

3D GRMHD Jet Simulations

[and other stuff]

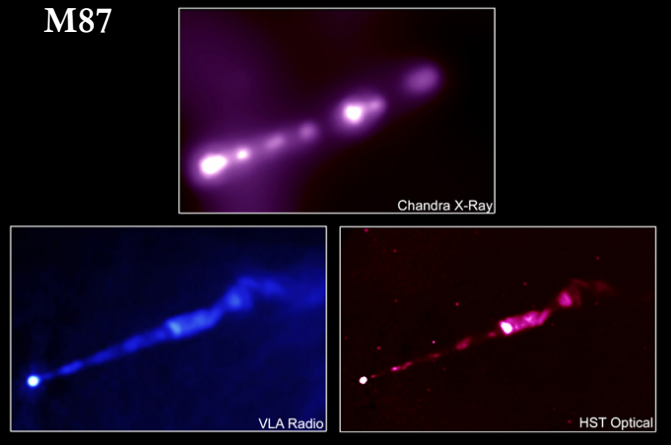
Jonathan McKinney
Stanford/KIPAC

Issues Addressed

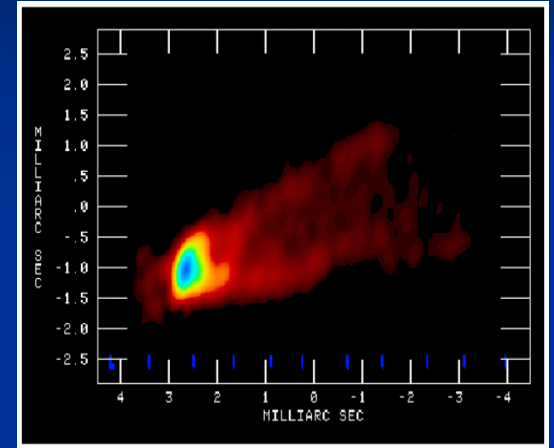
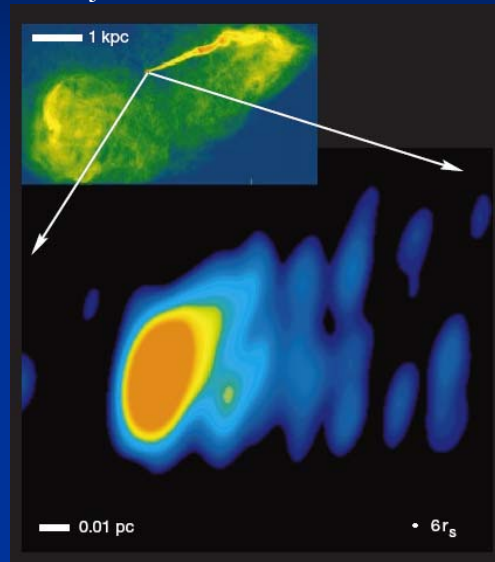
- Radio Loud/Quiet Dichotomy
 - Caused by Environment, Spin, Galaxy Evolution?
- Magnetosphere near BH
 - How different from NS?
- Jet Launching and Stability during Accretion
 - What is Dependence on Field and Turbulence?
 - What helps or hurts stability?

AGN Jets

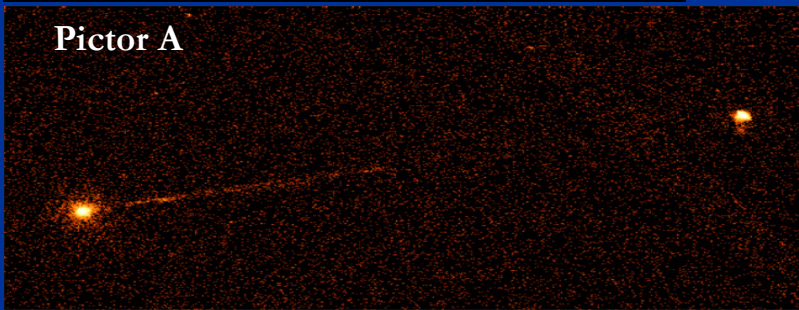
M87



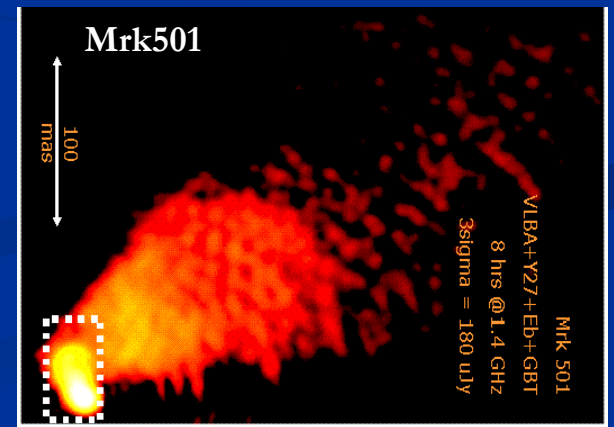
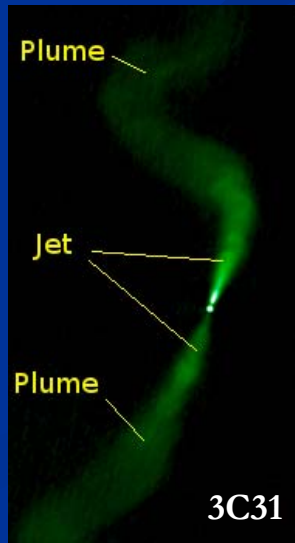
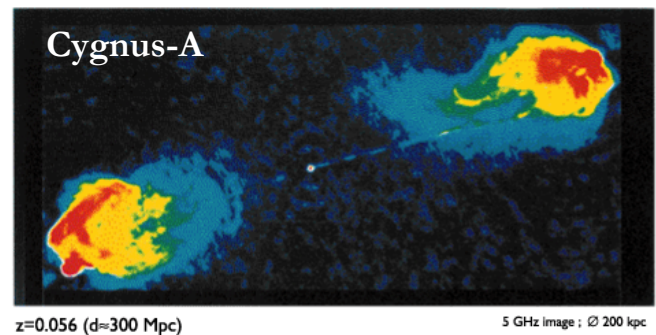
Junor/Biretta/Walker



Pictor A



Radio Image of Cygnus-A (FR-II)



The Sample

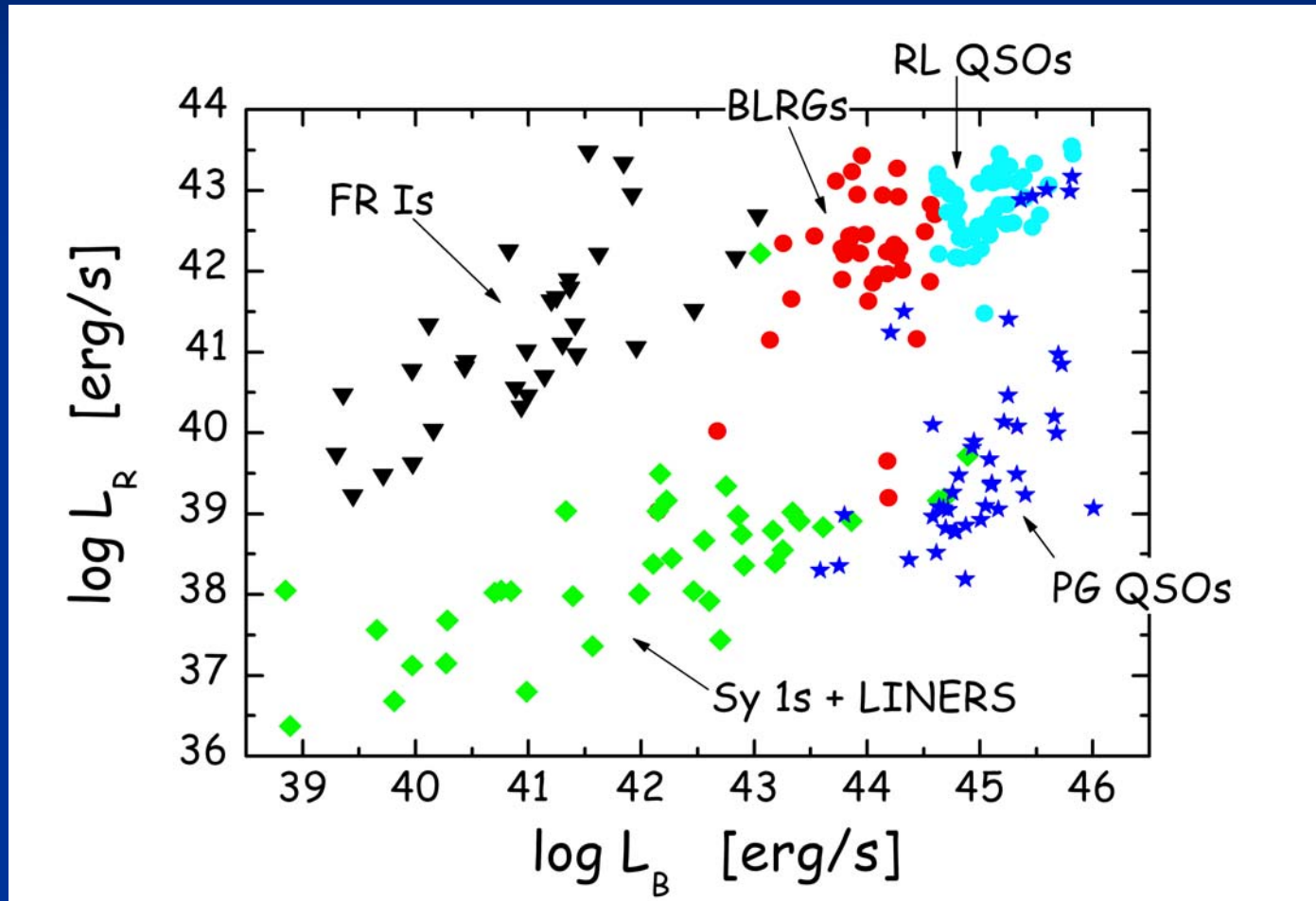
Want to compare the main jet parameter, i.e. bulk kinetic power L_j , with the main parameters of the central engine, namely M_{BH} and L_{acc} , for AGNs covering many decades in radio and disk luminosities. Hence, our sample has to be by definition *heterogeneous and incomplete*.

Select sources (**BLRGs, BLRQs, Sy1s, LINERs, FR Is, PG QSOs**) for which:

- (i) the optical flux of the unresolved nucleus is known;
- (ii) the total radio flux is known (including extended emission);
- (iii) the black hole mass can be estimated.

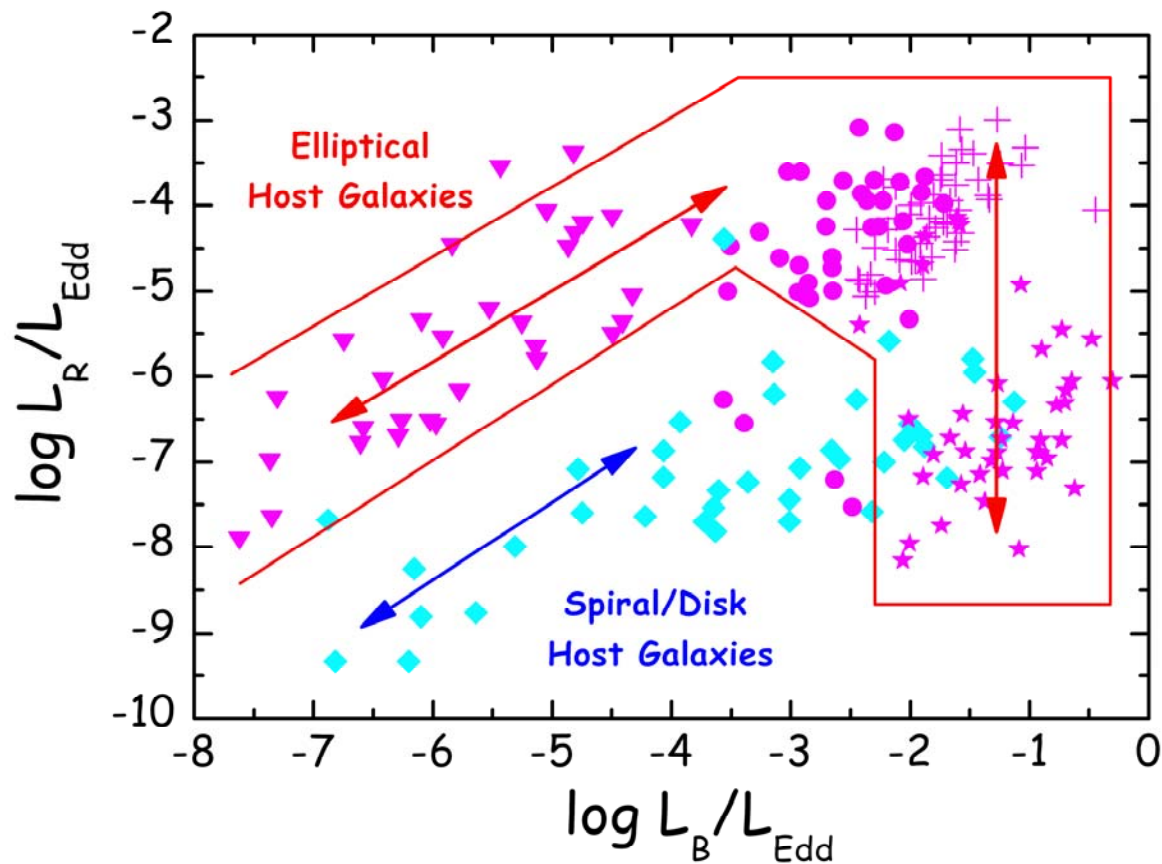
Want to avoid complications due to significant beaming and obscuration, and hence we exclude blazars (**OVVQs, HPQs, FSRQs, BL Lacs**) as well as type-2 AGNs (**NLRGs, Sy2s**).

Two sequences on L_B-L_R plane

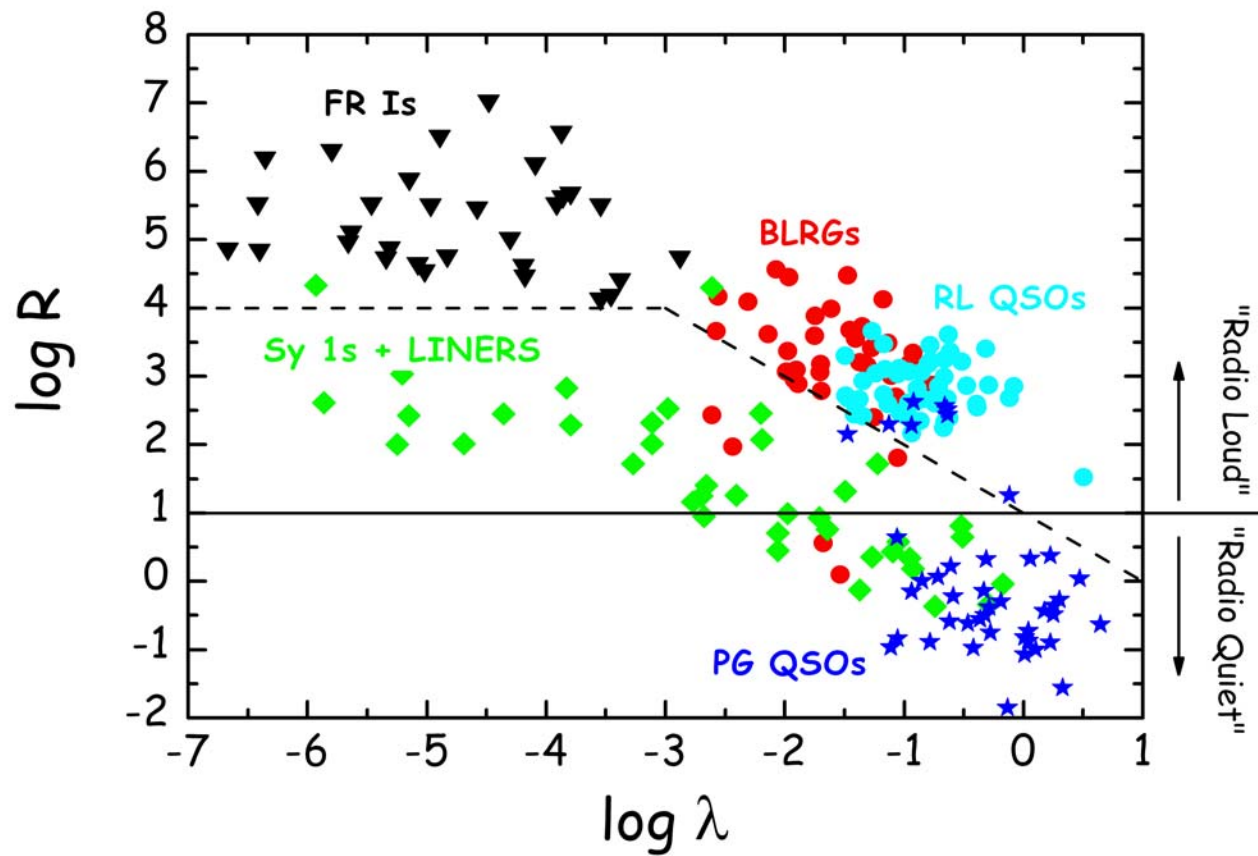


(Sikora, Stawarz, & Lasota 2007)

