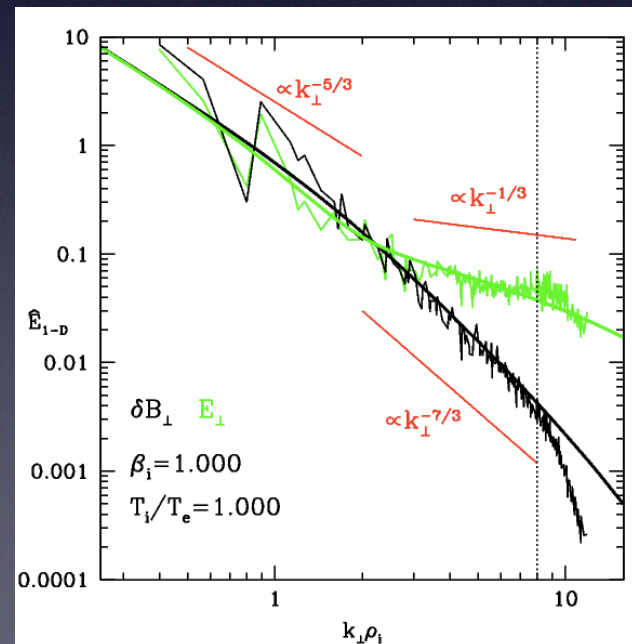
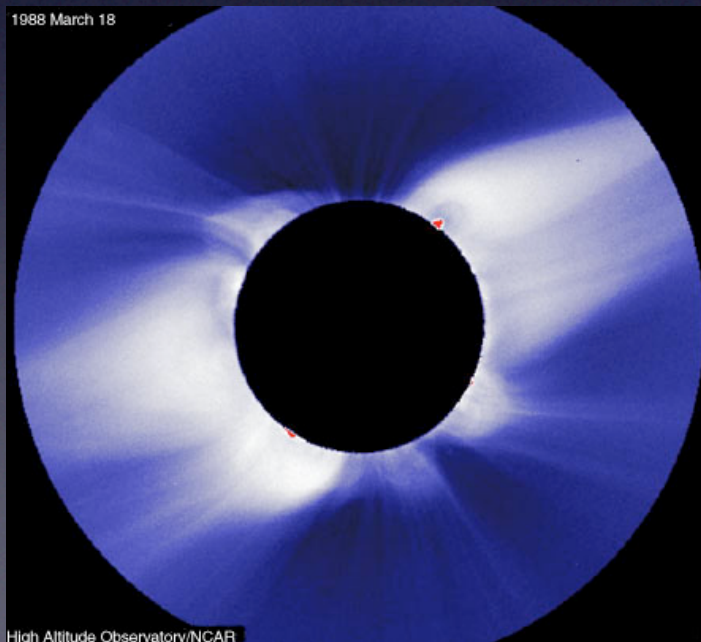


# The Heating & Acceleration of the Solar Wind

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Collaborators: Steve Cowley (UCLA), Bill Dorland (Maryland),  
Greg Hammett (Princeton), Greg Howes (Berkeley), Alex Schekochihin (Imperial)



# Overview

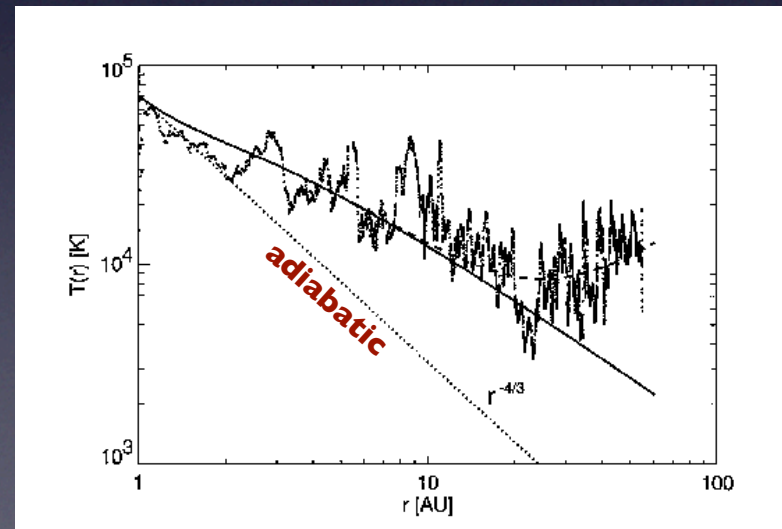
- Brief Observational & Theoretical Background
- Alfvénic Turbulence Theory (weak & strong)
  - Comparison to *In Situ* Observations at  $\sim$  AU
  - Transition to Kinetic Alfvén Wave Cascade at  $\sim$  the Ion Larmor Radius
- Particle Heating by Alfvénic Turbulence
  - Comparison to the Fast & Slow Winds
  - The Puzzle of the High Frequency Cascade (or the lack thereof ....)
  - Possible Solutions



# Background

- Heating required to accelerate the solar wind  
Parker 1958
- Early models invoked  $e^-$  conduction but  $T_p \gtrsim T_e$  in fast wind
- Local ( $r \sim R_\odot$ ) & extended ( $r \sim \text{few-}10^3 R_\odot$ ) heating required
- Extended heating favors waves
- Alfvén waves: primary observed fluctuation & least damped MHD mode in collisionless plasmas  
e.g., Belcher & Davis 1971; Barnes 1956

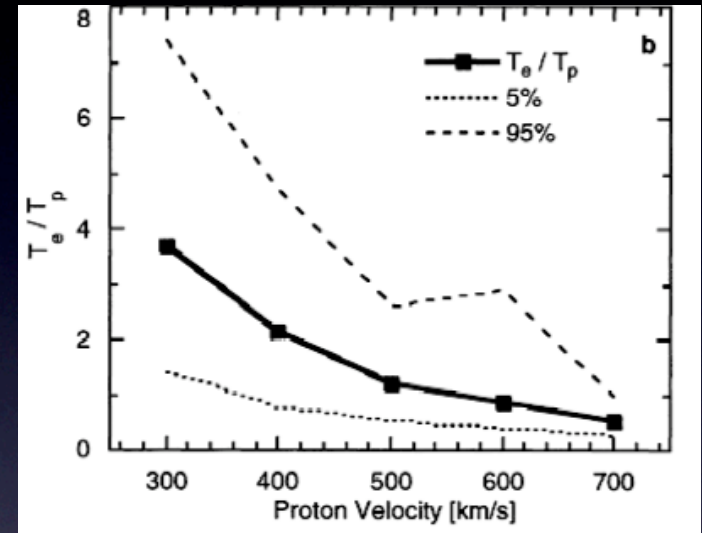
Voyager Temp Profile



Matthaeus et al. 1999

# Thermodynamic Constraints on Heating

- *In situ*: must dist. btw. Fast & Slow Wind
- **Fast**:  $T_{\text{ion}} \gtrsim T_p \gtrsim T_e$  &  $T_{\perp,i} \gtrsim T_{\parallel,i}$
- **Slow**:  $T_e \gtrsim T_p$  &  $T_{\parallel,i} \gtrsim T_{\perp,i}$  (?)



