

Boundary induced amplification and nonlinear instability of interchange modes

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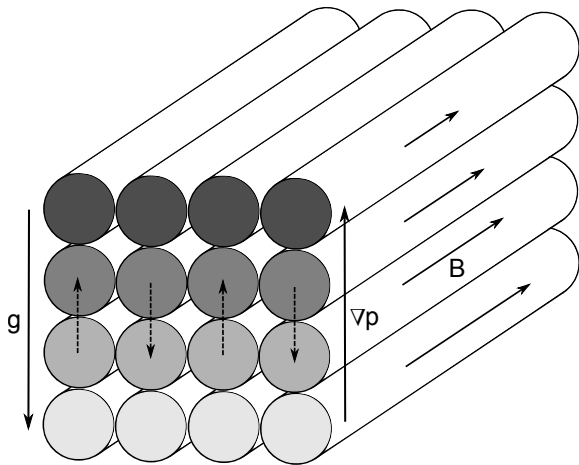
Stability, Energetics, and Turbulent Transport
in Astrophysical, Fusion, and Solar Plasmas

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Outline

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 - ▶ Ideal MHD interchange mode
 - ▶ Motivation
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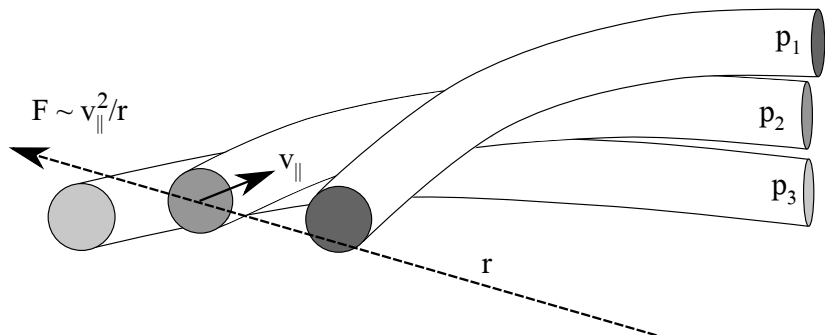
Interchange mode



Growth rate:

$$\gamma_g^2 \equiv g \frac{\rho_0'}{\rho_0}$$

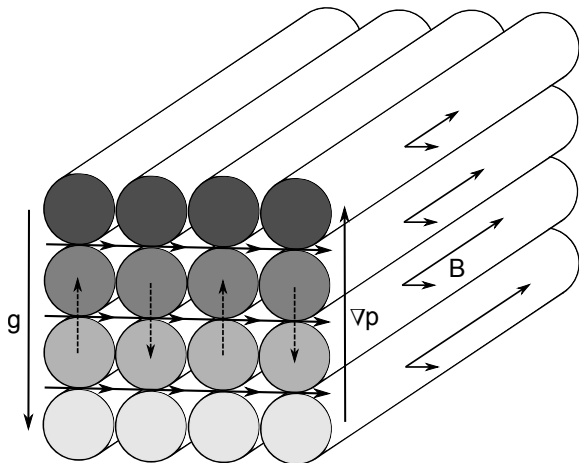
Toroidal device



Unstable when:

$$p_1 > p_2 > p_3$$

Transverse field stabilization



Dispersion relation:

$$\omega^2 = k^2 V_{Ay}^2 - \gamma_g^2$$

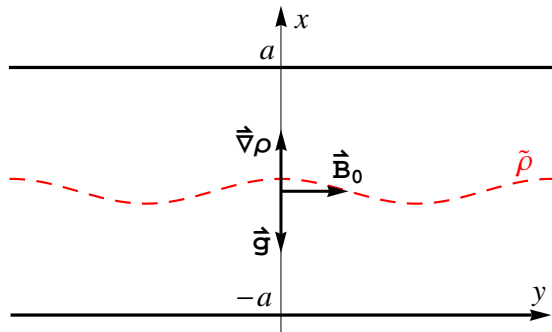
Motivation

Behaviour at marginal stability for ideal interchange modes is important in:

- ▶ Tokamaks,
- ▶ Reversed Field Pinches,
- ▶ Stellarators,
- ▶ etc...

Understanding the tradeoff between the deviation from marginality and residual convection is useful for B-field coil design in stellarators.