The same two sequences emerge on (L_B/L_{Edd}) - (L_R/L_{Edd}) plane



Radio-Loud / Radio-Quiet



Main parameters

$$L_B \equiv v_B \times L_{vB} \quad , \quad \lambda_B \equiv 4400 \text{ \AA}$$

nuclear B-band luminosity
by assumption $L_{\text{acc}} = 10 \times L_B$

$$L_R \equiv v_R \times L_{vR} \quad , \quad v_R \equiv 5 \text{ GHz}$$

total jet radio luminosity
by assumption $L_j \propto L_R$

$$R \equiv L_{vR}/L_{vB} \approx 10^5 \times (L_R/L_B)$$

radio-loudness parameter

$$L_{\text{Edd}} = 4\pi GM_{\text{BH}}m_p c/\sigma_T \approx 10^{38} \times (M_{\text{BH}}/M_{\odot}) \text{ erg/s}$$

Eddington luminosity

$$\lambda \equiv L_{\text{acc}}/L_{\text{Edd}} = 10 \times (L_B/L_{\text{Edd}})$$

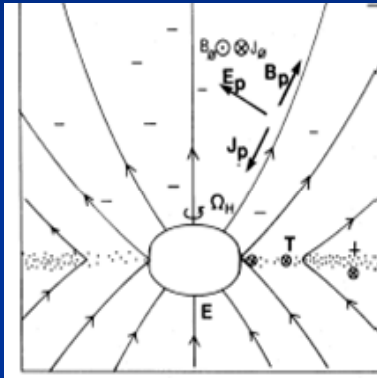
accretion rate

Possible Solutions to Dichotomy

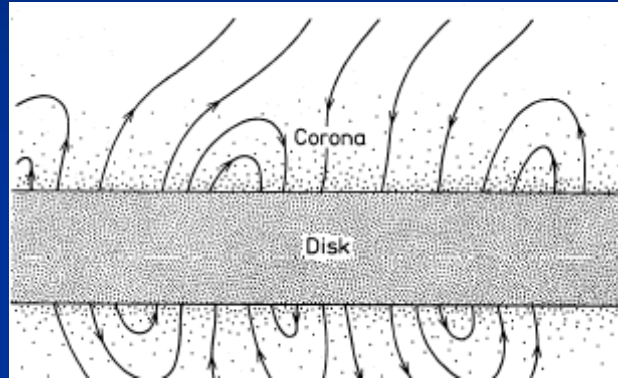
- Changes in Field Geometry or Confinement
- Mass-loading (Meier et al. 97: Magnetic Switch)
- Variation in amount of BH/Disk Magnetic Flux
 - Flux trapping (Reynolds 06, Garafalo 09)
 - Difference in Disk Thickness (Meier 01)
- Non-Dipolar Fields (Beckwith/McKinney)
- Radio power may not map to jet power:
 - ISM interaction and radio generation

Disk-BH-Jet Connection

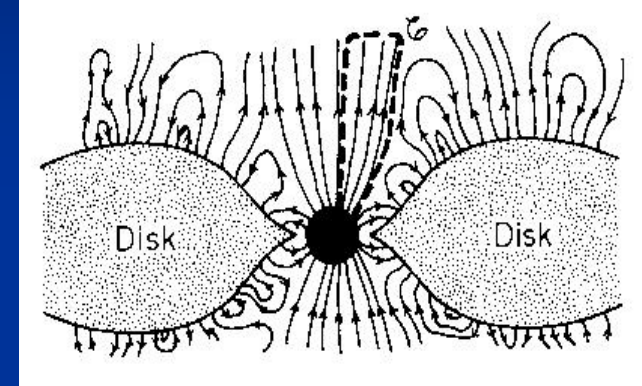
BZ77



Blandford & Payne '82

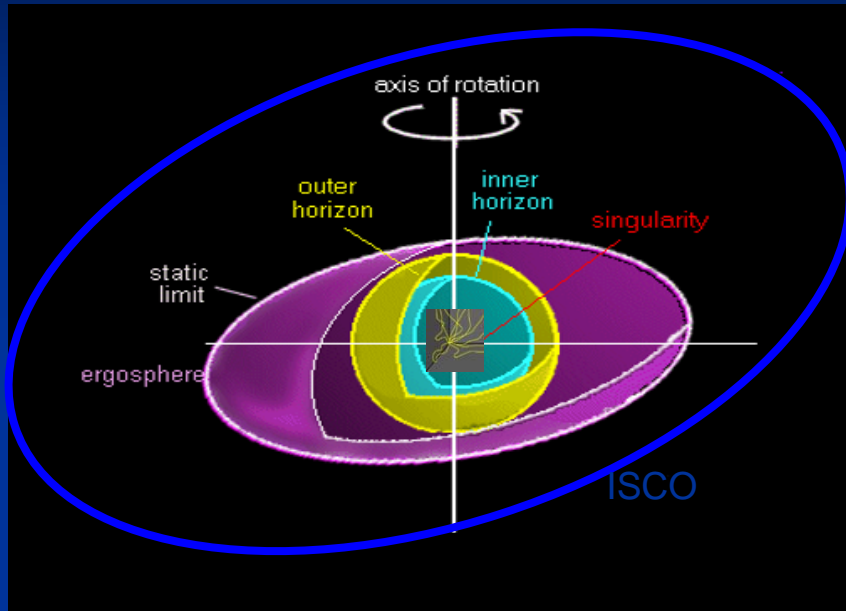


MacDonald & Thorne '82



- What is the structure of the electro-magnetosphere?
 - vs. spin? (Meissner Effect)
 - vs. jet mass-loading? (σ)
 - vs. accreted field geometry? (dipolar vs. quadrupolar, etc.)
 - in presence of non-axisymmetric turbulence?

Black Holes



Michell 1783 Escape velocity:

$$mv^2/2 = GMm/r \rightarrow v = \sqrt{2GM/r}$$

$$v = c \rightarrow r = 2GM/c^2$$

No-Hair Theorem:

Mass: M, Spin: J, Charge: ~~Q~~

Horizon: $r_H = M \pm \sqrt{M^2 - a^2}$

Static Limit: $r_S = M \pm \sqrt{M^2 - a^2 \cos^2 \theta}$

Inner-most stable circular orbit (ISCO): $3r_H$ for $a=0$, $1r_H$ for $a=M$

Photon Sphere: Inside, objects cannot orbit at all ($\sim 3/2r_H$ for $a=0$)

Static Limit: Varies from $1r_H$ to $2r_H$ for $a=M$ (ergosphere inside)

Horizon: Inside, objects must fall

Singularity: Reached in finite time (superstring theory?)

Equations of Motion

- Mass Conservation (ρ_0)
- Maxwell's Equations (E, B, ρ, J)
- Energy-Momentum Conservation ($u, p, \text{ and } v$)

Closure:

- EOS: $p(u)$
- Ohm's law: $J(E, B, \eta_i) : \text{e.g. } J_{co} = \sigma E_{co} \ \& \ E_{co} = 0$
- Viscosity + Resistivity

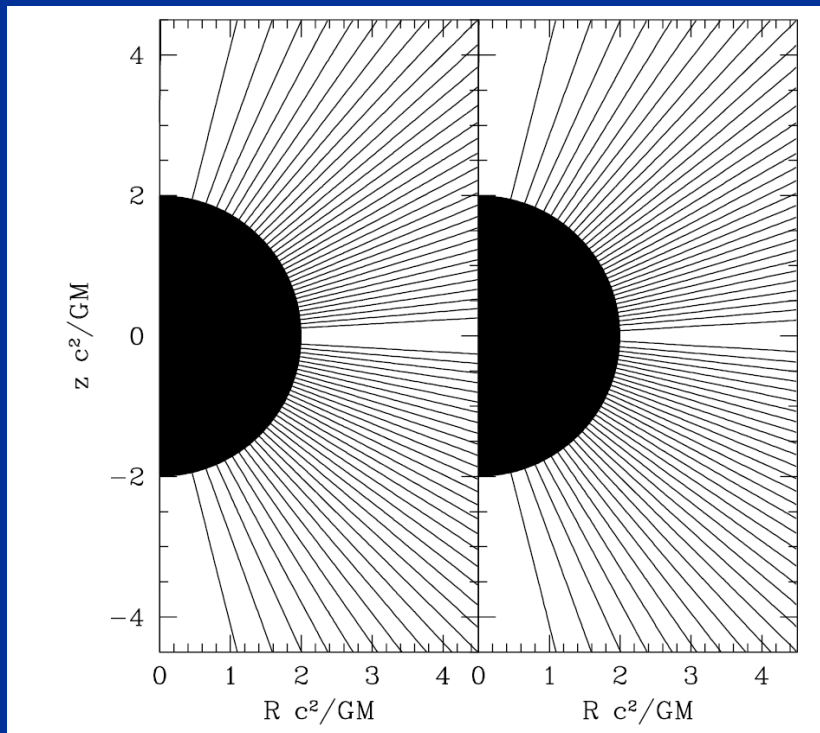
... in GR

BH Magnetospheres 101

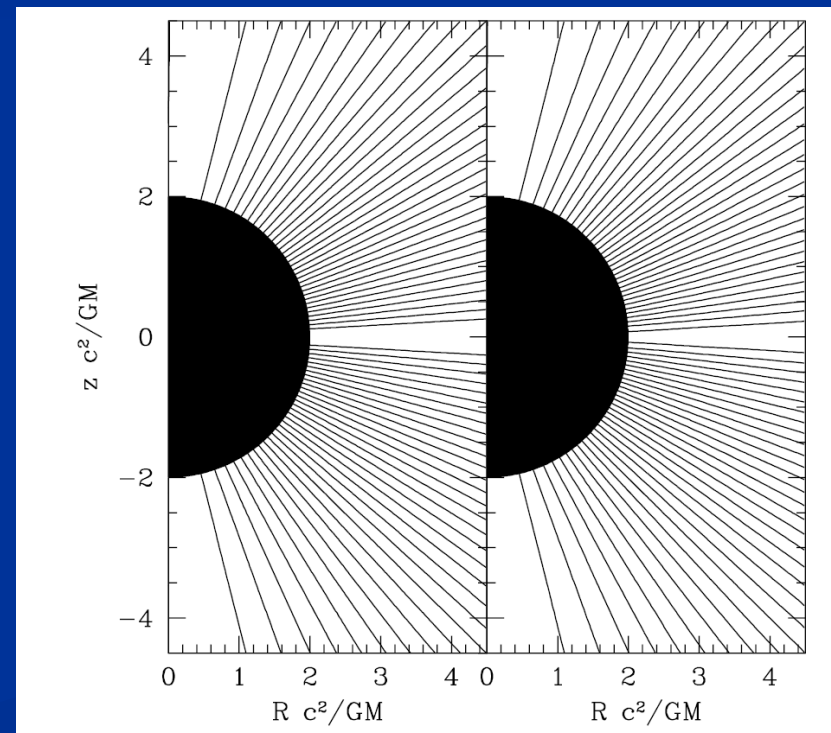
$$F_{\mu\nu} = A_{\nu,\mu} - A_{\mu,\nu}$$

$$A_\phi - A_\phi^0 = \int_s B \cdot ds = \int_\theta \sqrt{-g} B^r - \int_r \sqrt{-g} B^\theta$$

$$A_\phi = (1 - \cos\theta)$$

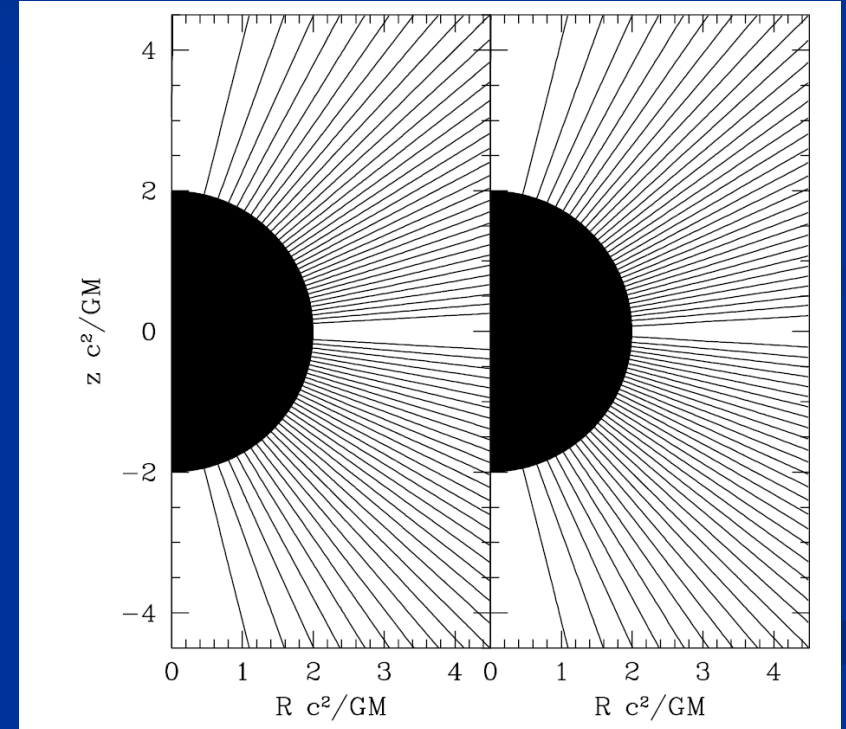
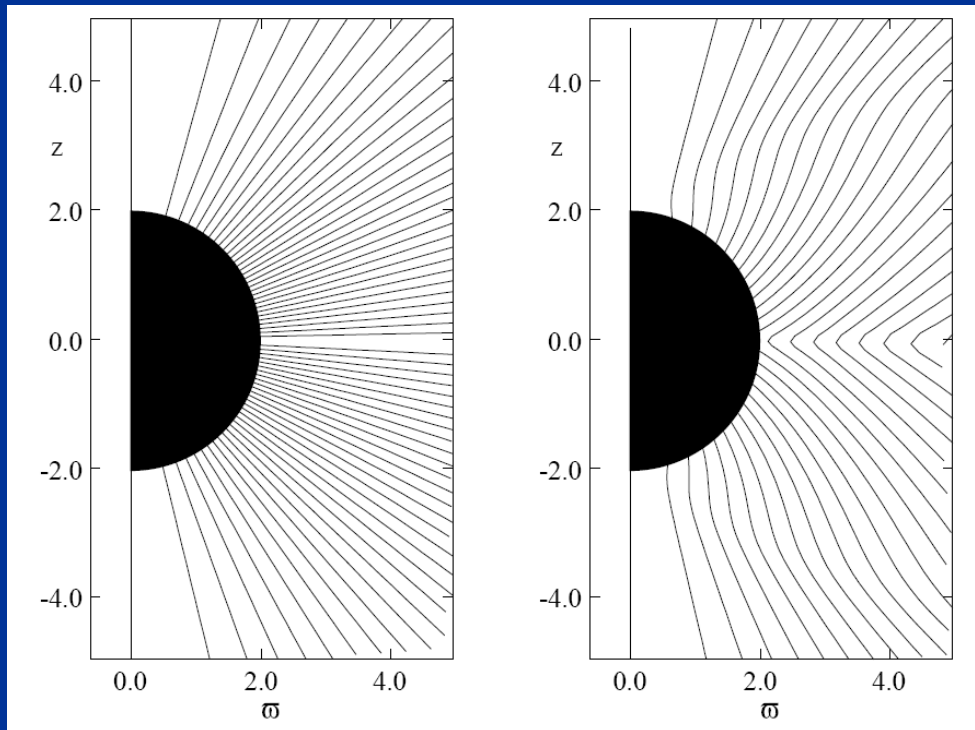