Newbury et al. 1998

- $\sim 1-4 R_{\odot}$ : constraints from UVCS/SOHO (in Coronal Holes = **Fast**)
- $T_{\perp,i} \gg T_{\parallel,i}$  (e.g.,  $O^{5+}$ , p)
- $T_i \gg T_p \gtrsim T_e$ ; preferential minor ion heating

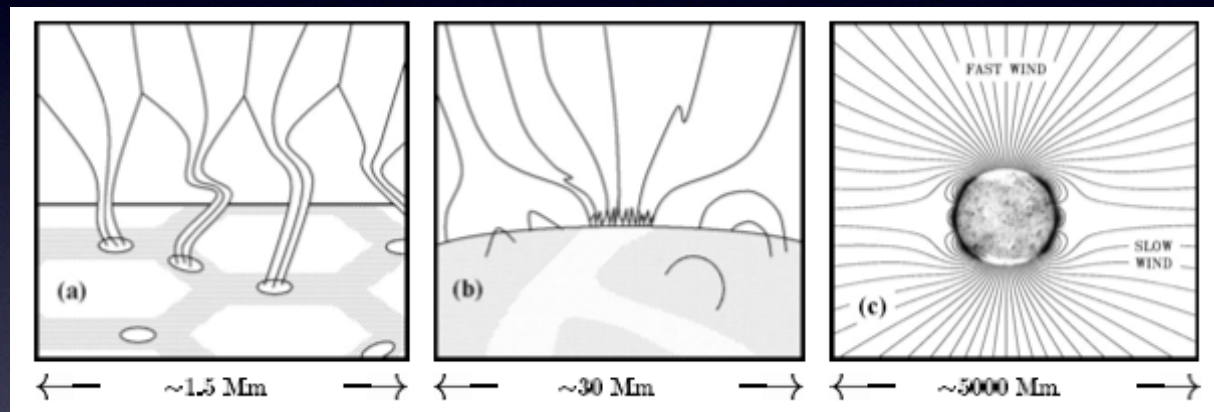
Kohl et al. 1997, 1998; Cranmer et al. 1999

**suggests  
ion cyclotron  
resonant heating**



# Wave Excitation/Launching

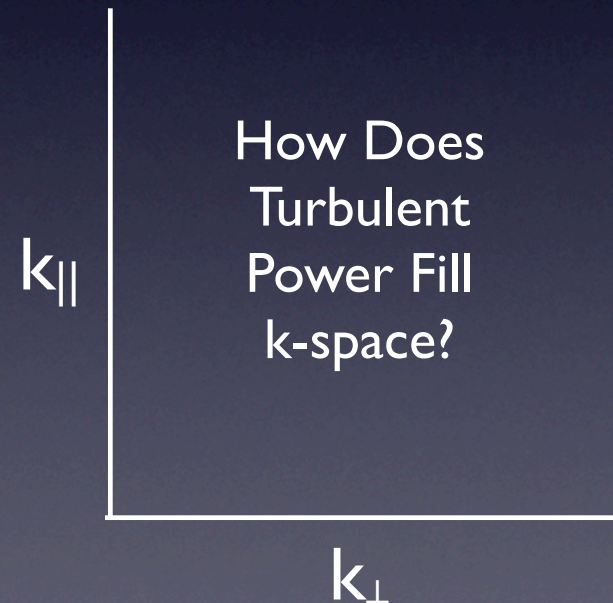
- **Small-scale Magnetic Activity → High Freq. Alfvén Waves**  
Axford & McKenzie 1992
  - ~ Hz and higher;  $f^{-1}$  spectrum often assumed
  - damp by ion cyclotron resonance: lower freq. waves damp at larger  $r$  (lower  $B$ )



- **Photospheric/Convective Motions → Low Freq. Alfvén Waves**  
e.g., Mattheus et al. 1999; Cranmer & van Ballegoijen 2005
  - ~ min & shorter
  - damp by **turbulent cascade** to small scales/high frequency

# MHD Turbulence

- Hydro:  $P(k) \sim k^{-5/3}$
- MHD: B-field defines local direction
  - $k = ??$ ;  $P(k) \sim k^{-??}$
- Focus on Incompressible MHD
  - Slow & Alfvén waves
  - Balanced Turbulence





# Incompressible MHD Turbulence

- View as interaction of Alfvén wave packets traveling at  $v = \pm v_A$  ( $\omega = |k_{\parallel}|v_A$ )  
e.g., Kraichnan 1965
- a single Alfvén wave packet is an exact non-linear soln of incompressible MHD  
→ **turbulence requires oppositely directed waves**
- solar wind: inward propagating waves generated by reflection of long-wavelength ( $\gtrsim$  density scale-height) outward propagating waves  
e.g., Matthaeus et al. 1999; Cranmer & van Ballegoijen 2005; Verdini & Velli 2007

- **weak** turbulence: non-linear (cascade) timescale  $\gg$  linear wave period

$$\omega_{nl} \ll \omega_{lin}$$

- **strong** turbulence: non-linear (cascade) timescale  $\sim$  linear wave period

$$\omega_{nl} \sim \omega_{lin}$$

# Weak MHD Turbulence

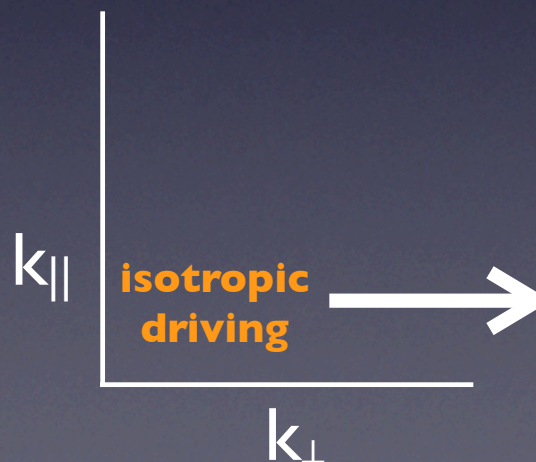
Shebalin et al. 1983; Goldreich & Sridhar 1995, 1997; Ng & Bhattacharjee 1996, 1997; Galtier et al. 2000

