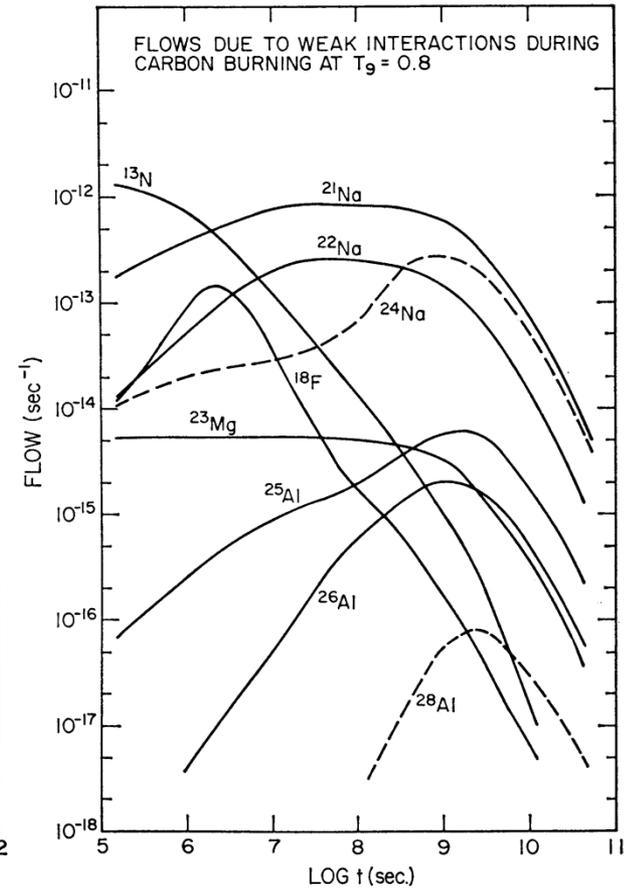


FIG. 11

FIG. 9.—Variation of isotopic abundance by mass with time during carbon burning at $T_9 = 0.8$.

FIG. 10.—Variation of abundance by mass of neutrons, protons, and α -particles with time during carbon burning at $T_9 = 0.8$.

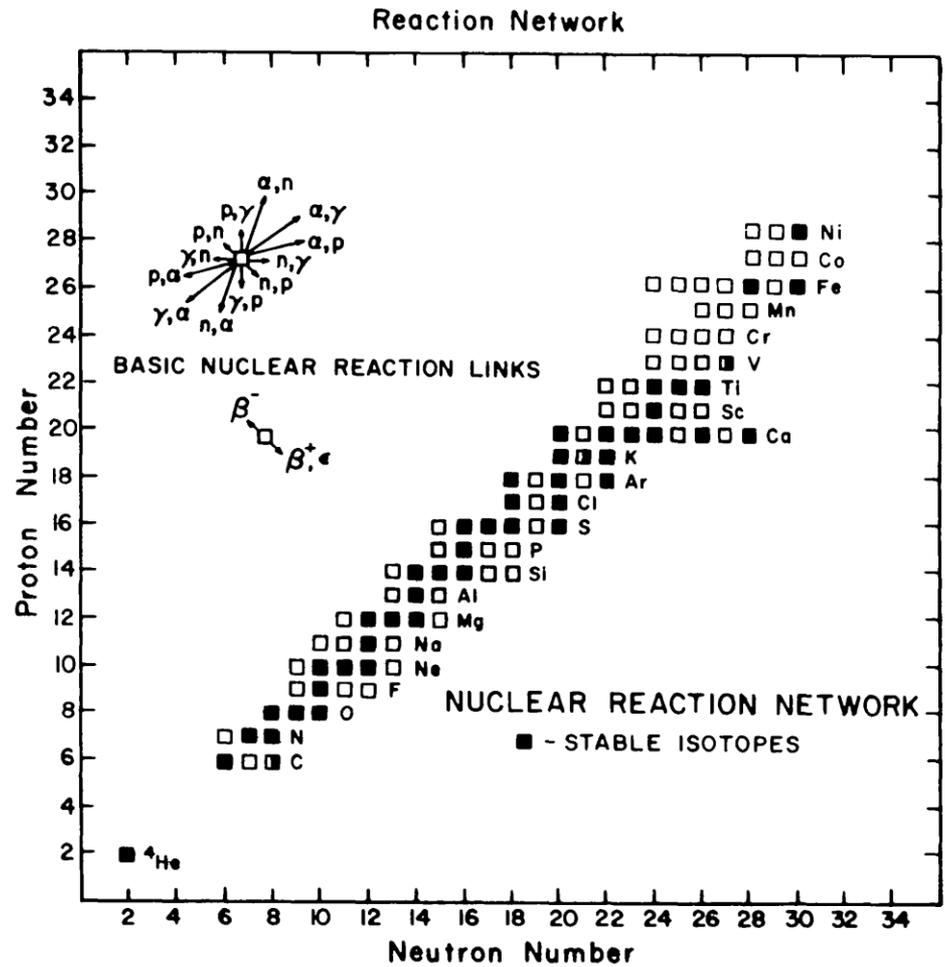
FIG. 11.—Variation of flows due to the weak interaction with time during carbon burning at $T_9 = 0.8$. Solid lines, β^+ decays; dashed lines, β^- decays.