Previous result

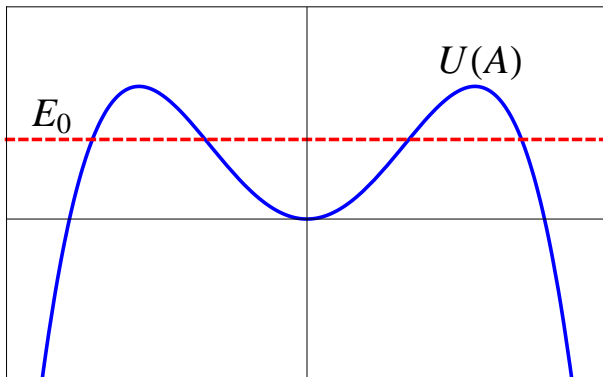


$$B_0 = B_c + b, \quad \tilde{\rho} \propto A \sim \frac{\partial_t}{\gamma_g} \sim \left| \frac{b}{B_c} \right|^{1/2} \ll 1$$

$$\boxed{\frac{1}{k^2} \ddot{A} = -2bA + \frac{1}{4} \frac{k^2 - 3}{k^2 + 1} A^3}$$

Amplitude dependent stability

$$E_0 = \frac{1}{2k^2} \dot{A}^2 + U(A), \quad U(A) = bA^2 - \frac{1}{16} \frac{k^2 - 3}{k^2 + 1} A^4$$



$$A_c = 2\sqrt{2} \sqrt{\frac{k^2 + 1}{k^2 - 3}} \times b^{1/2}$$

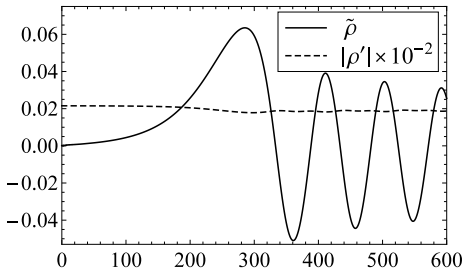
Numerical marginal conditions amplitude test

Linearly stable

$$\Delta B/B_c \approx 0.04\%$$

Small amplitude

$$a_0 = 10^{-4}$$

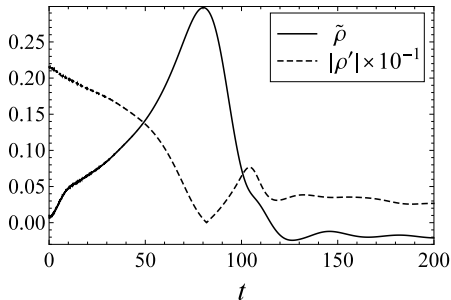


Linearly stable

$$\Delta B/B_c \approx 0.04\%$$

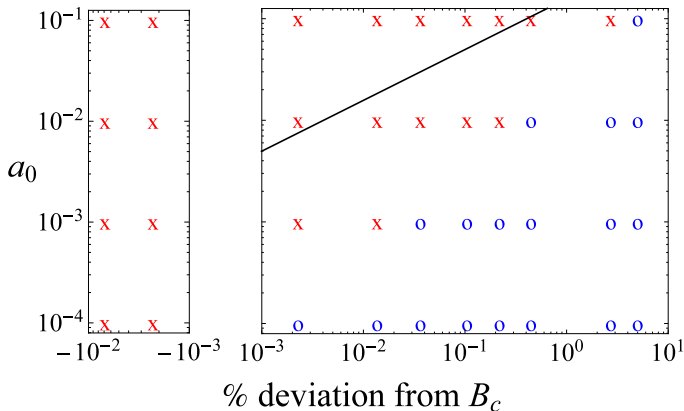
Large amplitude

$$a_0 = 10^{-2}$$



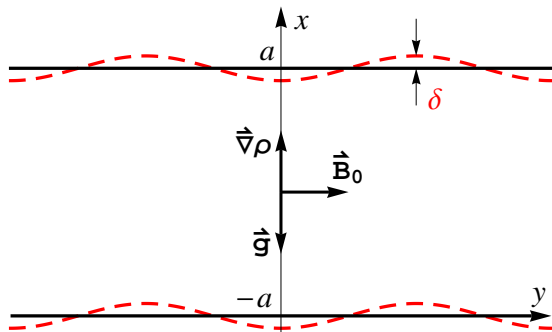
Numerical simulation result¹

Stable runs are marked with a blue circle and unstable runs with a red 'x'. The theoretical boundary is shown as the line $a_0 \propto \sqrt{b}$.



¹*Phys. Plasmas* **18** 122103 (2011).