Pure Monopole $a=0$



Split-Monopole $a=0$

Split-Monopole $a=0$ or small a



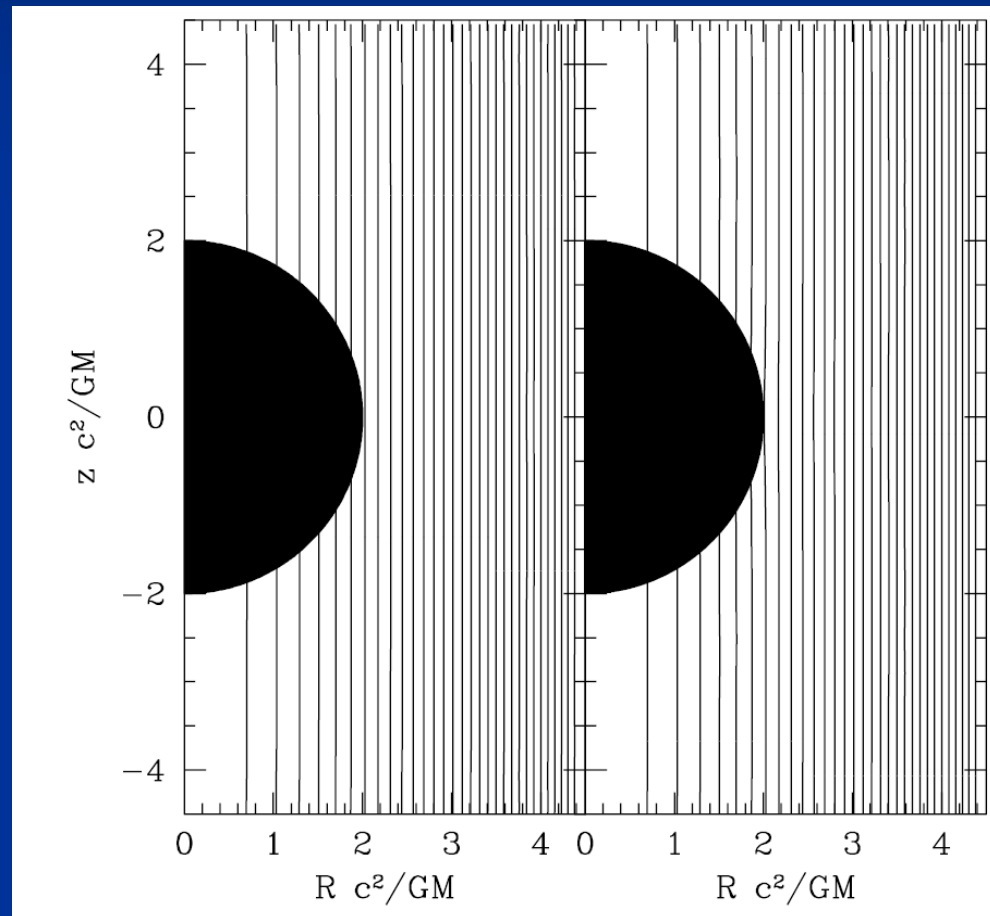
Komissarov 04 – $5GM/c^3$

McKinney 06 – $10^4GM/c^3$

Wald Solution: $a=0$ or small a

$$A_\mu = B_0(2at_\mu + \phi_\mu)$$

$$J^\mu = F_{;\nu}^{\mu\nu} = 0$$



McKinney 06 – $10^4 \text{GM}/\text{c}^3$

Wald $a=0.9$

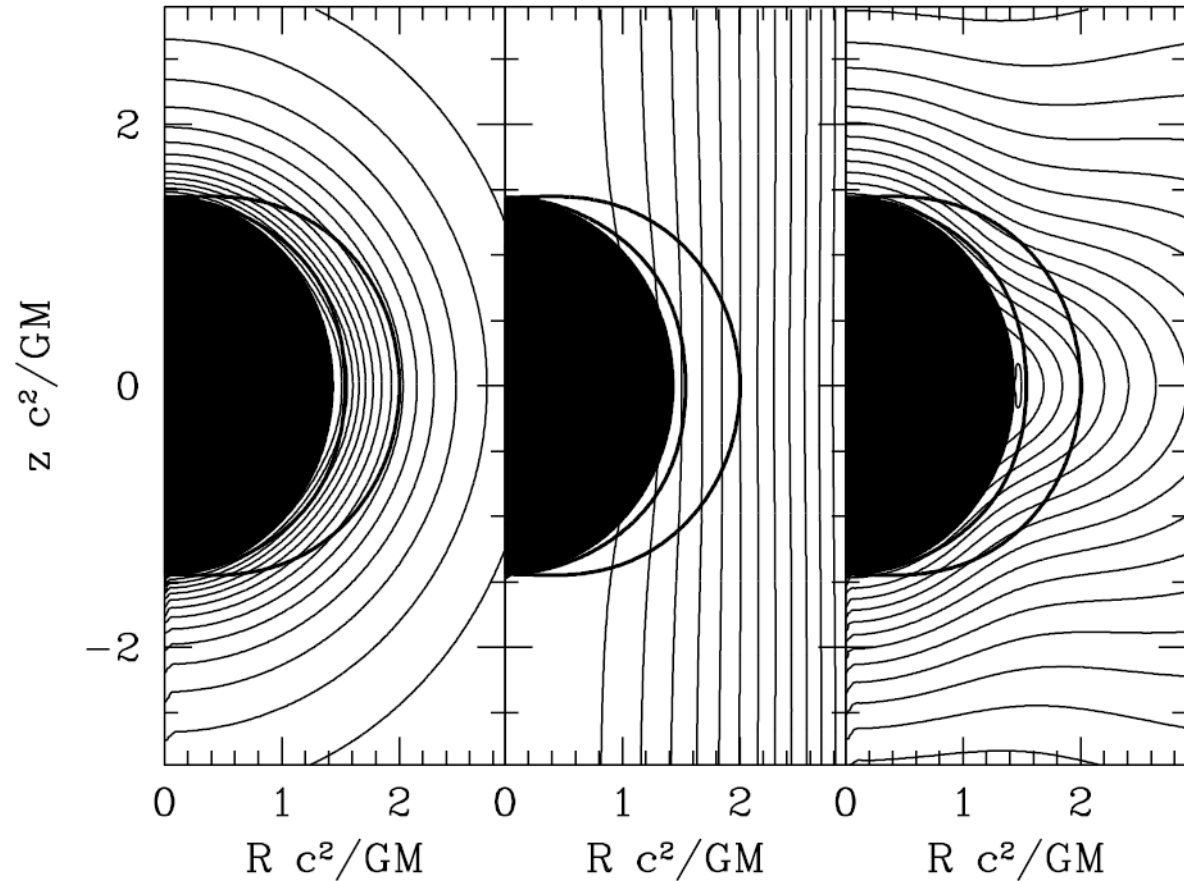
Light Surface:

$$u \cdot u = -1$$

$$u^t \rightarrow \infty$$

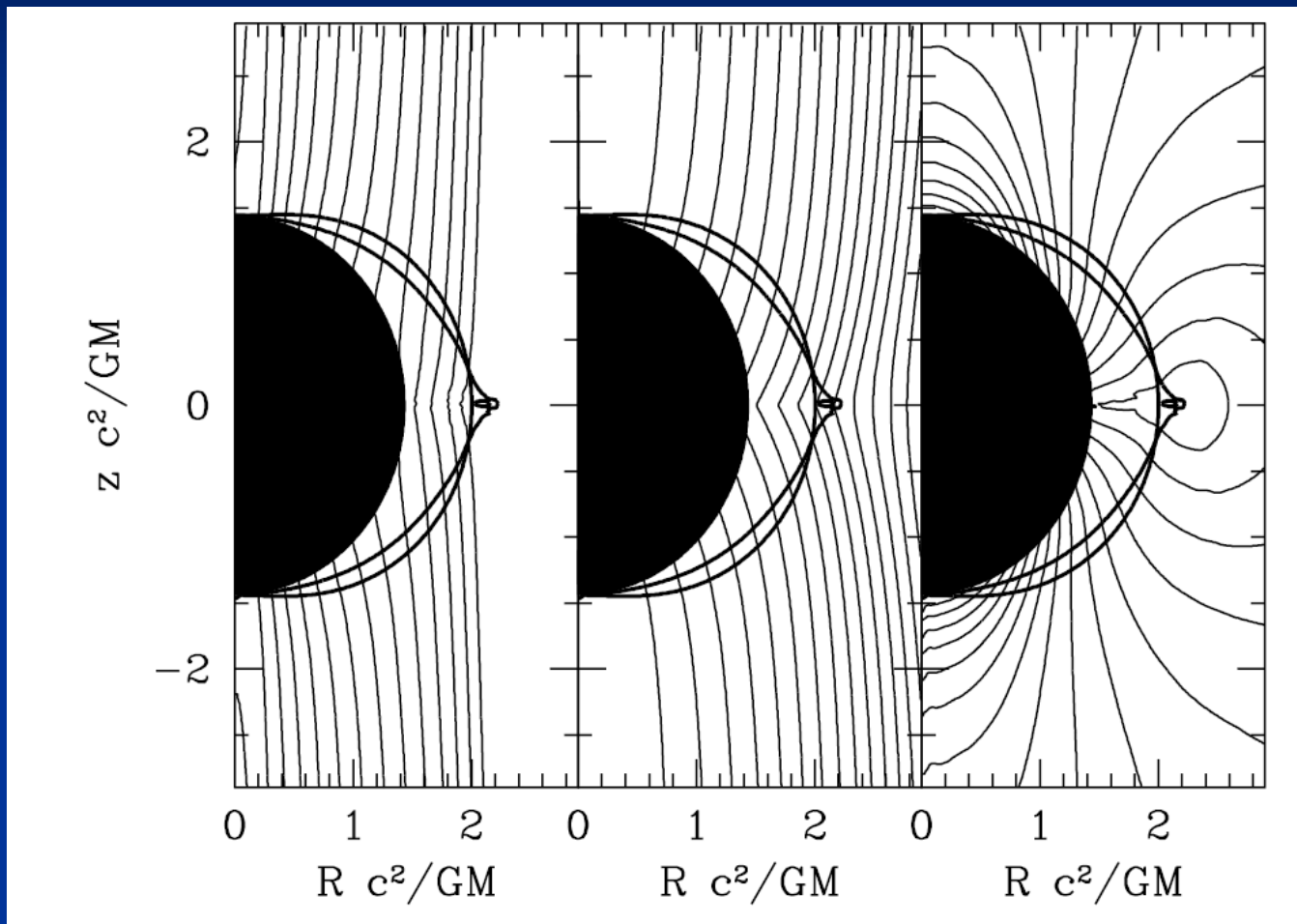
$$g_{tt} + 2\Omega g_{\phi t} + \Omega^2 g_{\phi\phi} = 0$$

$$\Omega = E_\theta/B^r = -E_r/B^\theta$$



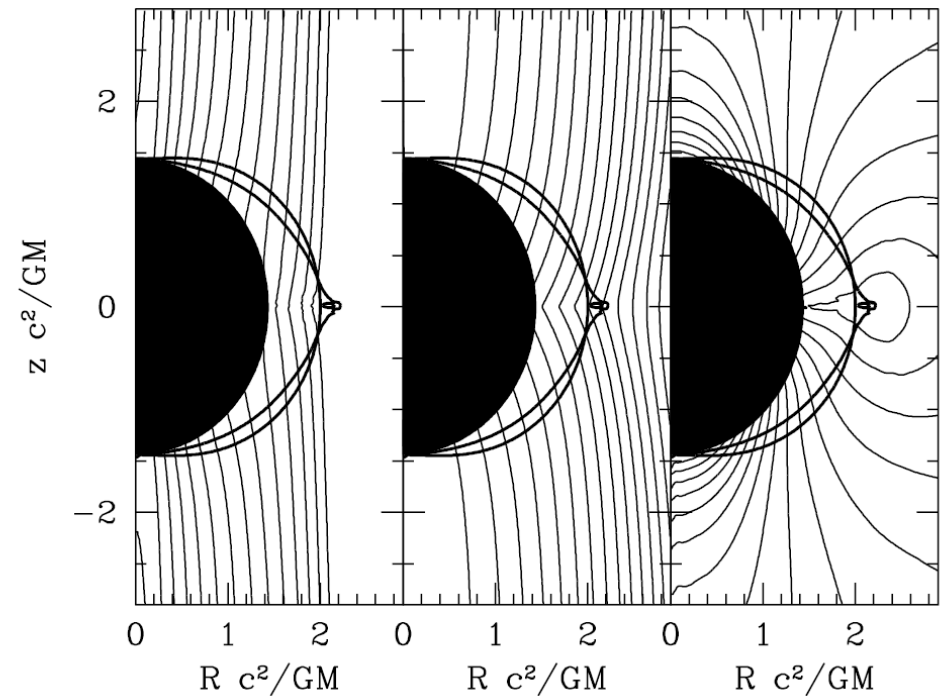
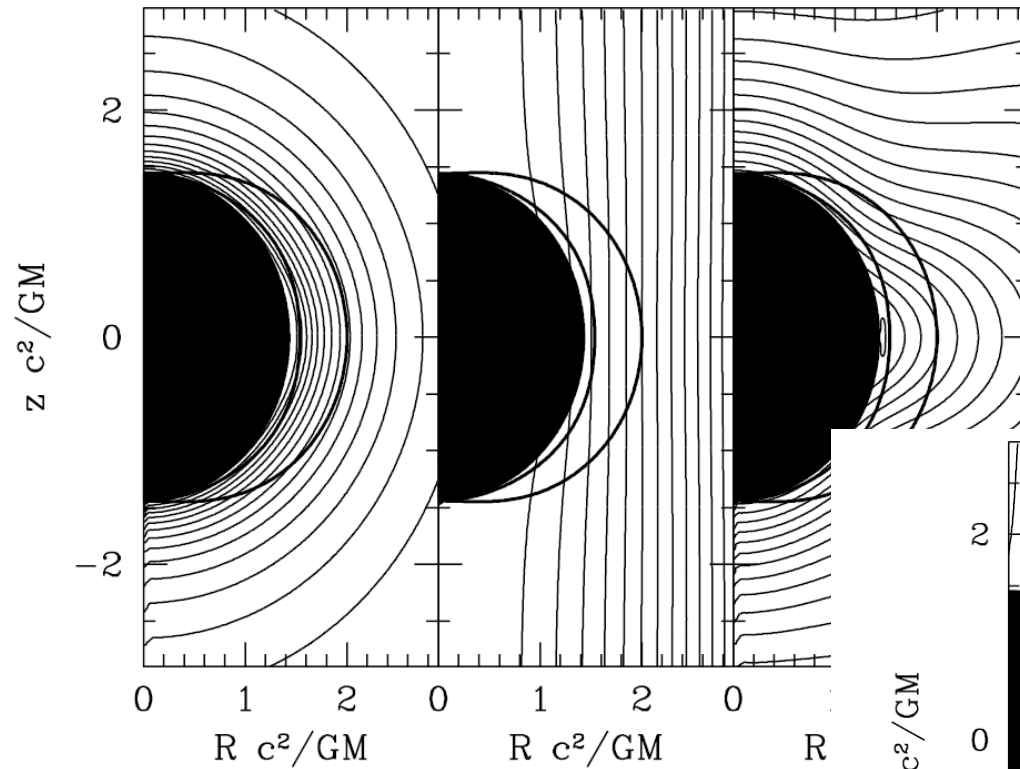
$$\Omega/\Omega_H: 0-0.67, A_\phi, b^2$$