Neon Burning

- neon and carbon burning have many similarities
- the main reaction is $^{20}\text{Ne}(\gamma, \alpha)^{16}\text{O}$
 - protons and neutrons are produced by (α, n) and (α, p) reactions
- behavior of abundances of neutrons, protons, and alpha particles is similar to that in carbon burning
- creates similar products to carbon burning, although ^{22}Ne and ^{23}Na are reduced in abundance
- carbon burning produces about three times more energy than neon burning

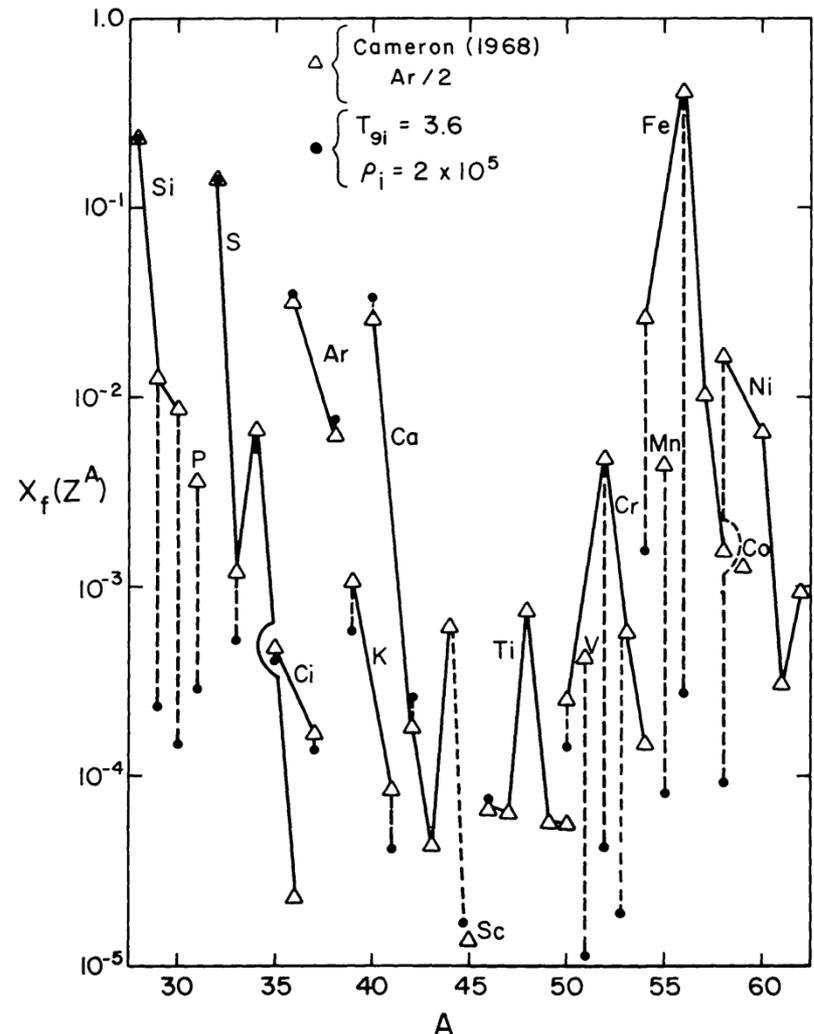
Oxygen Burning

- quasi-equilibrium behavior
 - less sensitive to reaction rates than carbon burning and more dependent on binding energies
- heavy elements ($Z > 28$) begin to photodissociate at the higher temperatures required for oxygen burning
- Woosley et al. 1973 did an analysis of explosive oxygen and silicon burning as an improvement on Truran and Arnett 1970, which was on explosive oxygen burning
 - initial composition for oxygen burning was about half ^{16}O and half ^{12}C
 - peak density was taken as $2 \cdot 10^5 \text{ g cm}^{-3}$
 - peak temperature was taken as $3.6 \cdot 10^9 \text{ K}$