- non-linear time  $\gg$  linear wave period  $\sim (|k_{||}| v_A)^{-1}$
- Momentum & Energy Conservation  $\rightarrow$

$$\vec{k}_1 + \vec{k}_2 = \vec{k} \quad \omega_1 + \omega_2 = \omega$$

$$\rightarrow k_{||,1} - k_{||,2} = k_{||} \quad \& \quad k_{||,1} + k_{||,2} = k_{||}$$

- **$k_{||}$  cannot increase: energy flows in the perp. direction**





# Strong MHD Turbulence

Higdon 1984; Goldreich & Sridhar 1995

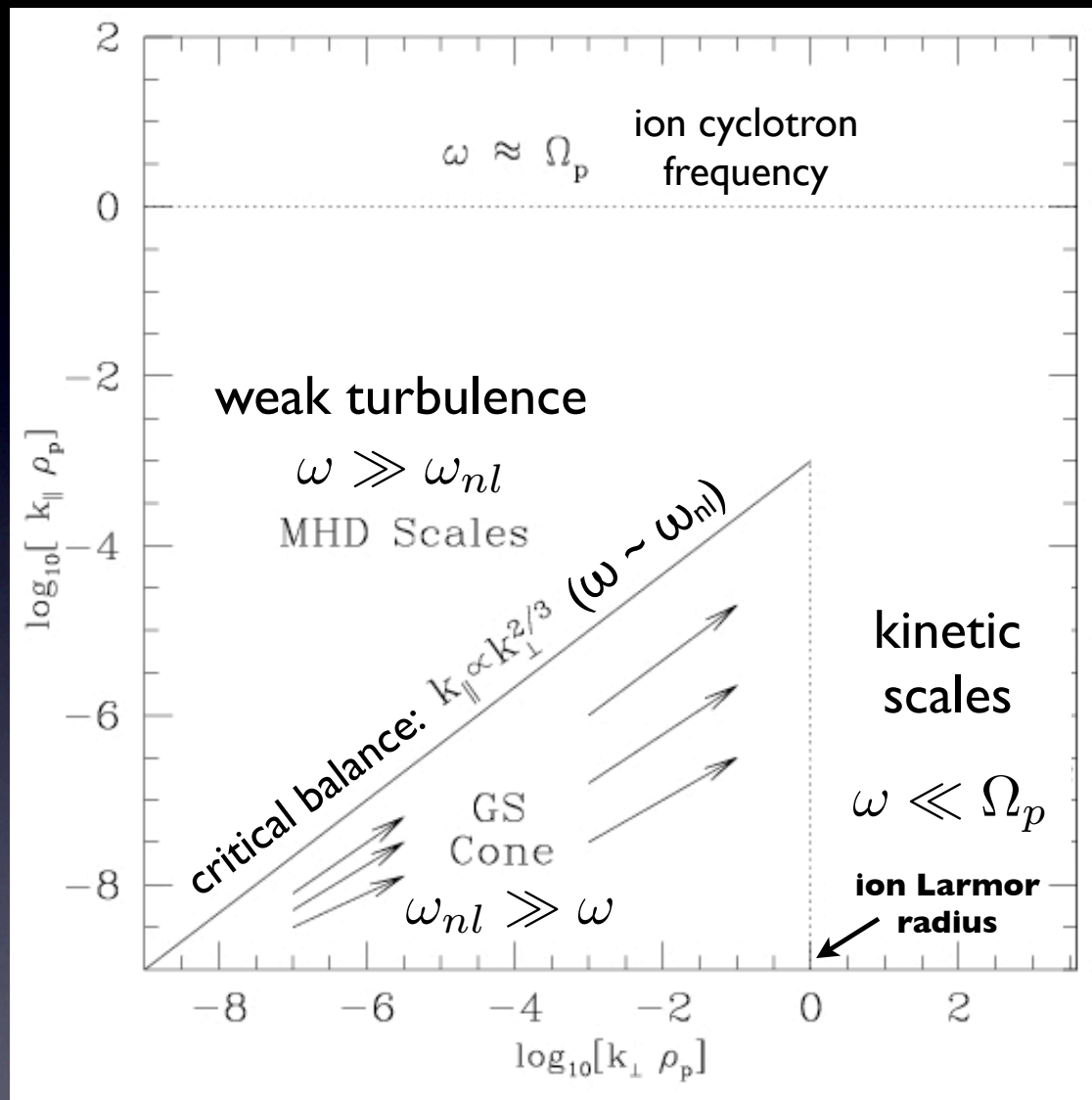
- non-linear interactions  $\sim (\mathbf{v} \cdot \nabla) \mathbf{v}$
- $\omega_{\text{nl}} \sim k_{\perp} \delta v_{\perp} \uparrow$  during weak turb.;  $\omega_{\text{lin}} = |k_{\parallel}| v_A$  unchanged
- **weak turbulence becomes strong:  $\omega_{\text{nl}} \sim \omega_{\text{lin}}$**
- “critical balance”: assume turbulence maintains  **$\omega_{\text{nl}} \sim \omega_{\text{lin}}$**   
Goldreich & Sridhar 1995

$$\rightarrow E(k_{\perp}) \propto k_{\perp}^{-5/3} \rightarrow \delta v_{\perp} \propto k_{\perp}^{-1/3}$$