Conductive* Wald $a=0.9$

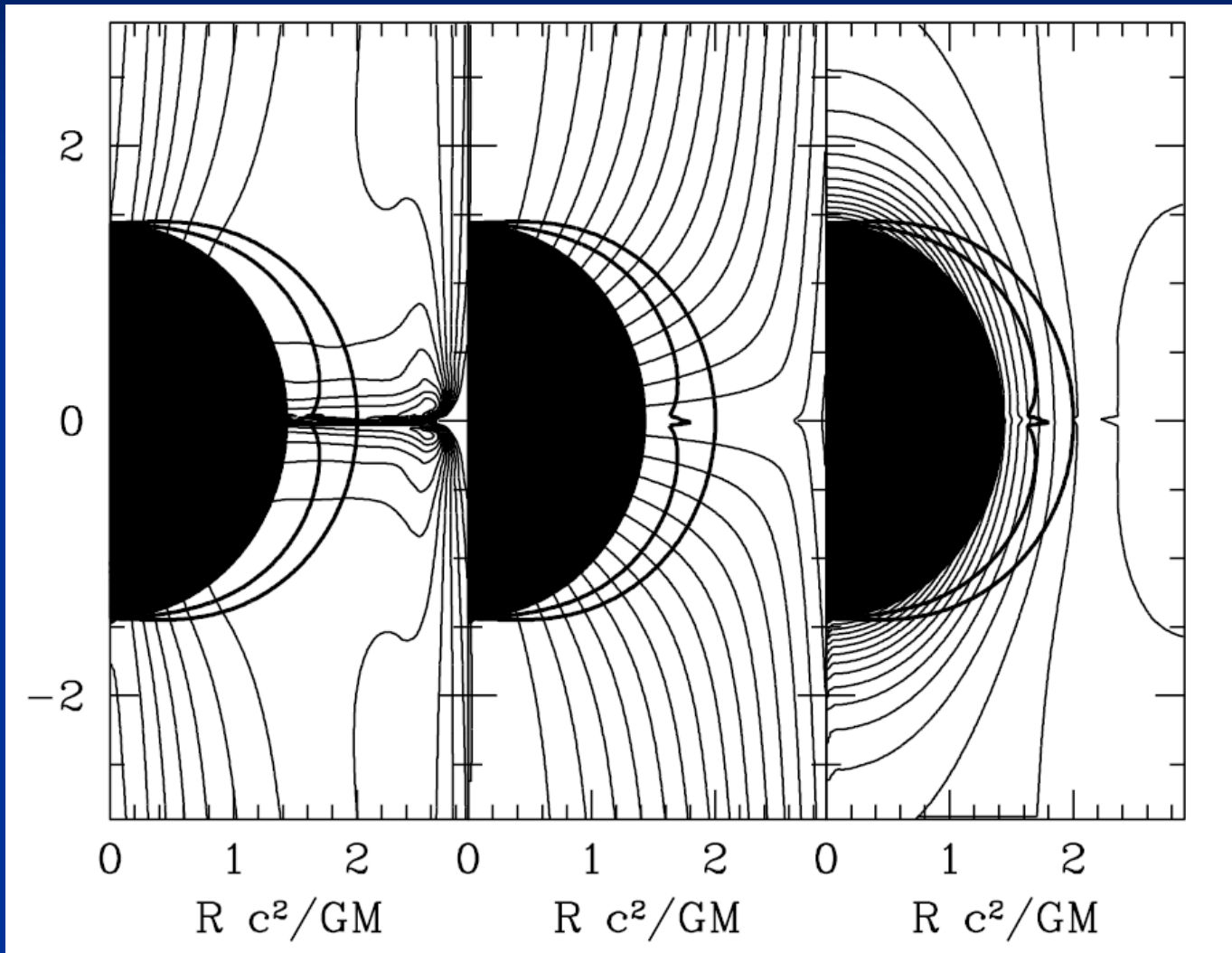


$$\Omega_F: 0-0.67, A_\phi, b^2$$

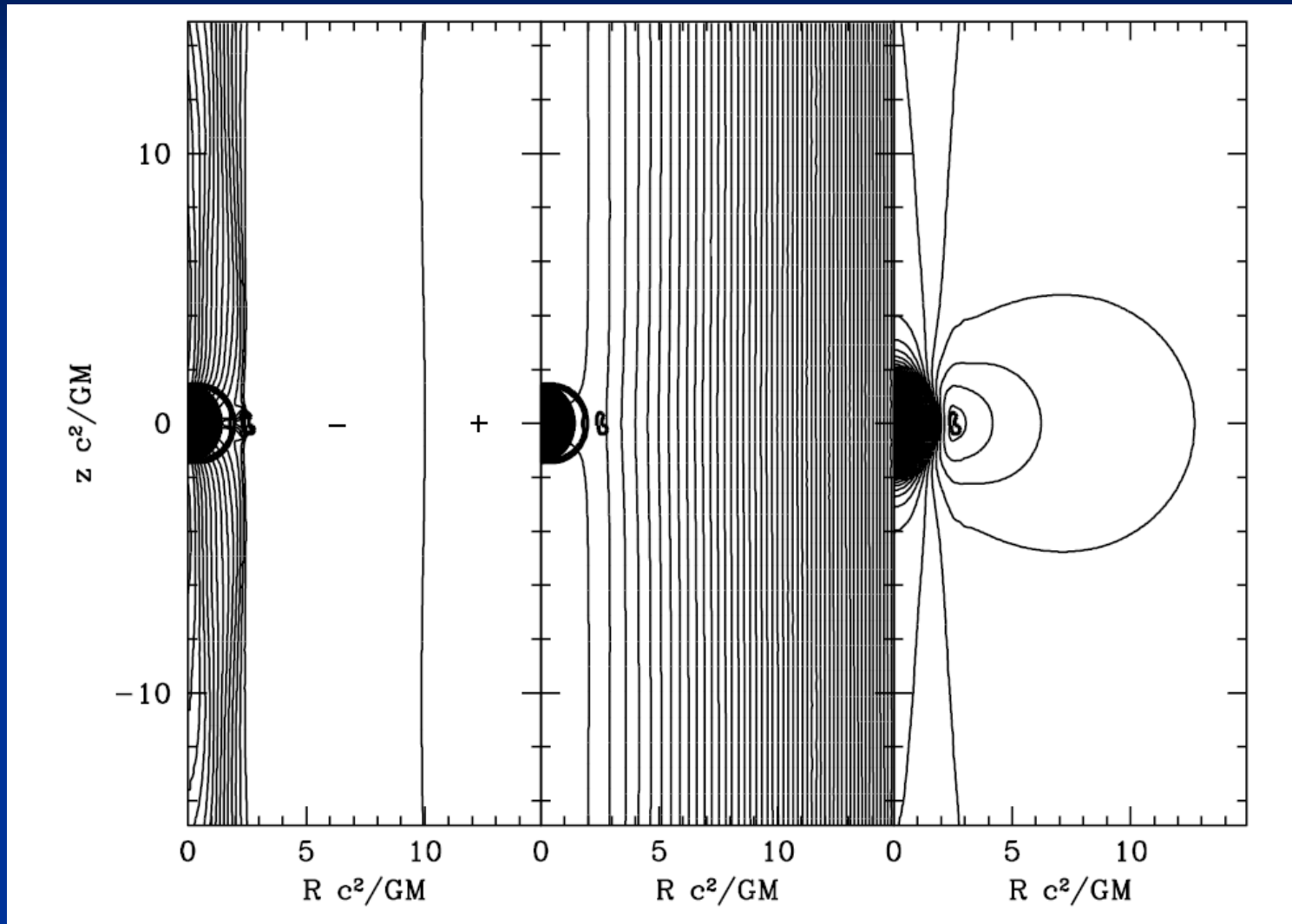
Vacuum vs. Conductive* Wald



Conductive** Wald $a=0.9$

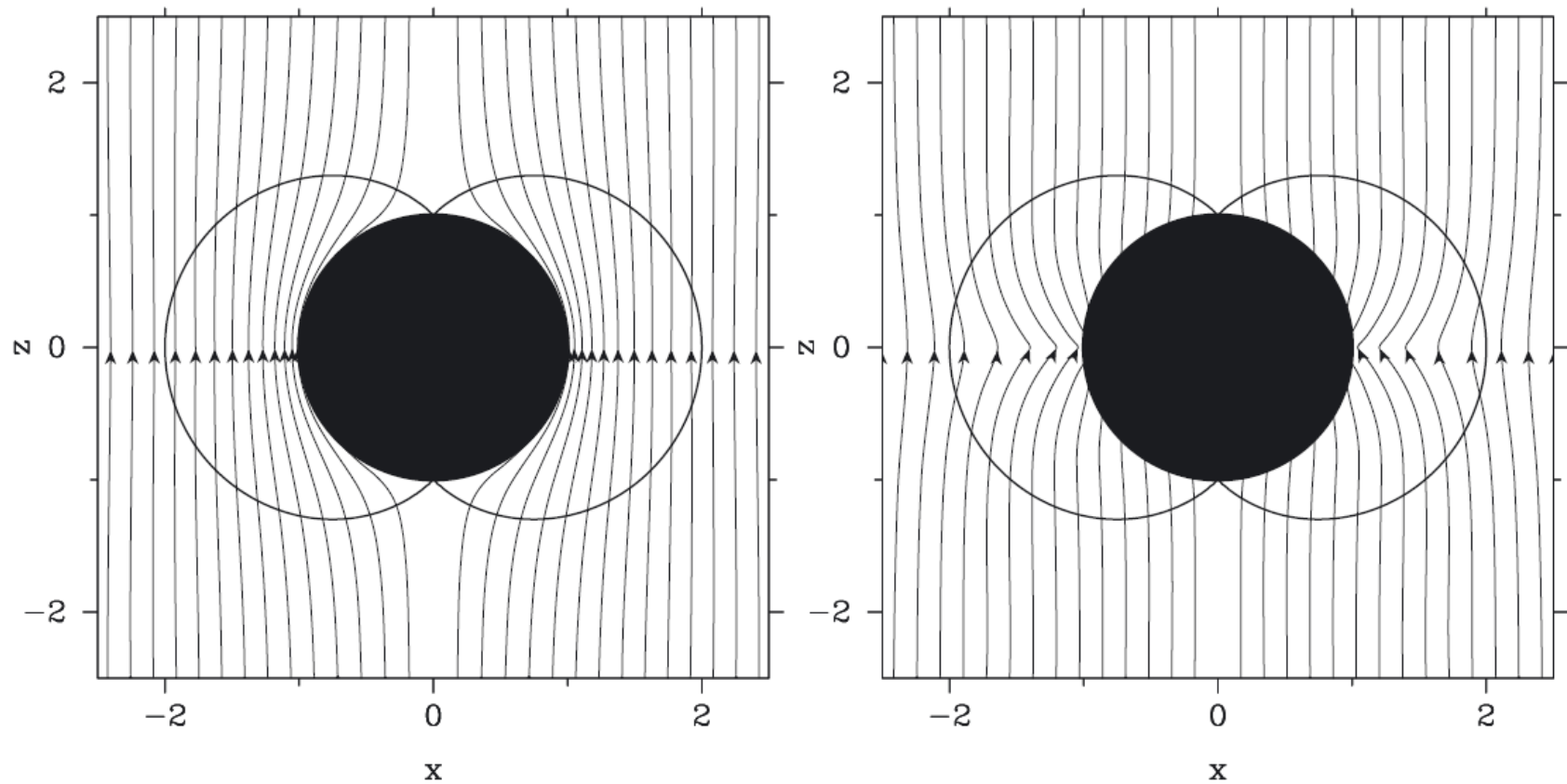


Conductive** Wald $a=0.9$



$a=0.999$ Force-Free GRMHD

Komissarov & McKinney 07 vs. King, Lasota, Kundt 75, et al.