Oxygen Burning

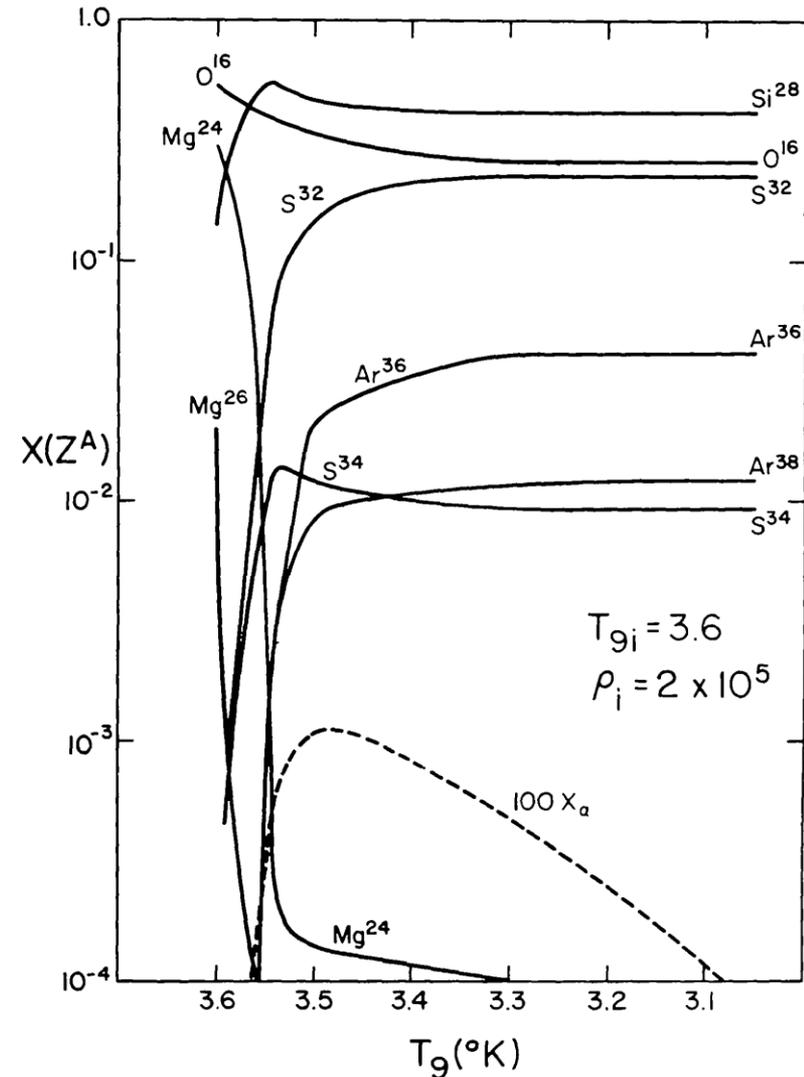
- calculation by Woosley succeeded in reproducing major abundances from ^{28}Si to ^{42}Ca , as well as the correct solar abundances of ^{46}Ti and ^{50}Cr
- calculations also agreed well with Truran and Arnett
- driving reactions:
 - $^{16}\text{O} (^{16}\text{O}, \alpha) ^{28}\text{Si}$
 - $^{16}\text{O} (^{16}\text{O}, p) ^{31}\text{P}$
 - $^{16}\text{O} (^{16}\text{O}, n) ^{31}\text{S}$
 - $^{16}\text{O} (^{16}\text{O}, d) ^{30}\text{P}$
 - $^{16}\text{O} (\gamma, \alpha) ^{12}\text{C}$
 - higher temperature and lower density favor this channel



Silicon Burning

- silicon burning also shows quasi-equilibrium behavior
- if ^{16}O is depleted before freezeout, nuclear evolution will become characterized by photodisintegration rearrangement reactions
- produces a sizable amount of iron peak nuclei
- overall silicon burning process:

$$^{28}\text{Si} \text{ (} ^{28}\text{Si, } 7\alpha \text{)} ^{56}\text{Ni}$$
- Woosley et al. start with initial conditions of temperature $4.7 \cdot 10^9$ K and density $2 \cdot 10^5$ g cm $^{-3}$
 - continue the silicon burning where the oxygen burning ended
 - end up with good approximation of solar abundances
- extreme silicon burning will merge with the e-process and approach nuclear statistical equilibrium