Anisotropic  
Kolmogorov

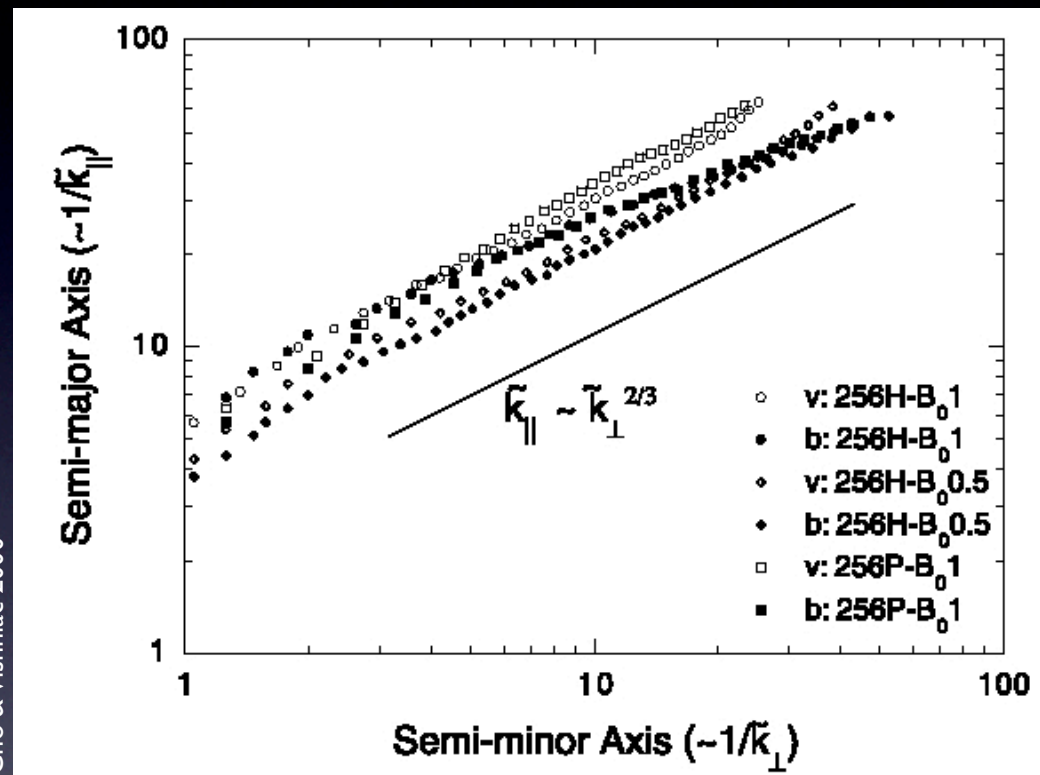
$$\text{critical balance} \rightarrow k_{\parallel} \propto k_{\perp}^{2/3}$$

Scale-Dependent  
Anisotropy





# MHD Simulations Support the Goldreich-Sridhar (GS) Model

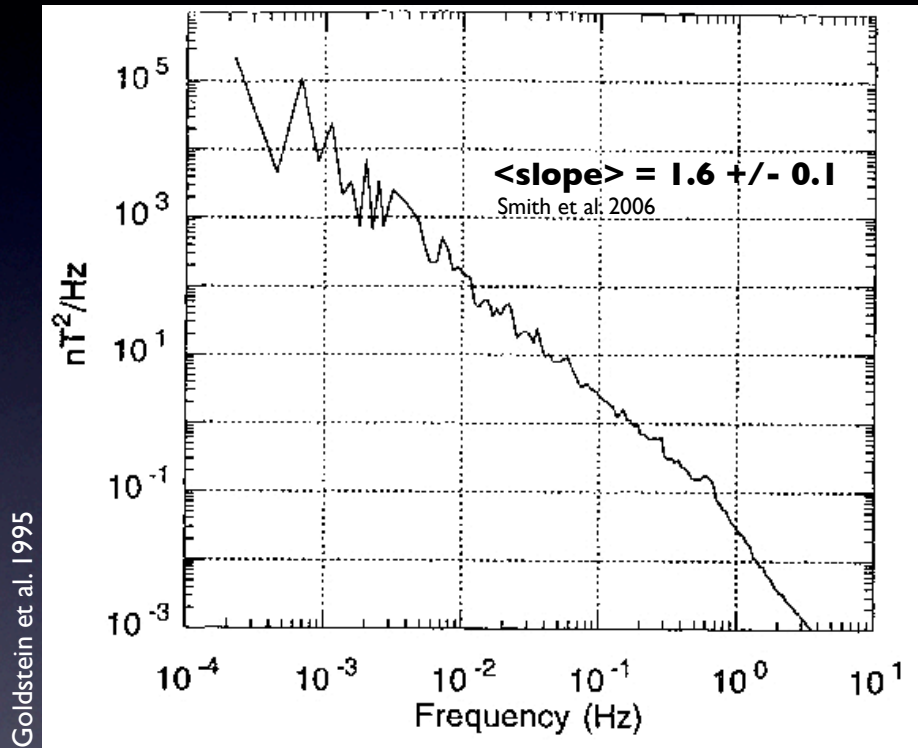


Compressible Sims show that Alfvén & Slow  
Modes Follow the GS Cascade

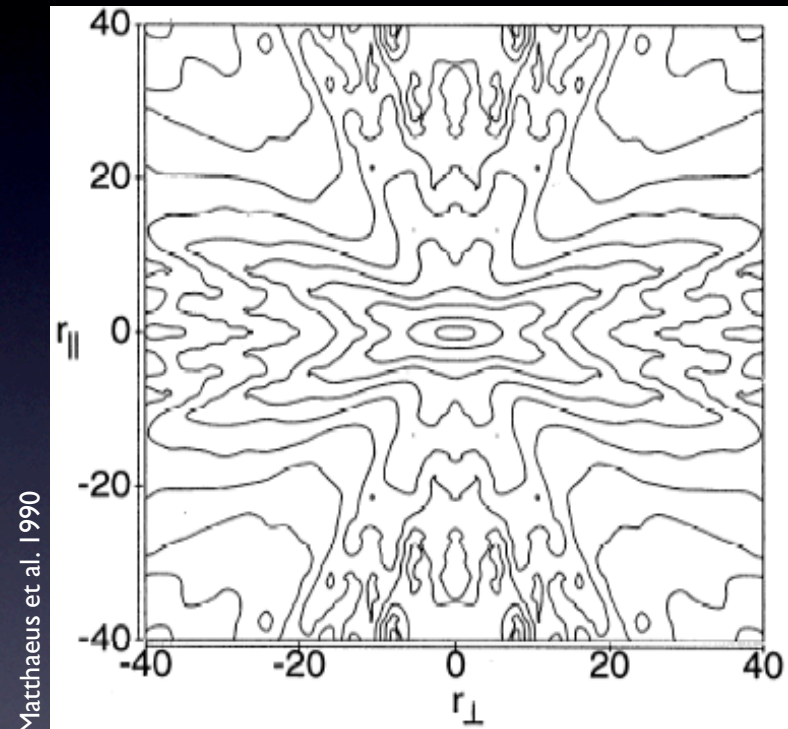
Some Fast Mode Energy Cascades to High Freq

Cho & Lazarian 2003; see also Chandran 2005

# Solar Wind Fluctuations



Magnetic field power spectrum  
consistent w/ Kolmogorov  
(above the ion Larmor radius)