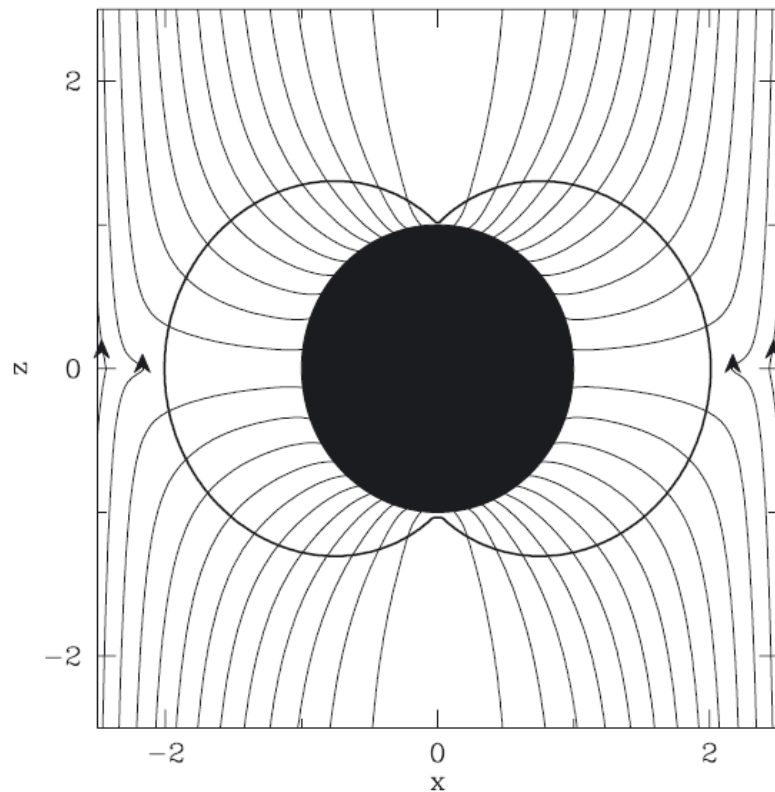


Vacuum

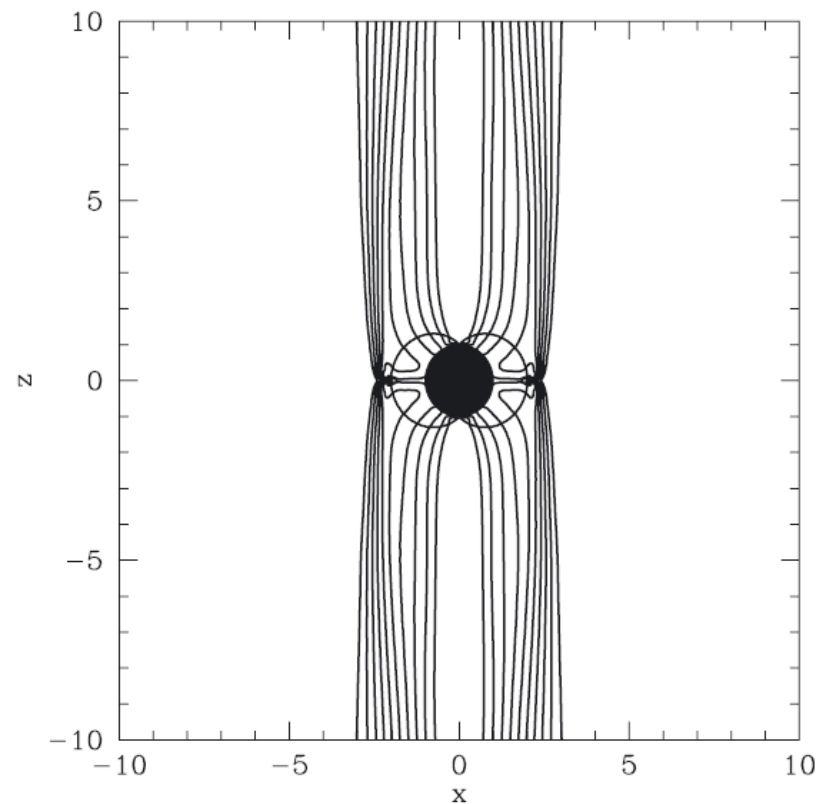
Conducting*

$a=0.999$ Force-Free GRMHD

Komissarov & McKinney 07

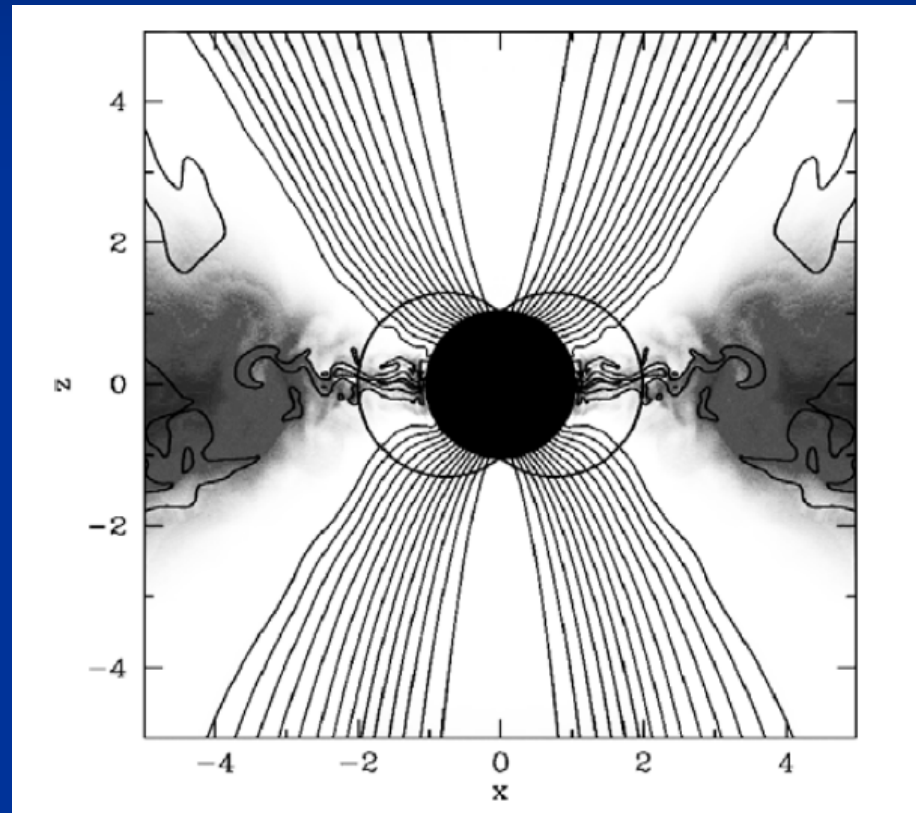
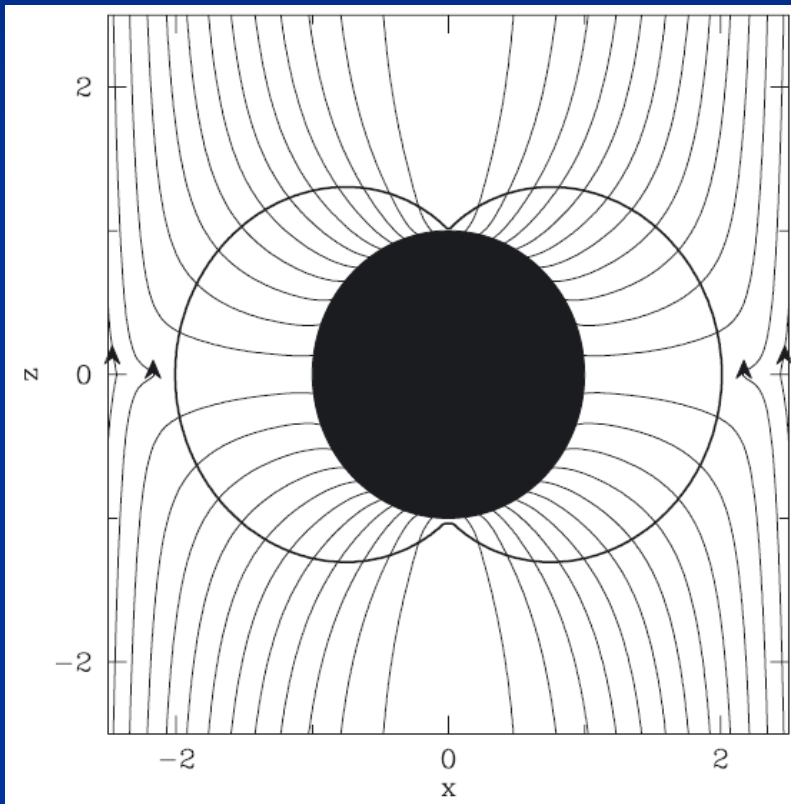


Conducting**

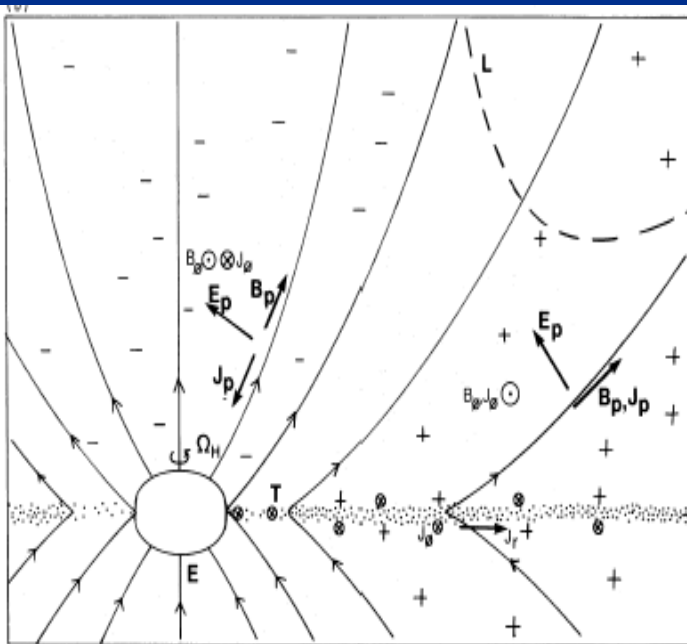


Ω_F Jet

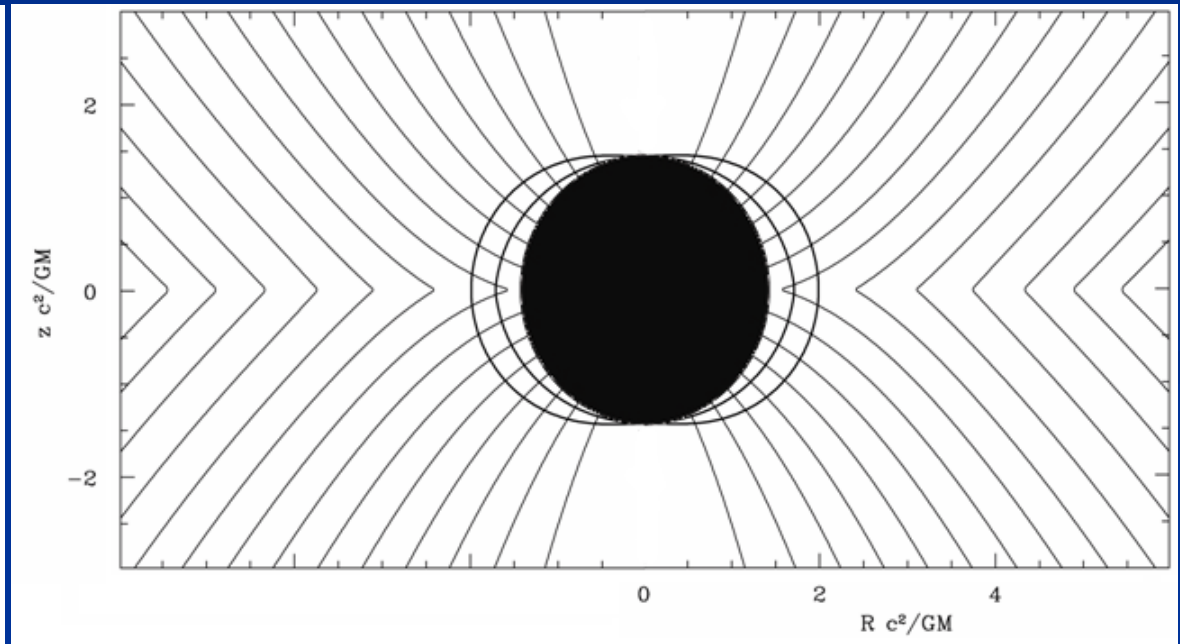
2D GRMHD Disk $a=0.999$



Disks w/ Paraboloidal Fields



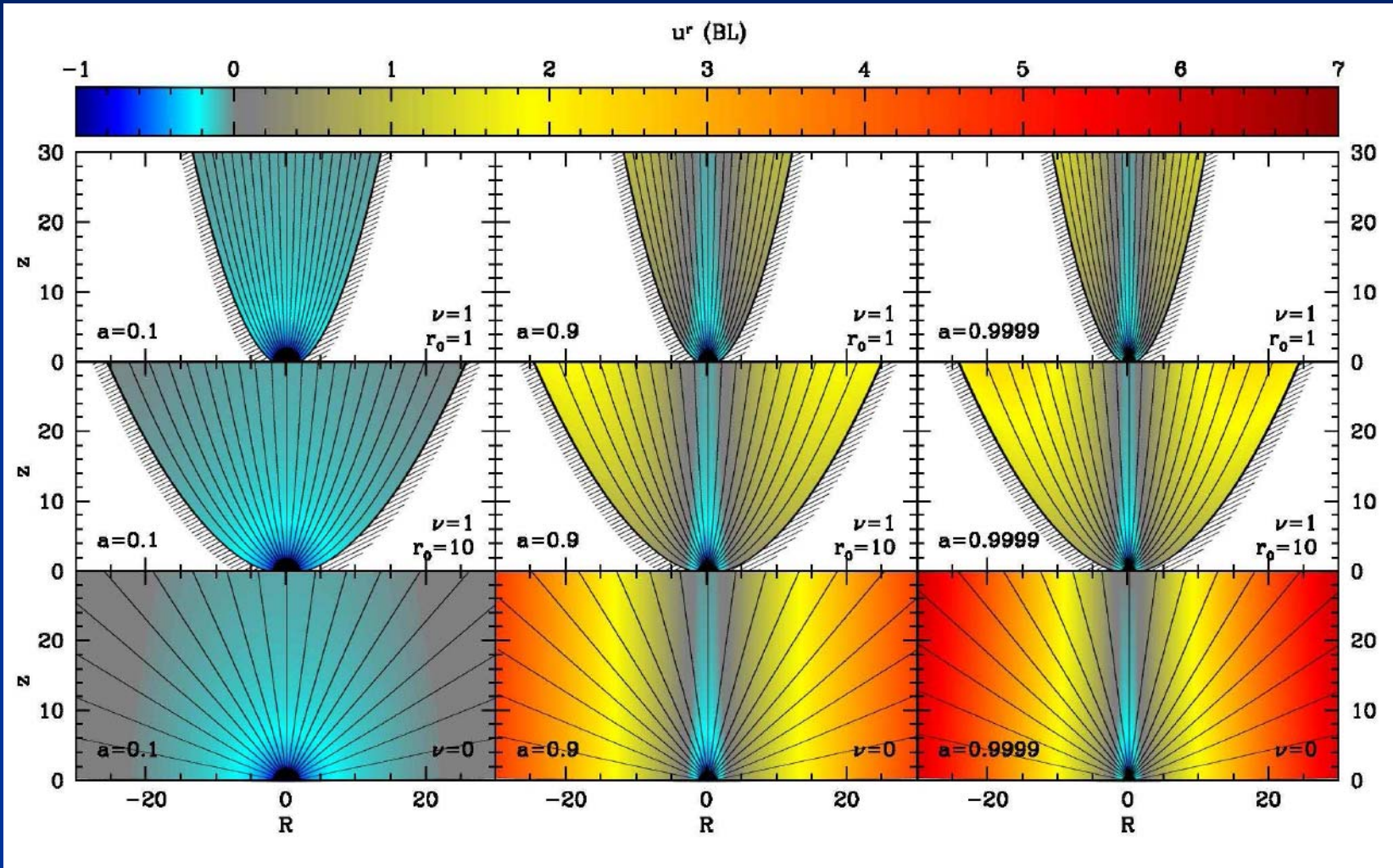
BZ77



McKinney & Narayan 2007

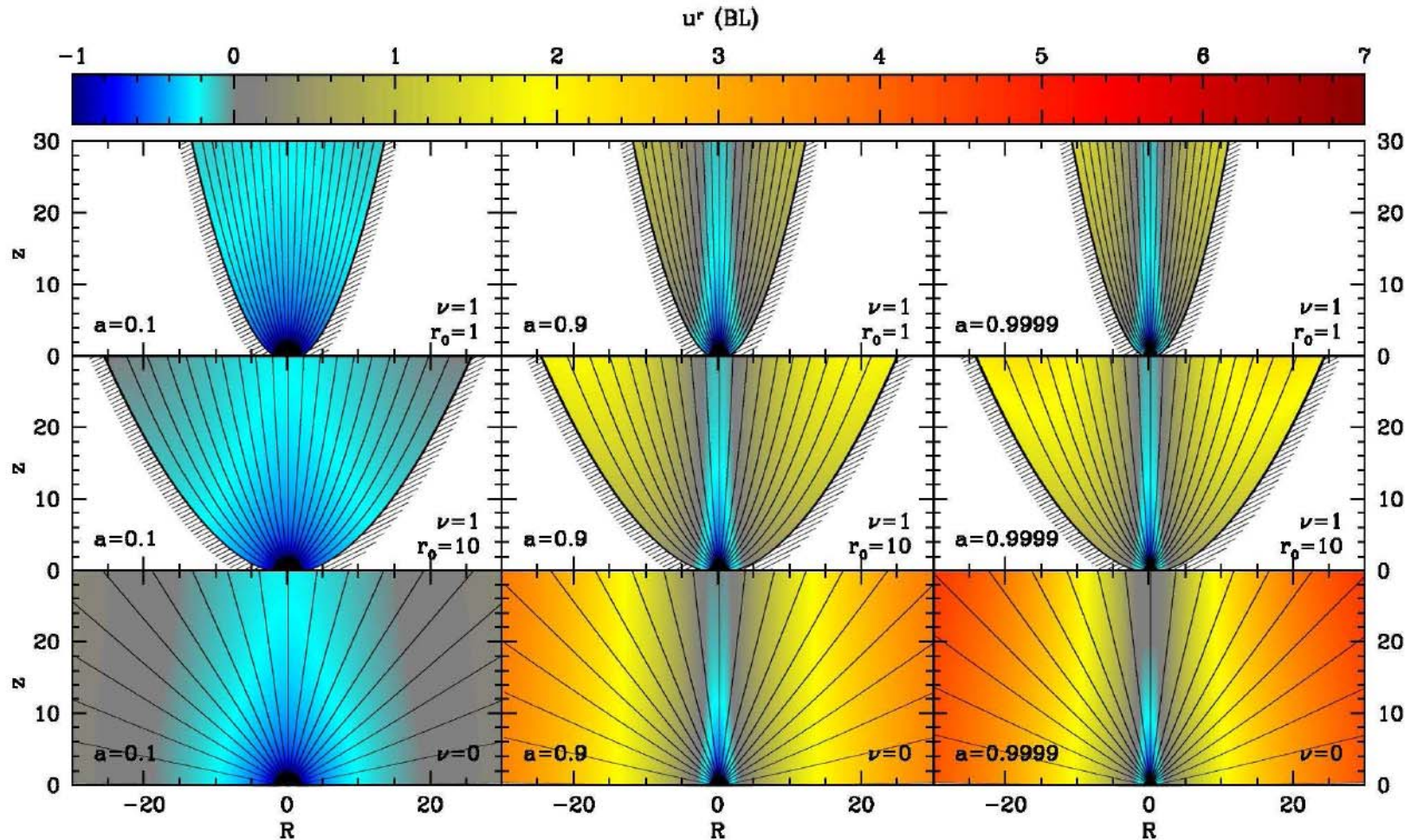
$$A_\phi \propto [(r + r_0)/(r_H + r_0)]^\nu (1 - \cos \theta)$$