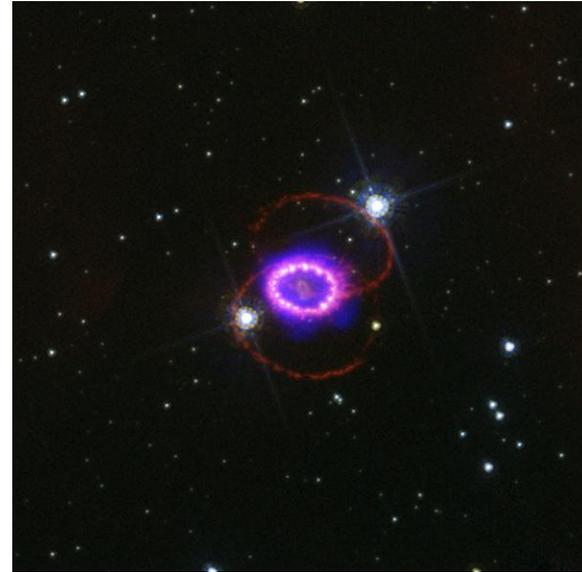
Observations of Nucleosynthesis in Core-Collapse Supernovae

○ Type II supernova SN 1987A

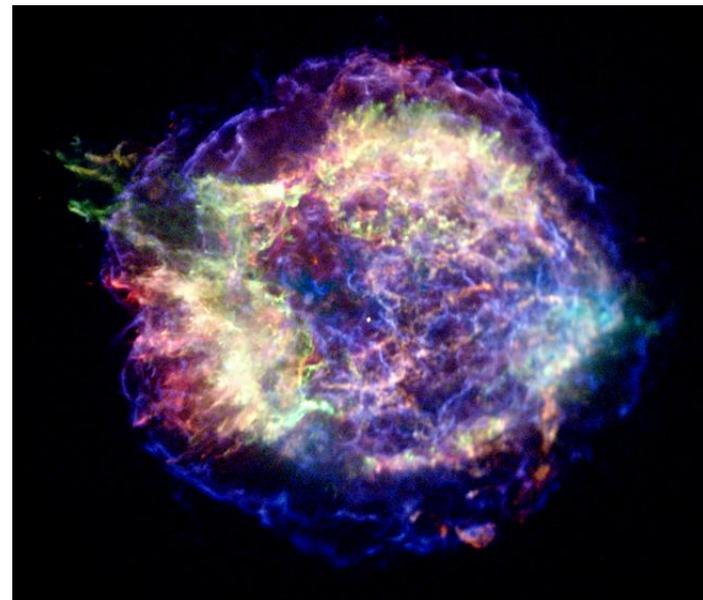
- light curve showed decline consistent with ^{56}Co
- gamma ray lines from ^{56}Co detected at 847 and 1238 keV
- confirmed by balloon-borne experiments
- line profiles of ^{56}Co imply that the supernova explosion was not spherically symmetric
- ^{57}Co decay detected

○ type II supernova Cassiopeia A

- obtained abundance estimates of S, Ar, Ca, Fe, and Ne, and upper limits for H, He, N, Mg, and C
- displayed expected quasi-equilibrium pattern expected of oxygen burning



*Credit: X-ray:
NASA/CXC/PSU/
S.Park &
D.Burrows. Optical:
NASA/STScI/CfA/P.
Challis*



*Credit:
NASA/CXC/MIT/U
Mass Amherst/
M.D.Stage et al.*

Simulations of Core-Collapse Supernovae

- although we do not know the precise mechanism by which a core-collapse supernova proceeds, nucleosynthesis in a core collapse supernova can be computed by artificially simulating a supernova explosion
 - injection of momentum through a piston that moves inward during the infall prior to the explosion and outward during the explosion
 - does not allow material in the ejecta to slow down and collapse back onto the star
 - underestimates fallback, leading to an overestimate of the amount of heavy elements, such as ^{56}Ni
 - “thermal bomb” – injection of thermal energy into the Fe core in a way such that the ejecta attains a kinetic energy of $\sim 10^{51}$ erg
 - designed to incorporate the energy increase in the convective region

Thermonuclear Supernovae (SNIa)

- probable mechanism for SNIa: delayed detonation
 - burning is subsonic (deflagration) in the inner core where densities are large and supersonic (detonation) in the outer core
 - pure detonation at densities larger than $\sim 10^7 \text{ g cm}^{-3}$ would cause the entire star to be incinerated to Ni
 - deflagration overproduces some neutronized isotopes, such as ^{54}Fe , ^{54}Cr , and ^{58}Ni
- ^{56}Ni lines (158, 750, 812 keV) prominent during first few days
 - 158 keV line is the most interesting to discriminate between models
 - almost undetectable in pure deflagration models
 - strong in pure detonation models
- ratio between fluxes of ^{56}Co 847 keV line (200 days after maximum) and 158 keV line (at maximum)
 - provides information about the ratio between ^{56}Ni in the ejecta and the external layers

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