Rippled boundary problem



$$\rho'_0 > 0 \text{ and } B_0 \text{ constant}$$

$$\frac{\delta}{a} \ll 1$$

Reduced equations in 2D

Complete, nonlinear set:

$$\begin{aligned}\partial_t \psi &= \{\psi, \varphi\} \\ \hat{\mathbf{z}} \cdot \vec{\nabla}_\perp \times \rho (\partial_t \vec{\mathbf{u}} + \{\varphi, \vec{\mathbf{u}}\}) &= \{\psi, \nabla_\perp^2 \psi\} + g \partial_y \rho \\ \vec{\mathbf{u}} &= \hat{\mathbf{z}} \times \nabla_\perp \varphi \\ \vec{\mathbf{B}}_\perp &= \hat{\mathbf{z}} \times \nabla_\perp \psi\end{aligned}$$

Where, in general, $\rho = \rho(\psi)$ and

$$\{f, h\} \equiv \partial_x f \partial_y h - \partial_y f \partial_x h.$$

“Simple-minded” linear calculation

Boundary condition:

$$\text{At } x = \pm a - \delta \cos(ky)$$

$$\partial_y \psi = -\delta k \sin(ky) \partial_x \psi$$

Quasistatic equilibrium:

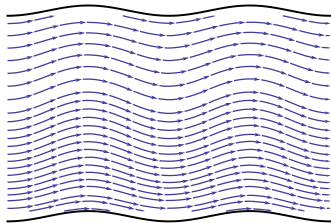
$$\psi = B_0 x + \tilde{\psi}$$

$$(B_0^2 \nabla_{\perp}^2 + g\rho'_0) \partial_y \tilde{\psi} = -B_0 \{ \tilde{\psi}, \nabla_{\perp}^2 \tilde{\psi} \}$$

To lowest order in δ/a :

$$\tilde{\psi} = \frac{\delta B_0}{\cos(k_x a)} \cos(k_x x) \cos(ky)$$

$$k_x^2 = \frac{g\rho'_0}{B_0^2} - k^2$$



Boundary induced amplification

Critical condition:

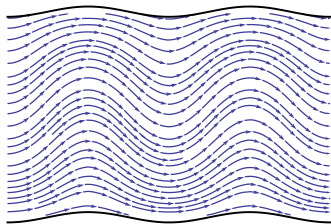
$$B_0 \rightarrow B_c \text{ when } k_x \rightarrow k_c \equiv \pi/2a$$
$$B_c^2(k_c^2 + k^2) \equiv B_c^2 k_{\perp}^2 = g\rho'_0$$

Amplification scaling:

$$B_0 = B_c + b \text{ and } k_x = k_c - \Delta k_x$$
$$b/B_c \sim \Delta k_x/k_c \ll 1$$

To lowest order in b/B_c :

$$\tilde{\psi} \approx \frac{\delta/a}{b/B_c} \frac{k_c}{k_{\perp}^2} B_c \cos(k_x x) \cos(ky)$$



Core penetration

Dirac delta gradient:

$$\rho'_0(x) = \rho'_0 a \delta(x)$$

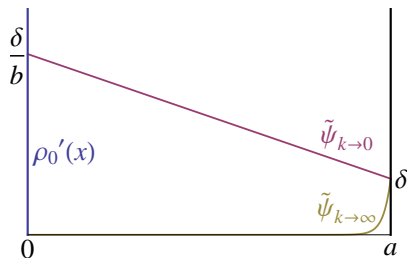
Solution:

$$\tilde{\psi} = B_0 \delta \frac{ka \cosh(kx) - \frac{1}{2} \kappa^2 a^2 \sinh(k|x|)}{ka \cosh(ka) - \frac{1}{2} \kappa^2 a^2 \sinh(ka)} \cos(ky), \quad \kappa^2 \equiv \frac{g\rho'_0}{B_0^2}$$

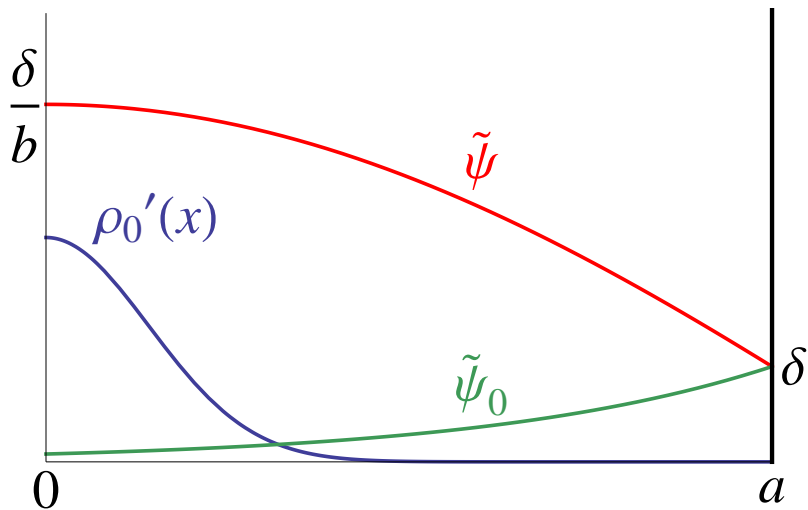
Limits:

$$\tilde{\psi}_{k \rightarrow 0} \propto B_0 \delta \frac{1 - \frac{1}{2} \kappa^2 a |x|}{1 - \frac{1}{2} \kappa^2 a^2}$$

$$\tilde{\psi}_{k \rightarrow \infty} \propto 2B_0 \delta e^{-ka} \cosh(kx)$$



Global amplification



Nonlinear scaling and expansion

Previous scaling:

$$|\tilde{\psi}/\psi_0| \sim \epsilon \equiv (b/B_c)^{1/2}$$
$$\Rightarrow \frac{\delta}{a} \sim \epsilon^3$$