Force-Free GRMHD



$$A_\phi \propto [(r + r_0)/(r_H + r_0)]^\nu (1 - \cos \theta)$$

GRMHD - walls

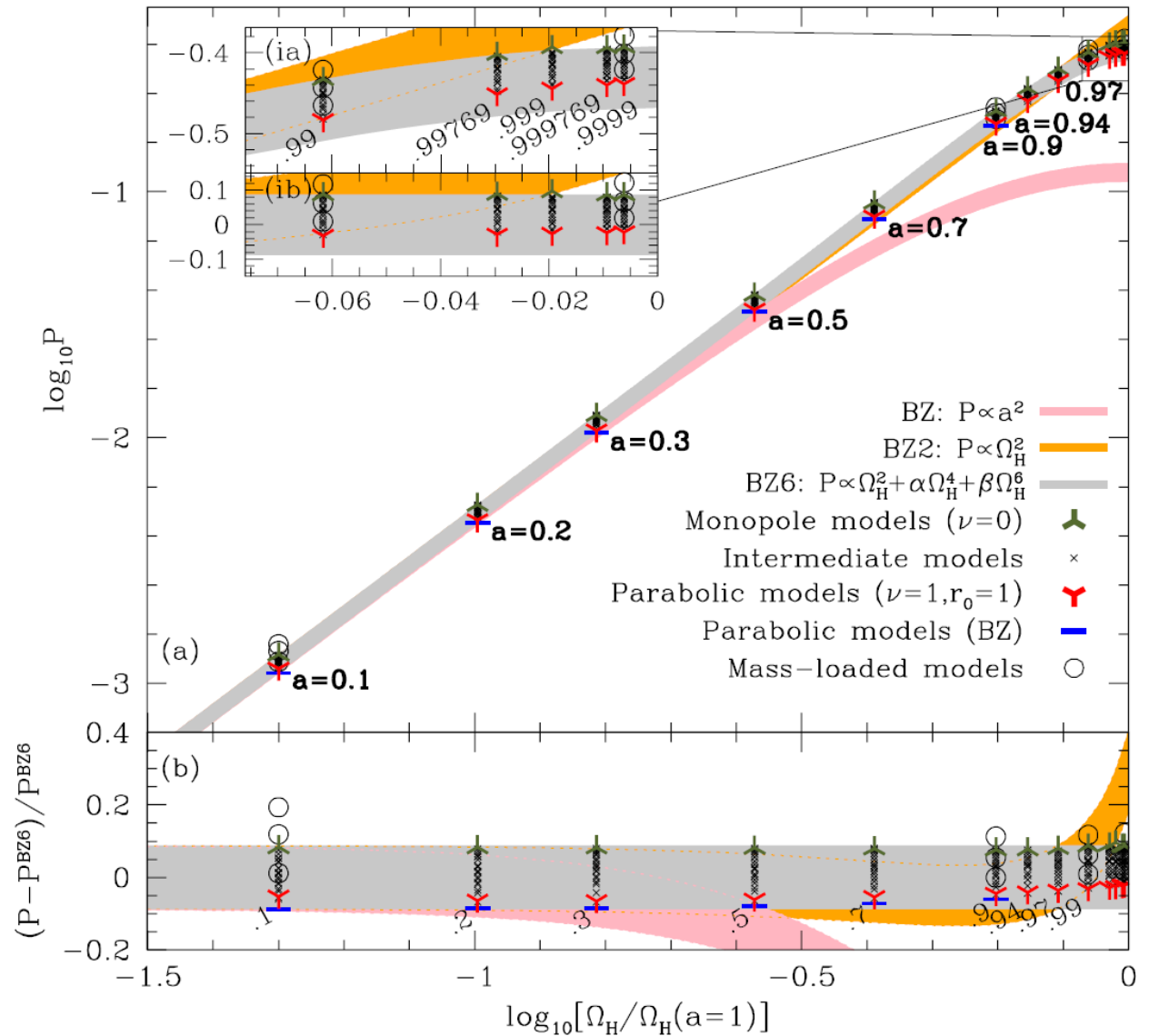


$$A_\phi \propto [(r + r_0)/(r_H + r_0)]^\nu (1 - \cos \theta)$$

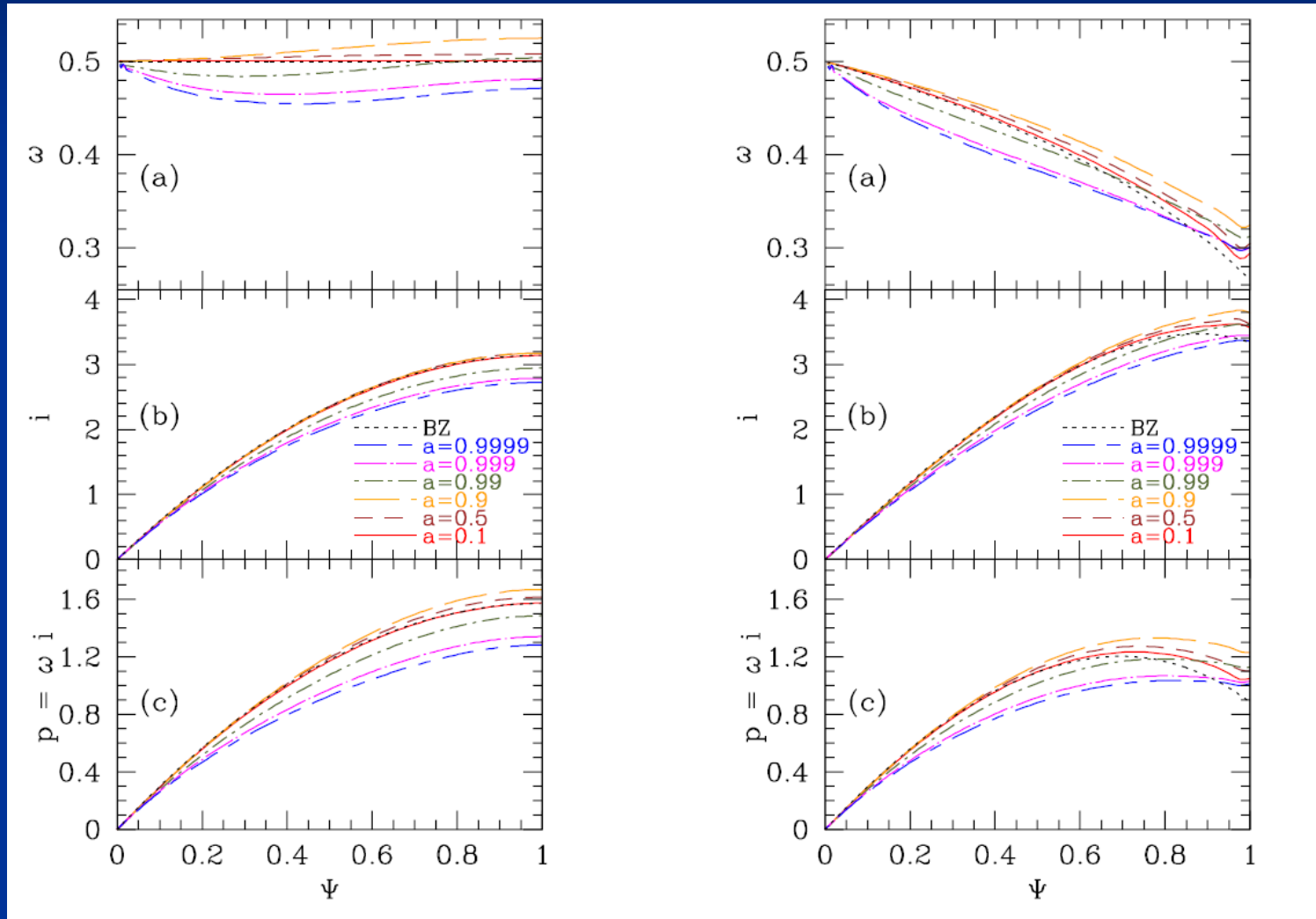
Power vs. Spin

BH Angular
Rotation Rate

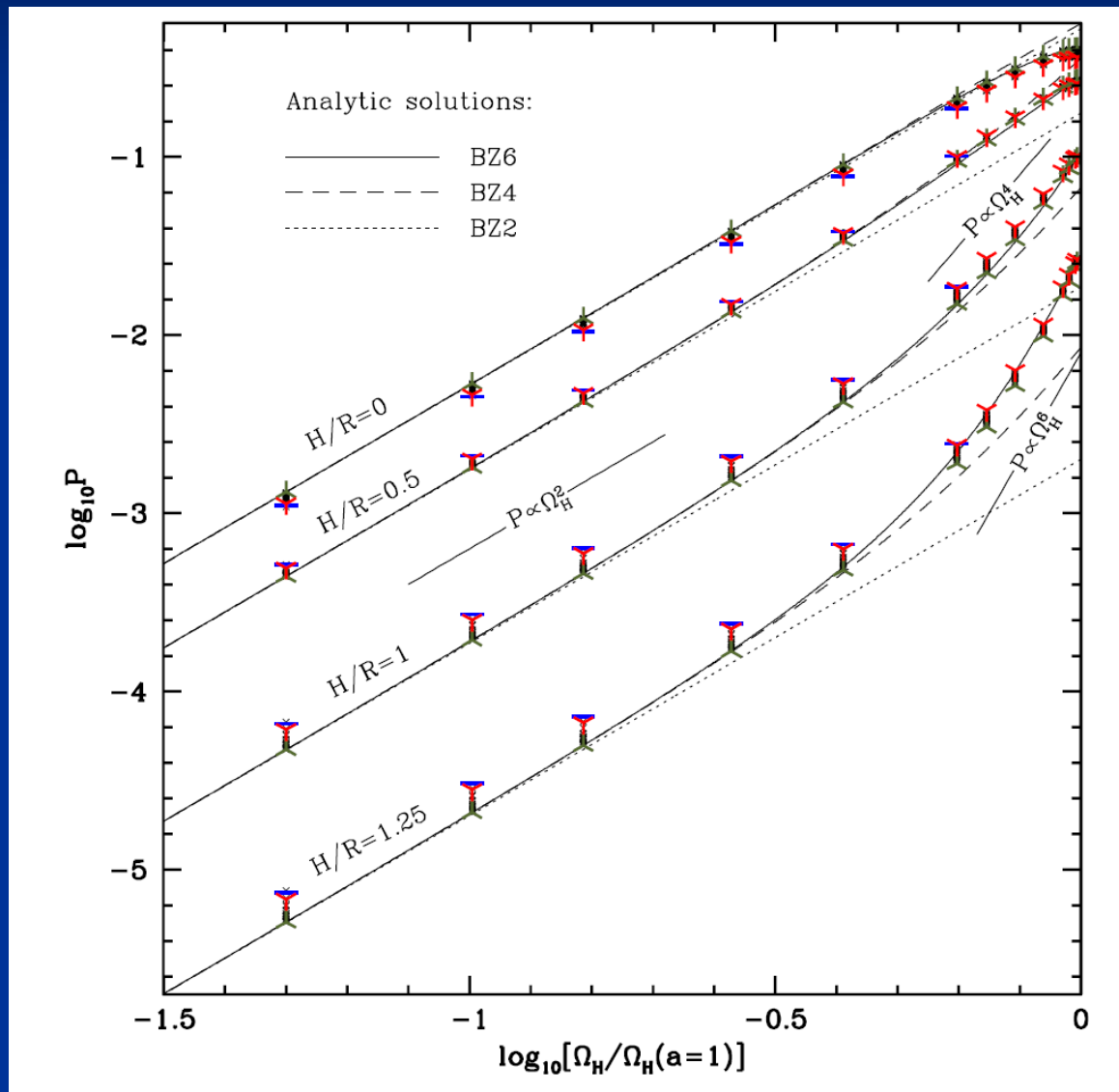
$$\Omega_H = ac / (2Mr_H)$$



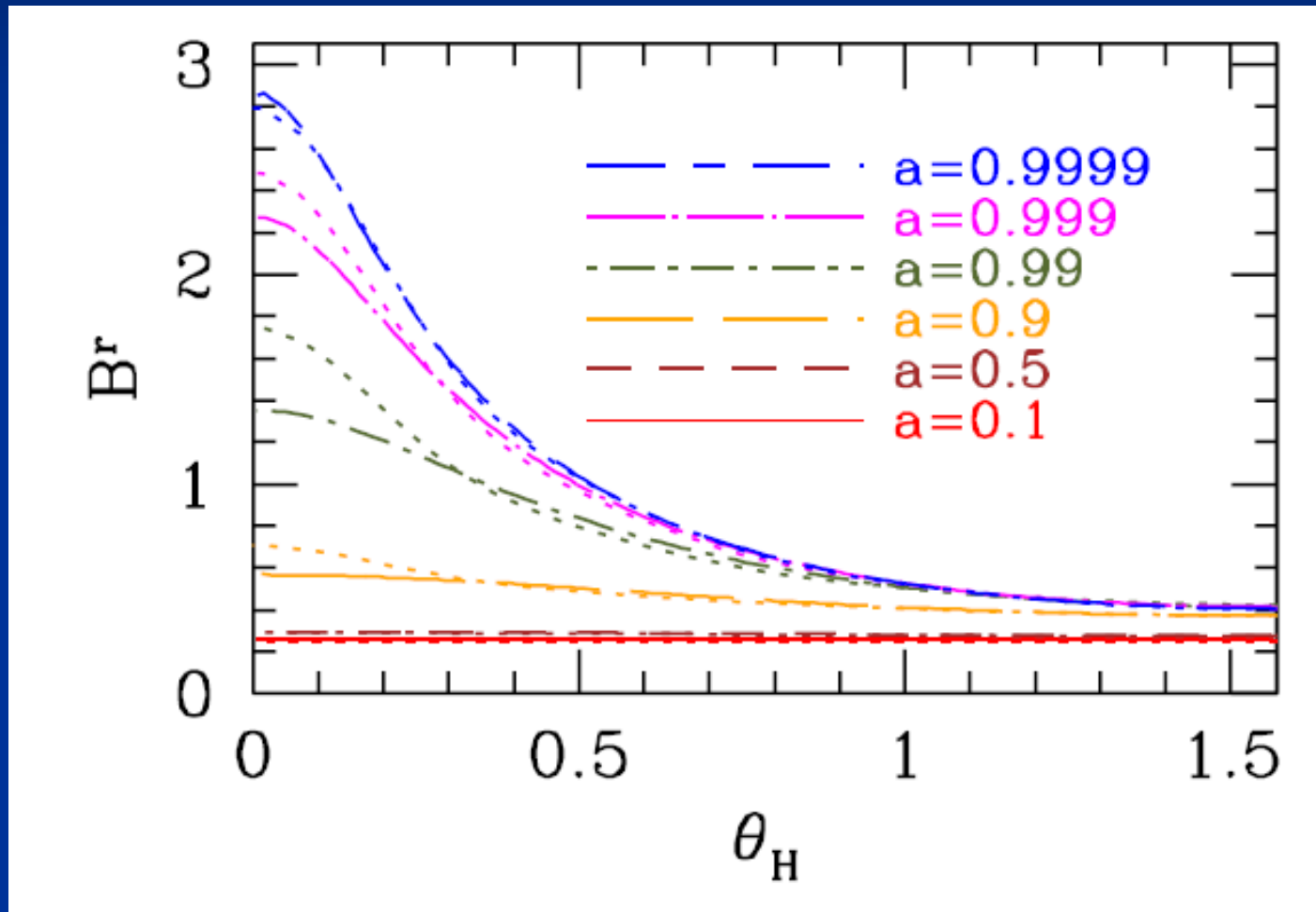
Conserved Quantities



Power vs. Angle vs. Spin



Field Strength vs. Angle vs. Spin



Derive Power vs. Spin vs. Angle

(1) $F_E(\theta) = [2(B^r)^2 \Omega (\Omega_H - \Omega) r M \sin^2 \theta] \Big|_{r=r_H} \sqrt{-g} B^r = A_\phi, \theta$