~ 90% of the Energy in  $\perp$  fluctuations

~ 10% in  $\parallel$  fluctuations

slow wind: more  $\perp$  fluctuations

fast wind: more  $\parallel$  fluctuations

Dasso et al. 2005

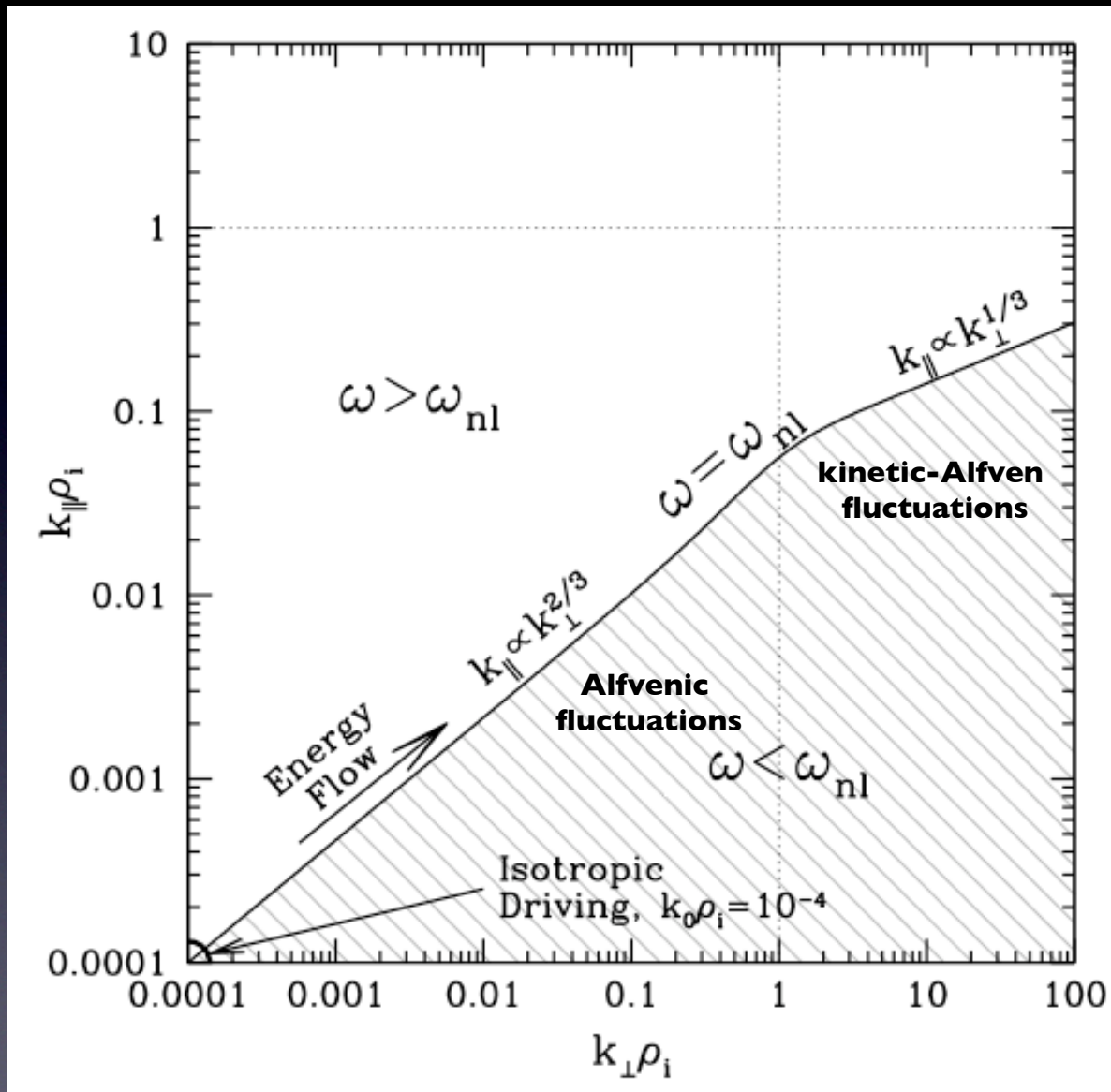


# Towards the Dissipation Range: The Transition to a Kinetic Alfvén Wave Cascade at $\sim \rho_i$

at  $k_{\perp} \rho_i \simeq 1$ ,  $\frac{\omega}{\Omega_i} \simeq \left(\frac{\rho_i}{L}\right)^{1/3} \beta_i^{-1/2}$   $L \equiv$  outer scale of turbulence

- **Solar Wind at 1 AU:**  $\omega/\Omega_i \simeq 0.04$  at  $k_{\perp} \rho_i \simeq 1$  ( $L \simeq 10^{11}$  cm)
- **Corona at  $\sim 2 R_{\odot}$ :**  $\omega/\Omega_i \simeq 0.03$  at  $k_{\perp} \rho_i \simeq 1$  ( $L \simeq 10^9$  cm)  
(fluctuations already anisotropic at the outer scale)
- $k_{\perp} \rho_i \gtrsim 1$  &  $\omega \lesssim \Omega_i$ , Alfvén waves  $\rightarrow$  Kinetic Alfvén Waves (KAWs)