Expand ψ in ϵ :

$$\psi = \psi_0 + \psi_1 + \psi_2 + \psi_3 + \dots$$

$$\psi_0 = (B_c + b)x$$

$$\frac{\psi_{i+1}}{\psi_i} \sim \epsilon$$

First order

Equation:

$$B_c^2 (\nabla_{\perp}^2 + k_{\perp}^2) \partial_y \psi_1 = 0$$
$$\partial_y \psi_1 |_{x=\pm a} = 0$$

Solution:

$$\psi_1 = \left[A + \frac{\delta/a}{b/B_c} \frac{k_c}{k_{\perp}^2} B_c \right] \cos(k_c x) \cos(ky)$$

Where A is the plasma response.

$$\frac{A}{B_c/k_c} \sim \epsilon$$

Second order

Equation:

$$B_c^2(\nabla_{\perp}^2 + k_{\perp}^2)\partial_y\psi_2 = -B_c\{\psi_1, \nabla_{\perp}^2\psi_1\} \\ = 0$$

$$\partial_y\psi_2|_{x=\pm a} = 0$$

Solution:

$$\psi_2 = -\frac{1}{4}\frac{k_c}{B_c}\left[A + \frac{\delta/a}{b/B_c}\frac{k_c}{k_{\perp}^2}B_c\right]^2 \sin(2k_c x)$$

Third order

Equation:

$$B_c^2(\nabla_{\perp}^2 + k_{\perp}^2)\partial_y\psi_3 + 2B_cb\nabla_{\perp}^2\partial_y\psi_1 = -B_c(\{\psi_1, \nabla_{\perp}^2\psi_2\} + \{\psi_2, \nabla_{\perp}^2\psi_1\})$$
$$\partial_y\psi_3|_{x=\pm a} = -\delta B_ck \sin(ky)$$

Secular term:

$$\psi_3 = B_c \frac{\delta}{a} x \sin(k_c x) \cos(ky)$$
$$-2B_cbA + \frac{k_c^2}{4} \frac{k^2 - 3k_c^2}{k_{\perp}^2} \left[A + \frac{\delta/a}{b/B_c} \frac{k_c}{k_{\perp}^2} B_c \right]^3 = 0$$

Time evolution

Time dependence:

$$A = A(t) \text{ where } \tau_A \partial_t \sim \epsilon$$

keep $\dot{\delta}$ "small"

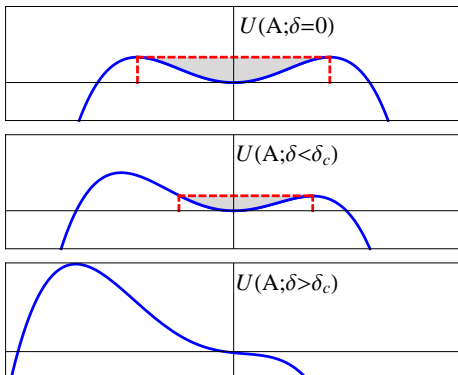
Result:

$$\frac{1}{k^2} \ddot{A} = -2bA + \frac{1}{4} \frac{k^2 - 3}{k^2 + 1} \left[A + \frac{2}{\pi} \frac{\delta}{b} \frac{1}{k^2 + 1} \right]^3$$

Normalized with k_c , B_c , and V_{Ac} equal to 1.

Boundary induced nonlinear instability²

$$U(A; \delta) = bA^2 - \frac{1}{16} \frac{k^2 - 3}{k^2 + 1} \left[A + \frac{2}{\pi} \frac{\delta}{b} \frac{1}{k^2 + 1} \right]^4$$



$$\delta_c = \frac{\pi}{2} \sqrt{\frac{32}{27} \frac{(k^2 + 1)^3}{k^2 - 3}} \times b^{3/2}$$

²Phys. Plasmas (Letters) **20** 020704 (2013).

Summary and conclusions

In marginally stable interchange mode systems boundary perturbations...

1. ...are amplified:

$$\tilde{\psi} \propto \delta/b$$

2. ...induce nonlinear instability.

Therefore, marginal system is highly sensitive to boundary perturbations:

$$\delta_c \propto b^{3/2}$$

Future work

Energy principle generalization:

- ▶ δW with nonlinear effects
- ▶ δW with boundary perturbations

Realistic geometry (shear, curvature, etc)