(2) Use $\Omega \approx \Omega_H/2 \rightarrow F_E$ to a^2

(3) Derive Monopole A_ϕ to $a^2 \rightarrow F_E$ to a^4

$$\Psi = \Psi_0 + 16\Omega_H^2 \Psi_2 + \Omega_H^4 \Psi_4$$

$$\Psi_0(\theta) = 1 - \cos \theta$$

$$\Psi_2(r, \theta) = f(r) \sin^2 \theta \cos \theta$$

(4) Numerically Motivate A_ϕ to $a^4 \rightarrow F_E$ to a^6

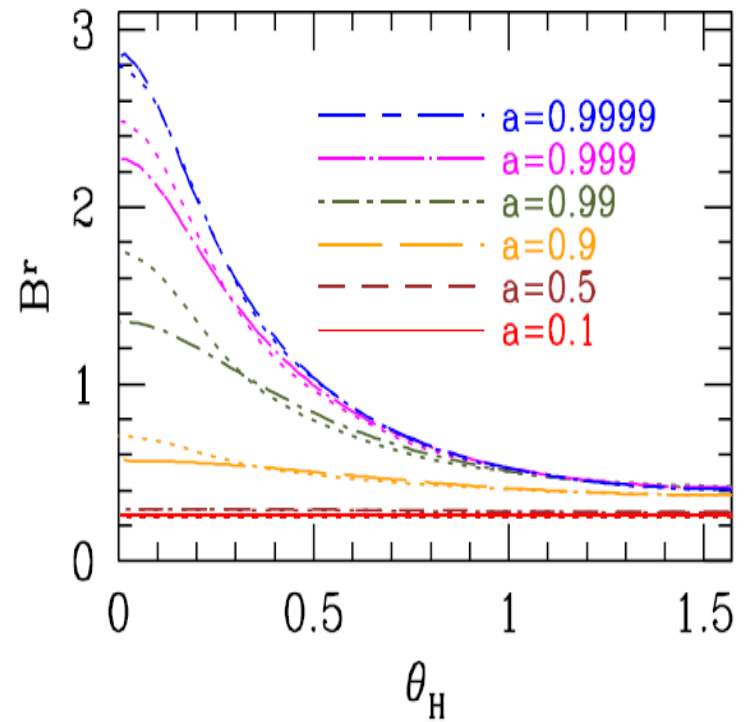
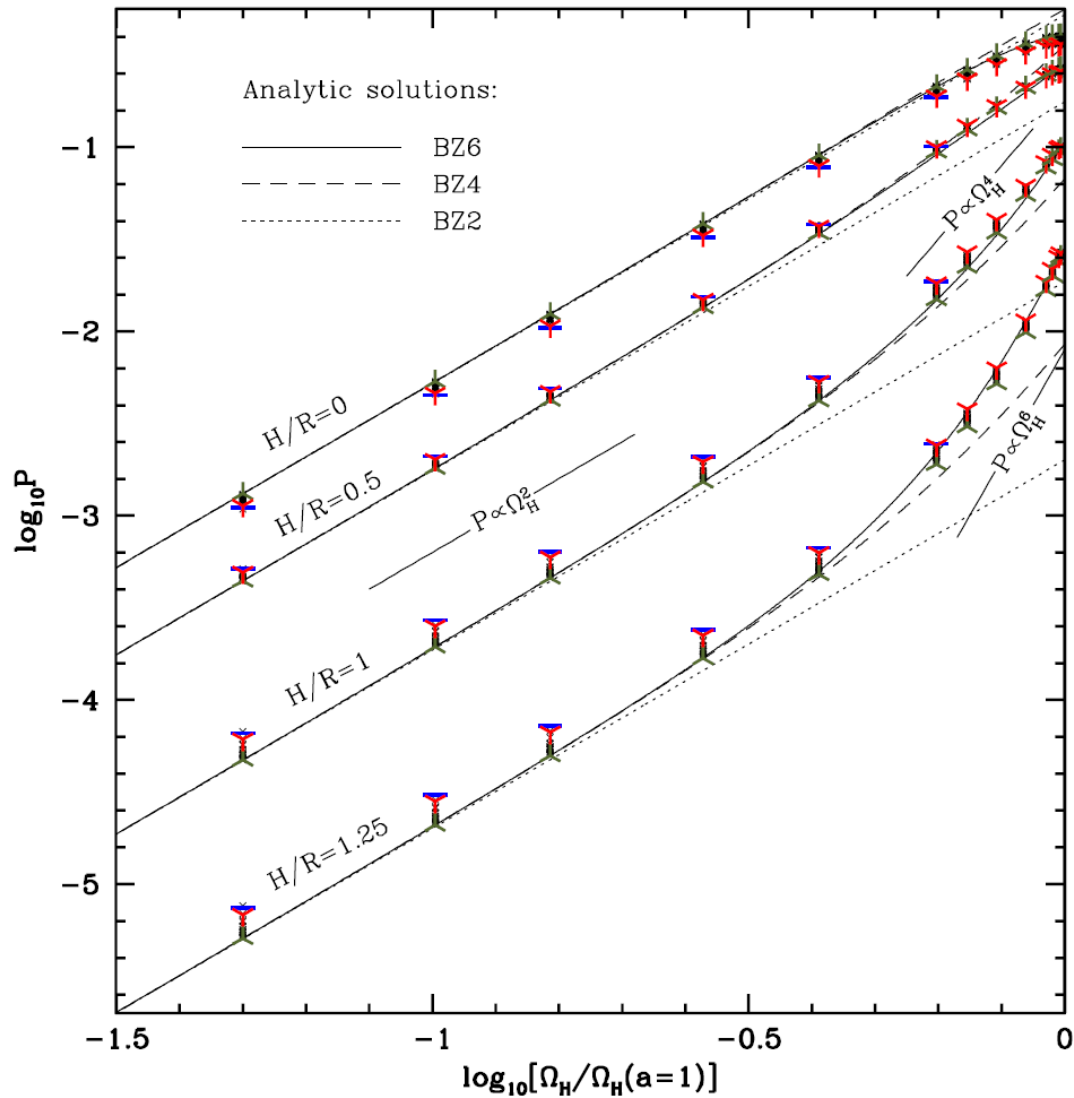
$$\Psi_4(\theta) = \sin^2(\theta) [c_1 \cos^{\alpha_1} \theta + c_2 \cos^{\alpha_2} \theta + c_3 \cos^{\alpha_3} \theta + c_4 \cos^{\alpha_4} \theta]$$

Result 1: P vs. Ω_F accurately fits all simulations

Result 2: Power subtended by smaller angles has

steeper dependence on Ω_F : $P \propto \Omega_F^{(2n)}$

Power Results Reviewed



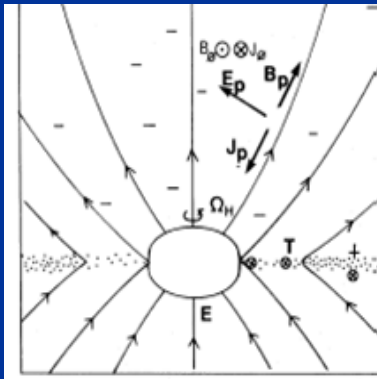
Review: Radio Dichotomy

- Total BH Power Depends upon Ω_F^2
 - Assumes fixed magnetic flux – may even be steeper
- Steeper Dependence at small solid angles (jet)
 - $H/R \sim 1$: $P \propto \Omega_F^4$: (consistent with McKinney 05)
 - $H/R \sim 1.25$: $P \propto \Omega_F^6$
 - (?) $H/R \sim 1.4$: $P \propto \Omega_F^8$
- Can BH Evolution Work? YES!
 - $H/R \sim 1-1.4$ (ADAF for high radio to optical)
 - $a \sim 0.2$ for radio quiet, $a \sim 1$ for radio loud
 - Consistent with spin evolution for elliptical vs. spirals (Volonteri 07)

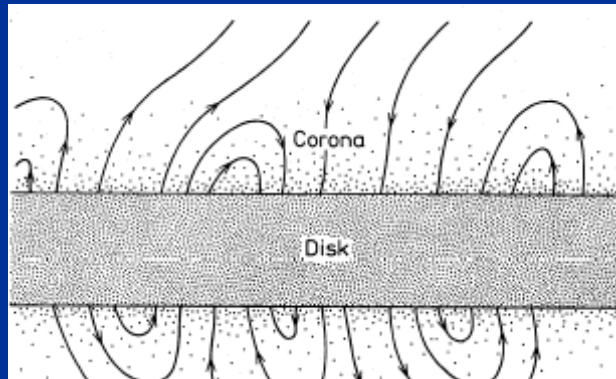
Disk-Jet Coupling Effects

- Old ideas: (Ghosh & Abramowicz 1997; Livio, Ogilvie, Pringle 1999)
 - $\alpha \sim 0.01 - 0.1$ in shearing box, predicts weak field near BH.
 - Sub-equipartition fields assumed near BH
 - Disk more powerful at producing EM jets

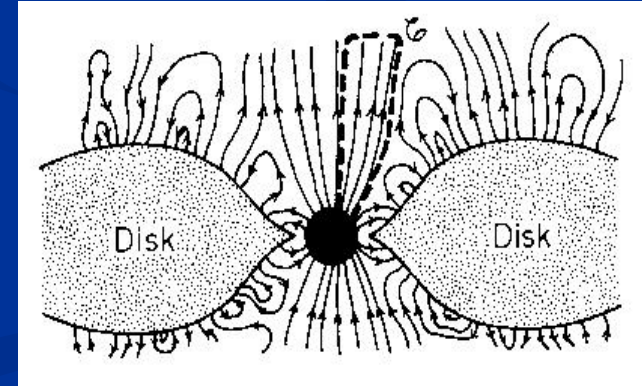
BZ77



Blandford & Payne '82

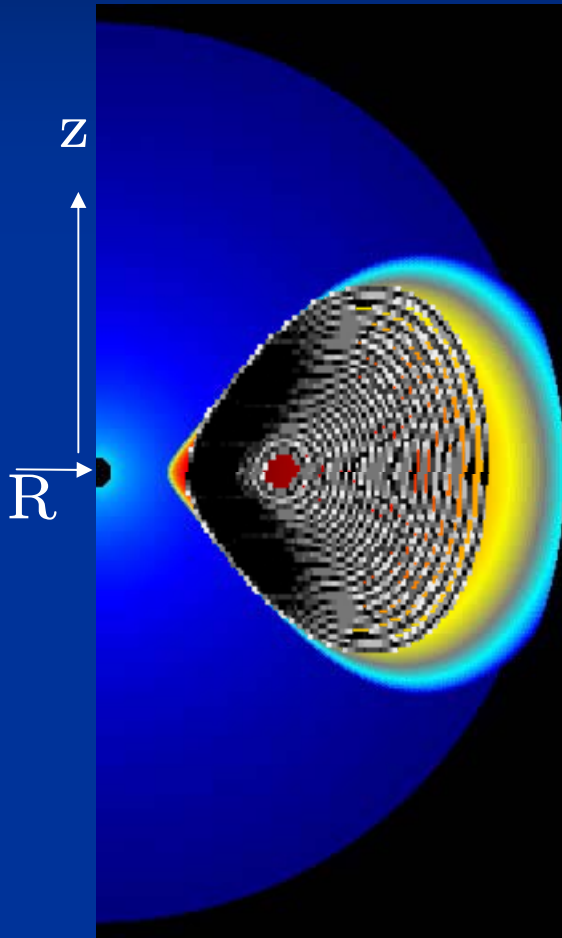


MacDonald & Thorne '82



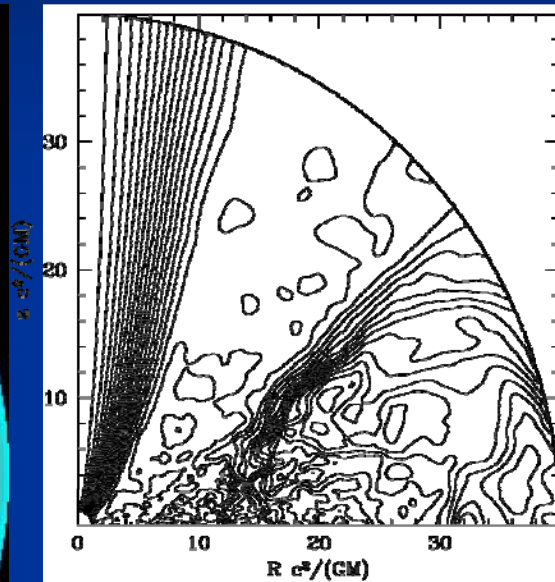
GRMHD Simulations

Log of mass density



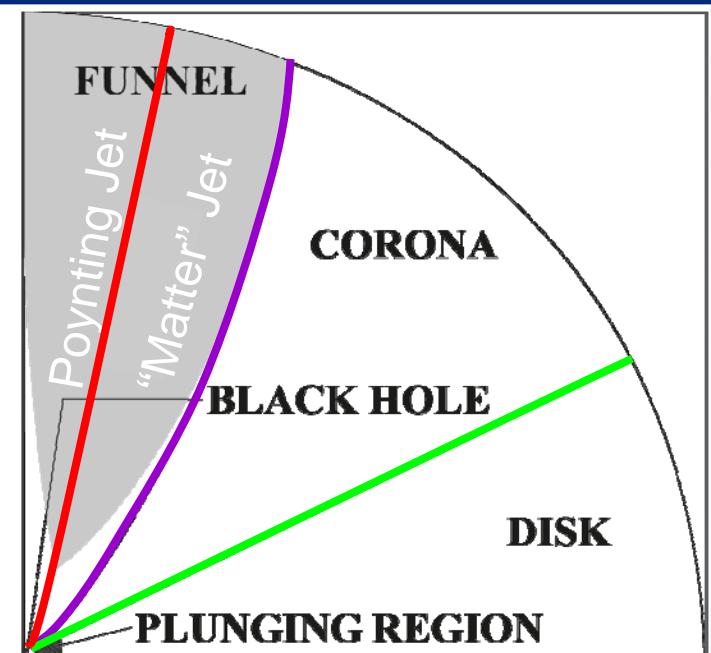
- Evacuated at poles
- Turbulent in equator

Poloidal Field



- Ordered at poles
- Random in equator

Flow Structure

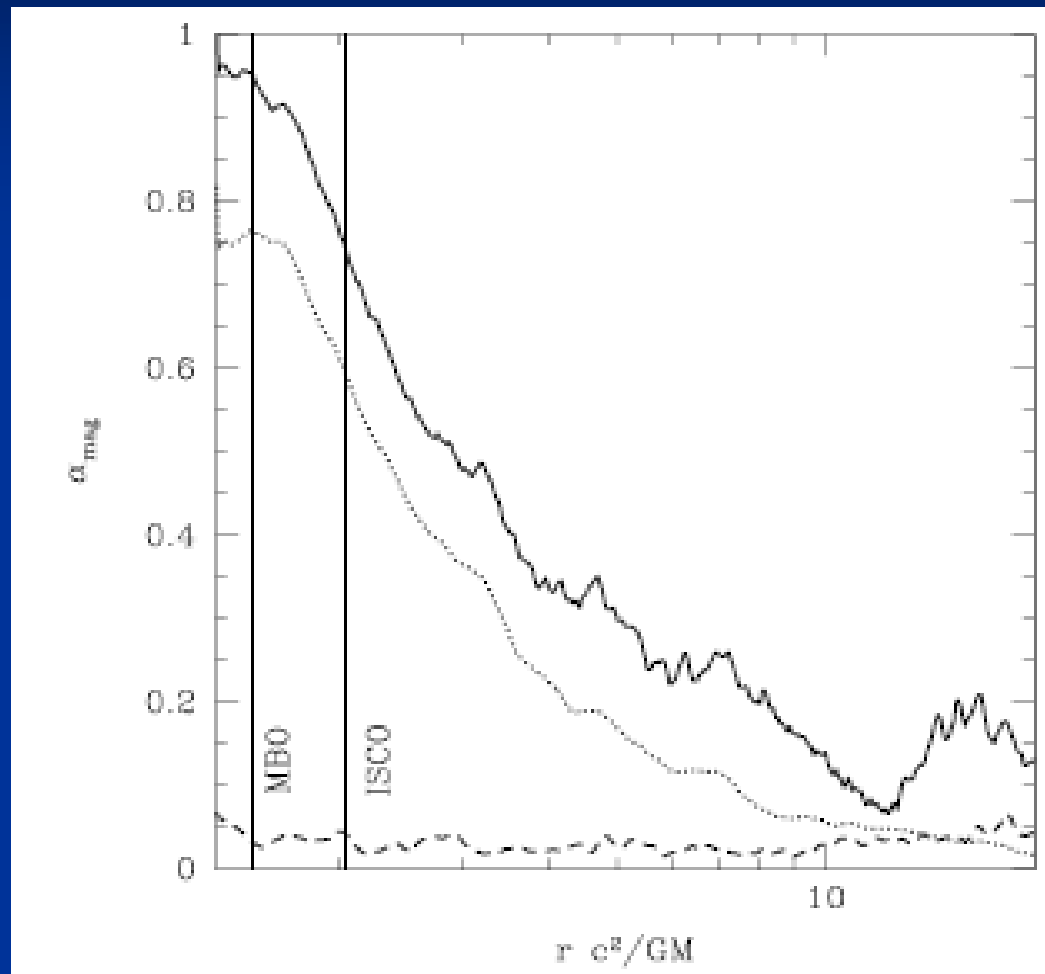


- CORONA:** MA~EM
- FUNNEL:** EM dominated
- JETS:** Unbound, outbound flow

McKinney & Gammie (2004)