strong Alfvén wave turbulence  $\rightarrow$  strong KAW turbulence



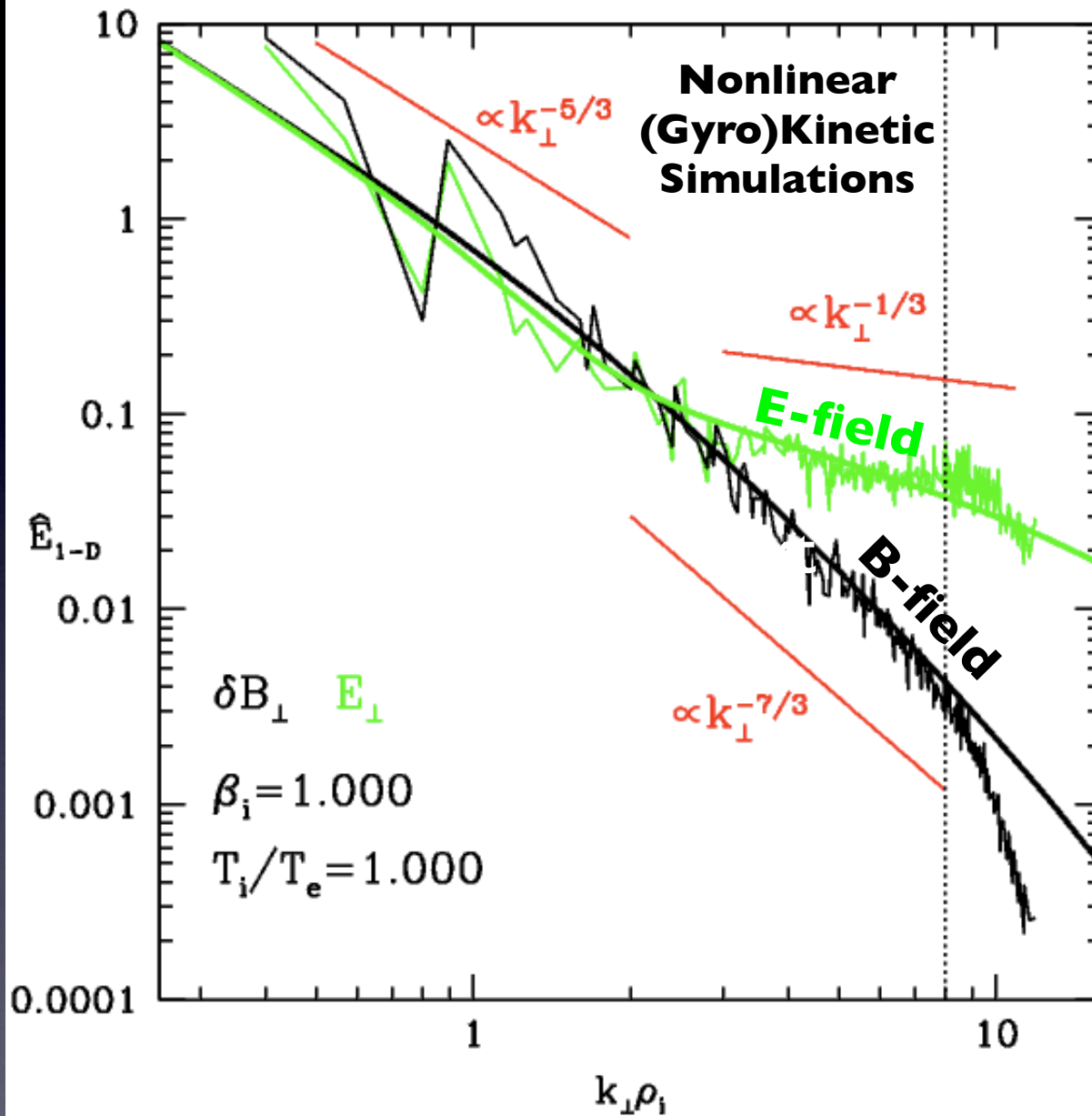
## Strong KAW Turbulence (sans damping)

$$E_B \propto k_{\perp}^{-7/3}$$

$$k_{\parallel} \propto k_{\perp}^{1/3}$$

Biskamp et al. 1999;  
 Cho & Lazarian 2004;  
 Schekochihin et al. 2007





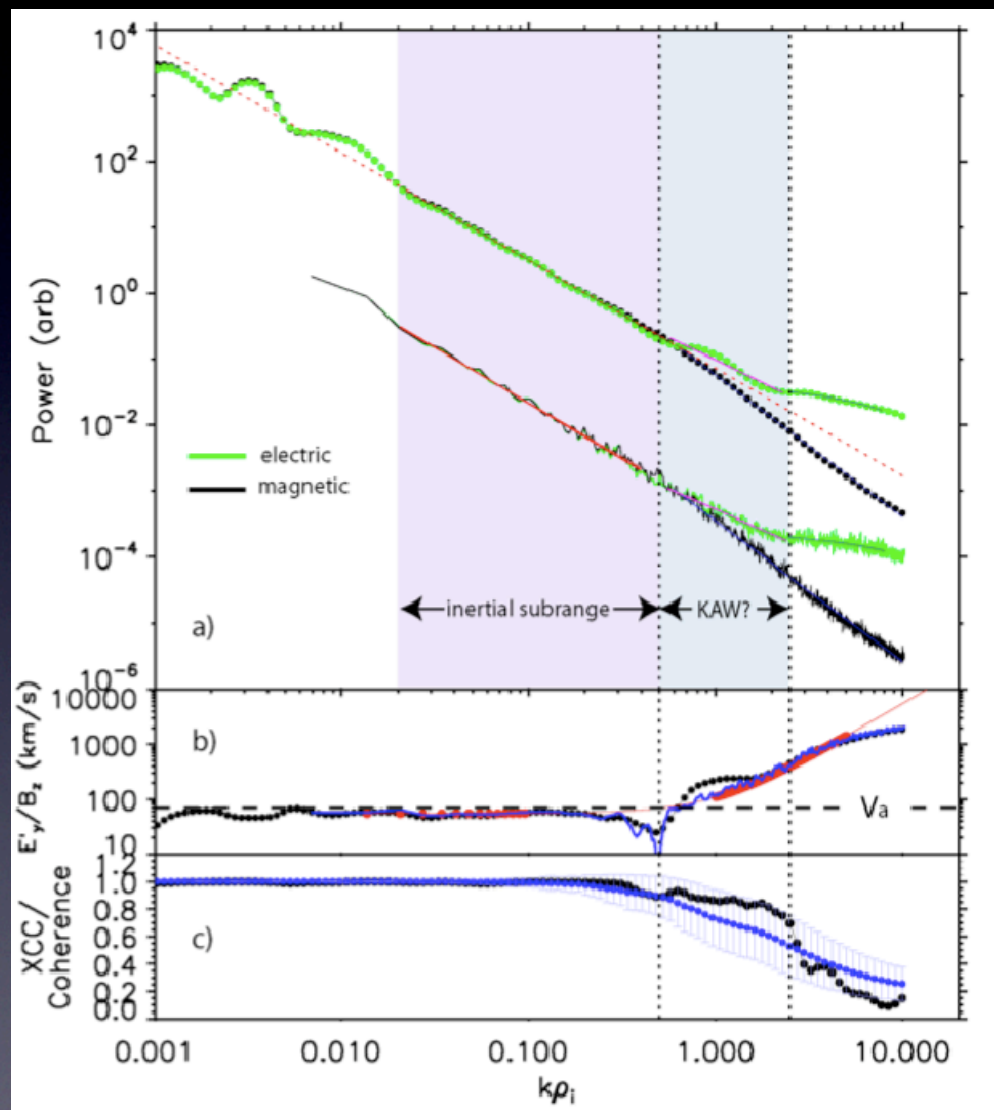
anisotropic  
low frequency  
turbulence  
**both above & below**  $\rho_i$  can be  
quantitatively  
modeled using a  
low freq. expansion  
of the Vlasov eqn

Howes et al. 2006; Schekochihin et al. 2007

**“gyrokinetics”**

# In Situ Measurements in the Solar Wind

(Bale et al. 2005)



**In Situ Measurements  
of E & B-fields  
with Cluster are  
Consistent with  
a transition  
to KAWs at small  
scales but *not* with  
the onset of ion  
cyclotron damping**



# Collisionless Damping of the Anisotropic Cascade

Quataert 1998; Leamon et al. 1998; Quataert & Gruzinov 1999; Cranmer & van Ballegoijen 2003; Gary & Nishimura 2004

- so long as  $\omega \lesssim \Omega_i$ 
  - no cyclotron resonance
  - magnetic moment  $\mu \propto T_{\perp}/B$  is conserved
  - $\rightarrow$  heating can only increase  $T_{\parallel}$
- cyclotron damping is strongly suppressed at  $k_{\perp}\rho_i \gtrsim 1$ 
  - $\rightarrow$  for cycl. damping to be impt,  $\omega \rightarrow \Omega_i$  at  $k_{\perp}\rho_i \lesssim 1$

# Collisionless Damping of the Anisotropic Cascade

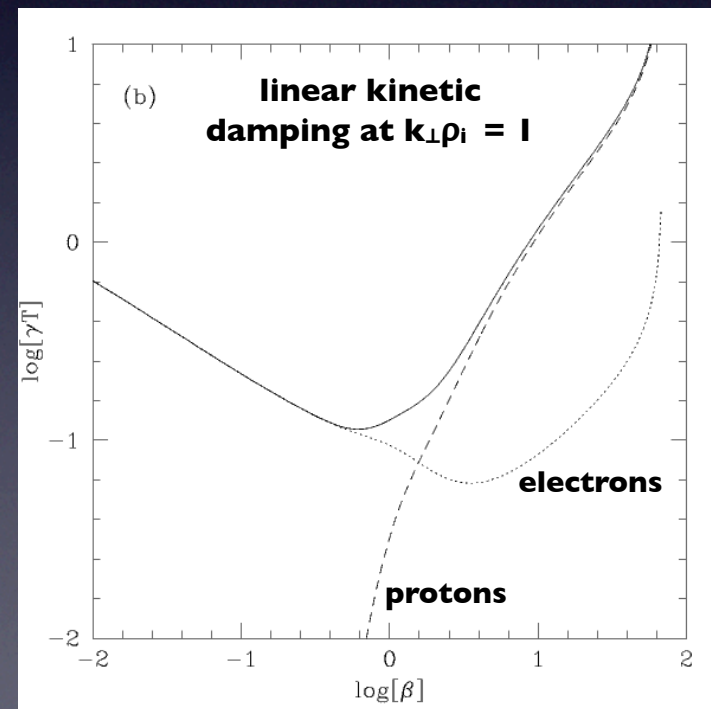
Quataert 1998; Leamon et al. 1998; Quataert & Gruzinov 1999; Cranmer & van Ballegoijen 2003; Gary & Nishimura 2004

- parallel heating via the Landau resonance:  $\omega = k_{\parallel} v_{\parallel}$
- both Landau damping ( $\delta E_{\parallel}$ ) & transit-time damping ( $\delta B_{\parallel}$ )  
 $\beta \lesssim 1$   $\beta \gtrsim 1$

- **primarily  $e^{-}$  heating for  $\beta \lesssim 10$**

- dominant source of  $e^{-}$  heating  
in solar wind (?); consistent with

$$T_e \gtrsim T_p \text{ in slow wind}$$





# The Puzzle ...

- How to get  $T_{\text{ion}} \gtrsim T_p \gtrsim T_e$  &  $T_{\perp,i} \gtrsim T_{\parallel,i}$ ? (Fast Wind)

- Outer scale  $\ll$  Assumed Values (unlikely ...?); Coupling of KAWs to Ion Bernstein, fast waves (unlikely)

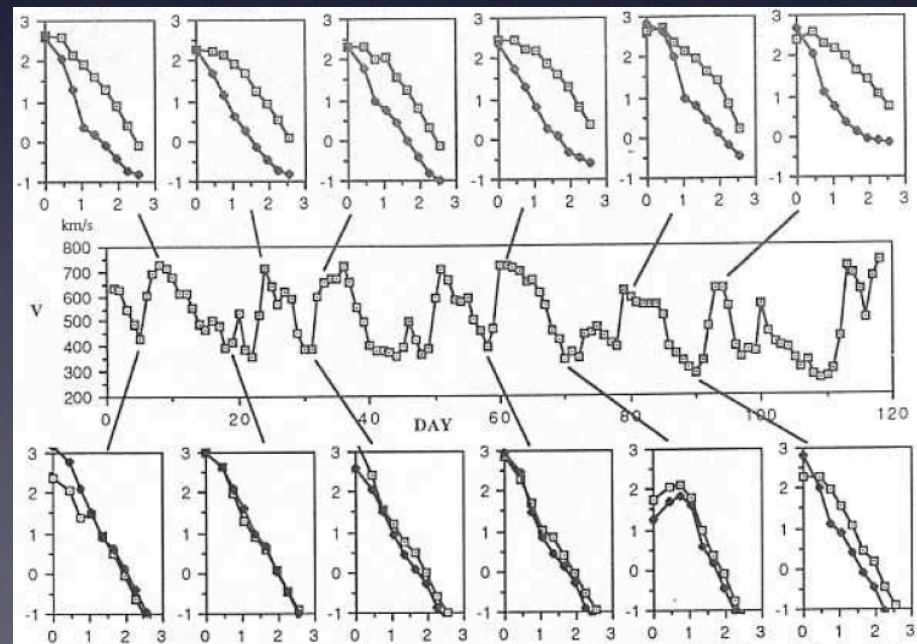
at  $k_{\perp} \rho_i \simeq 1$ ,  $\frac{\omega}{\Omega_i} \simeq \left(\frac{\rho_i}{L}\right)^{1/3} \beta_i^{-1/2}$   $L \equiv$  outer scale of turbulence

- Fast Waves Cascade to High Frequencies for  $\beta \ll 1$  (Chandran '05)

- but Alfvénic fluctuations dominate at  $\sim$  AU ...

- Imbalanced Turbulence?  
anisotropy is the same  
cascade slows down

→ **less** likely to reach  $\sim \Omega_i$

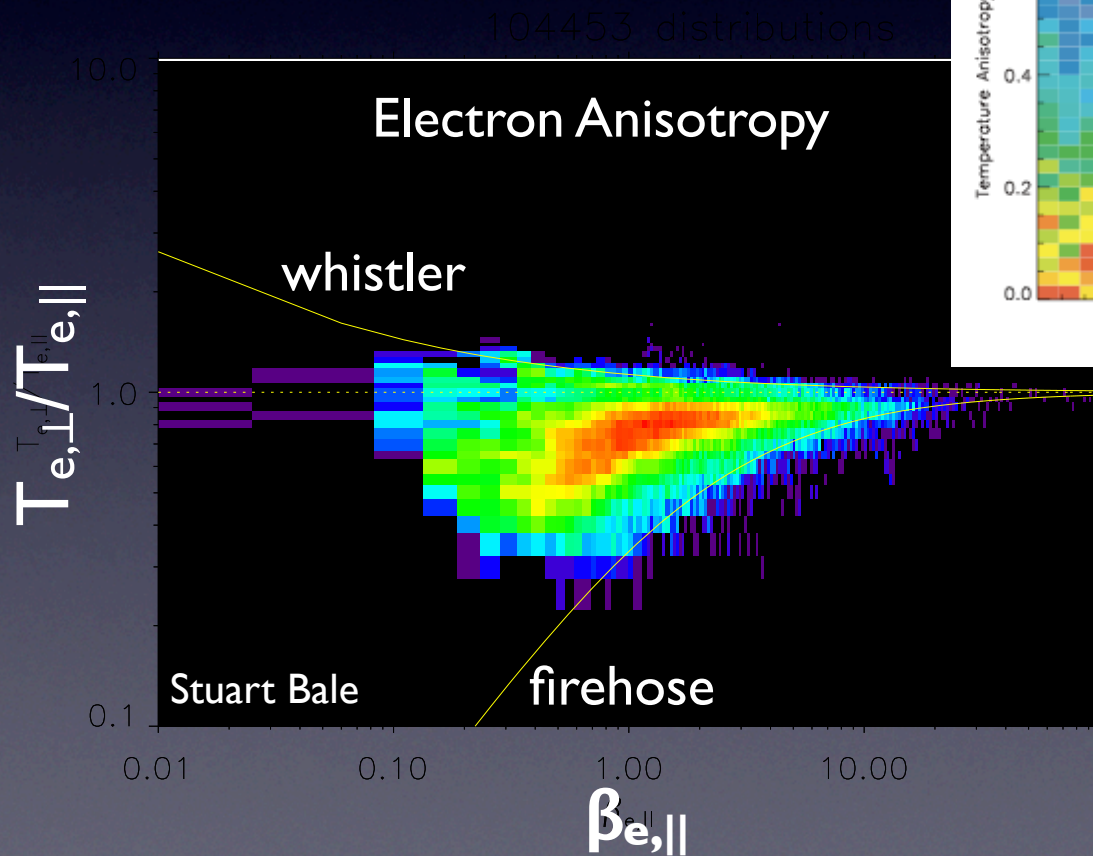
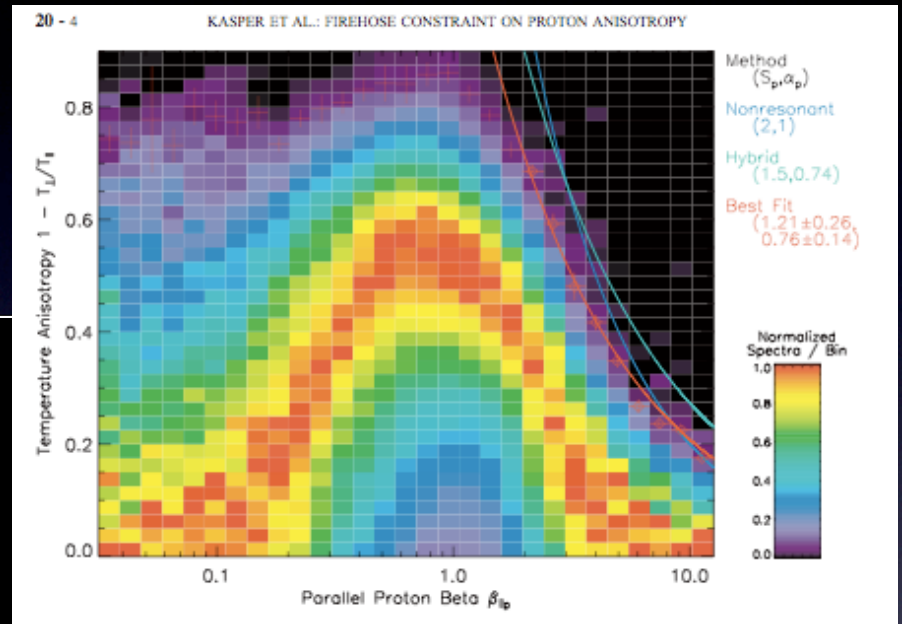


# The Puzzle ...

- How to get  $T_{\text{ion}} \gtrsim T_p \gtrsim T_e$  &  $T_{\perp,i} \gtrsim T_{\parallel,i}$ ? (Fast Wind)
  - Imbalanced Turbulence: cascade slows down
    - other non-linearities impt?
    - e.g., Alfvén waves steepen on a timescale  $\sim \omega_{\text{lin}}^{-1} (B/\delta B)^2$
  - Secondary Instabilities
    - electron-ion:  $T_e \rightarrow T_i$  (e.g., e- beams)  
Cranmer & van Ballegoijen 2003; Gary & Nishimura 2004
    - velocity space:  $T_{\perp} \leftrightarrow T_{\parallel}$



# Velocity-Space Instabilities



proton  $\parallel$  anisotropy  
limited by firehose  
Kasper et al. 2002

# Summary

- Alfvénic Turbulence is the Most Promising Source of Heating in the Extended Corona and Solar Wind
- Strong MHD Turbulence (Alfvénic)
  - ✓ Anisotropic Kolmogorov Turbulence: critical balance  $\rightarrow k_{\parallel} \propto k_{\perp}^{2/3}$
  - ✓  $k_{\perp} \rho_i \sim 1$ : Alfvén Wave Cascade  $\rightarrow$  Kinetic Alfvén Wave Cascade
    - **NOT** cyclotron damping:  $\omega \sim 0.03-0.2 \Omega_i$  even at  $k_{\perp} \rho_i \sim 1$
    - Confirmed by Cluster Electric Field Measurements
  - KAW Cascade  $\rightarrow$  Electron  $\parallel$  Heating at  $k_{\perp} \rho_i \sim 0.3-10$  ( $\beta \sim 10^{-3}-1$ )
- Puzzle:  $T_{\text{ion}} \gtrsim T_p \gtrsim T_e$  &  $T_{\perp,i} \gtrsim T_{\parallel,i}$  (Fast Wind)
  - Smaller Outer Scale? Fast Waves? Secondary Instabilities? Addtl non-linearities? Sweep + Cascade in  $\omega$ ? ...