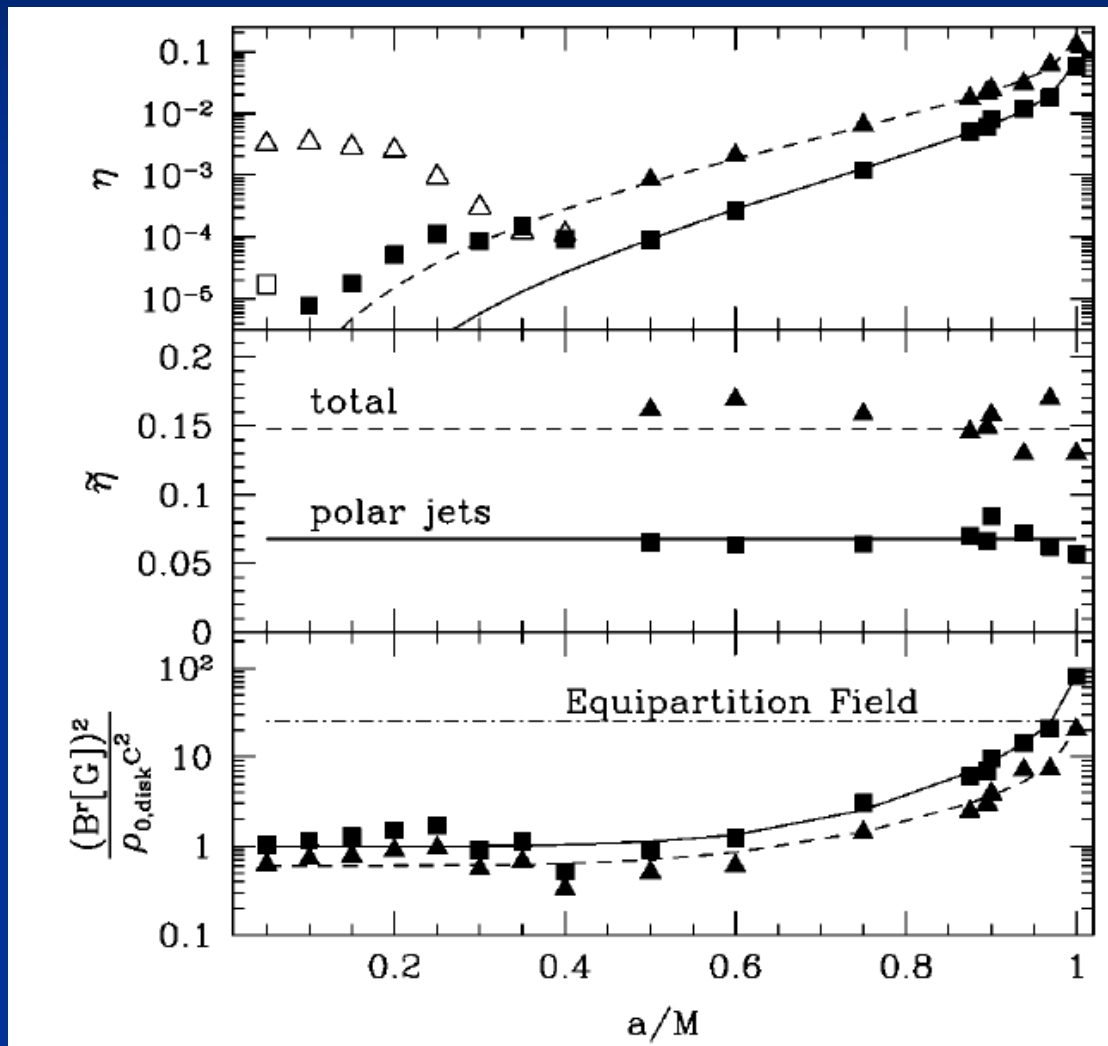
DeVilliers, Hawley, Krolik (2003-2004)

$\alpha \sim 1$ in plunging region

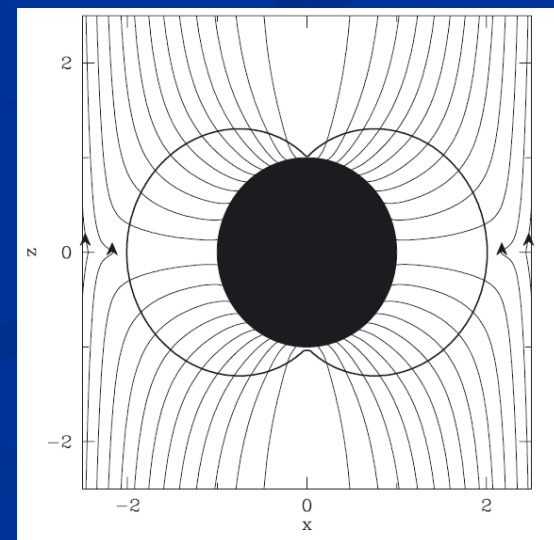
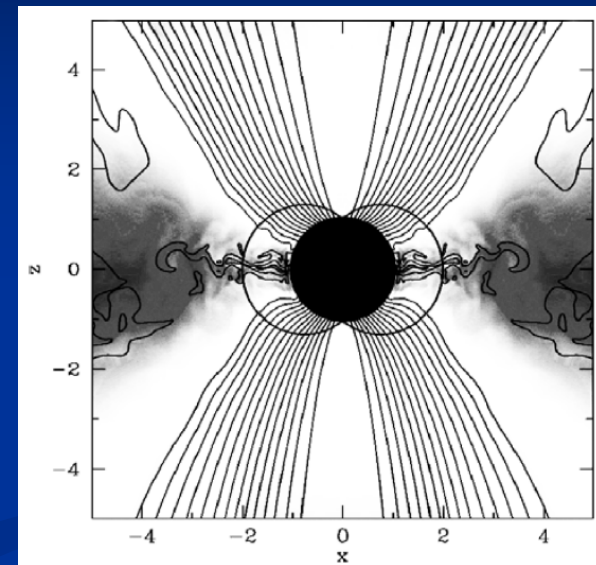


McKinney & Narayan (2007)

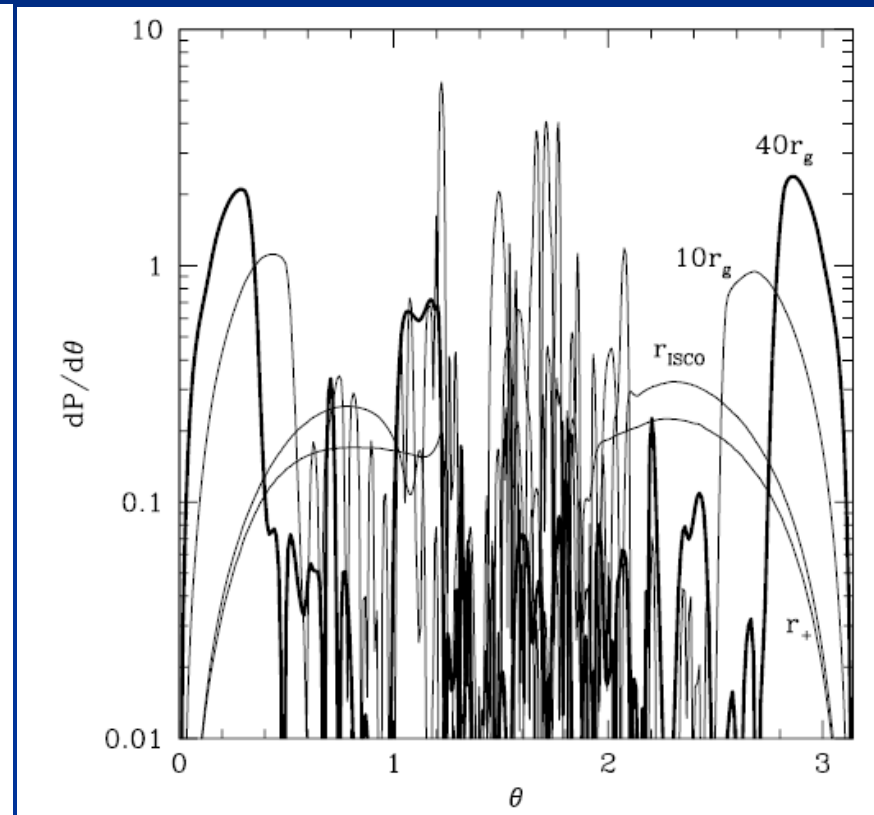
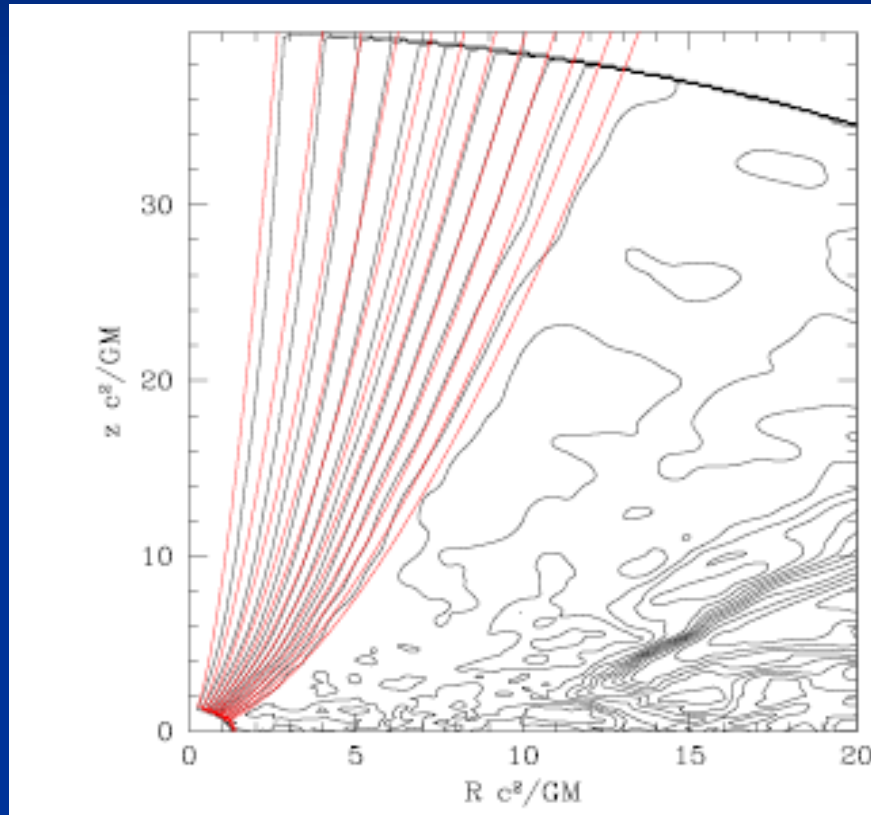
Field becomes super-equipartition for high spin



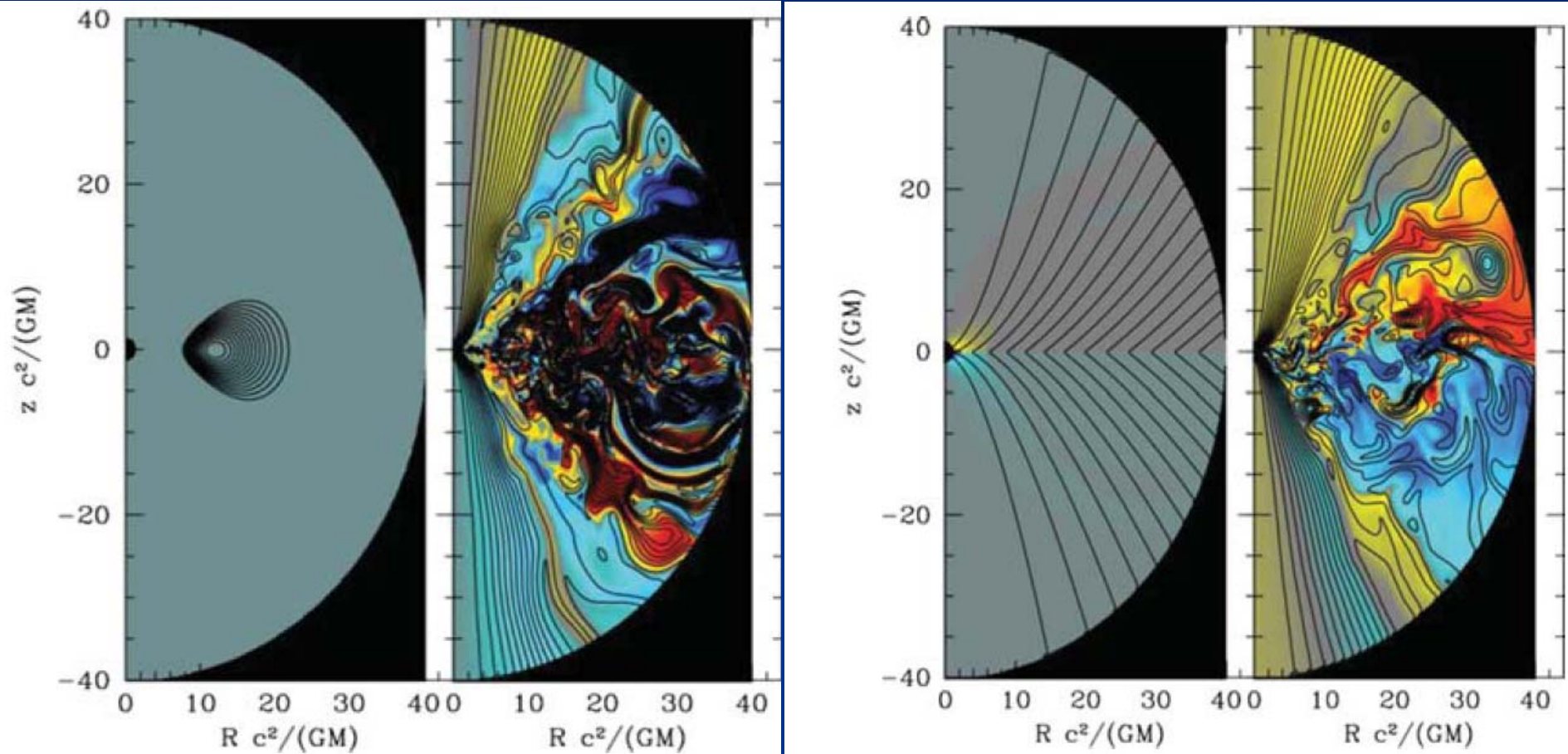
McKinney (2005)



Disk Jet degraded by mass-loading BH cleans field of mass

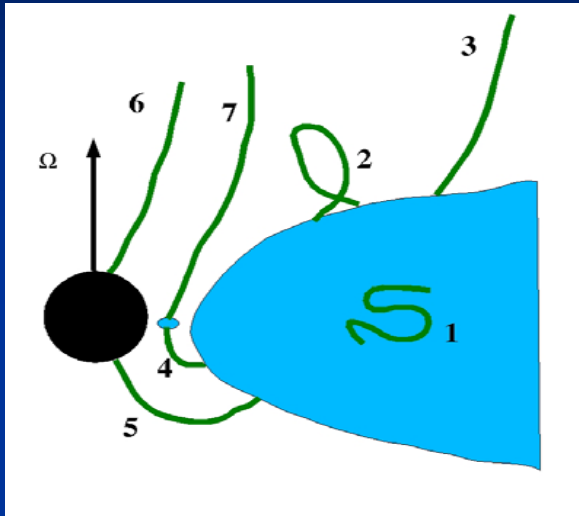


Disk Jet degraded by mass-loading BH cleans field of mass

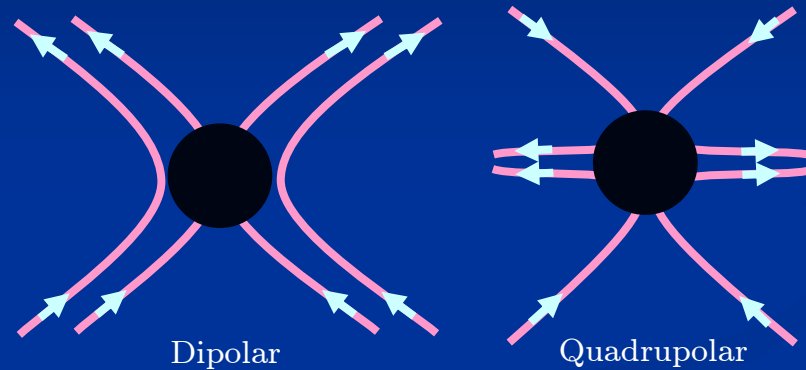


McKinney & Narayan (2007)

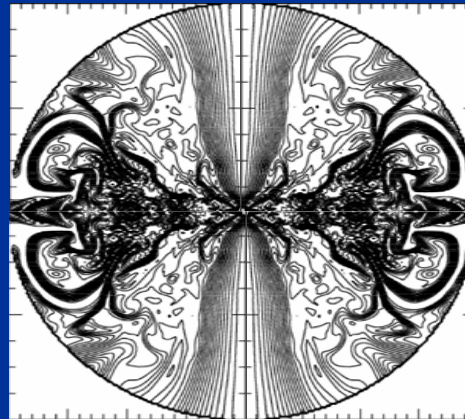
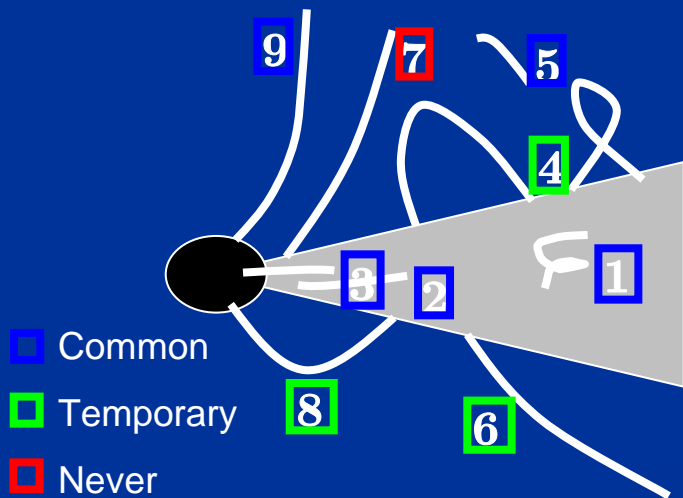
Emergent Magnetic Field Geometry



Blandford '02

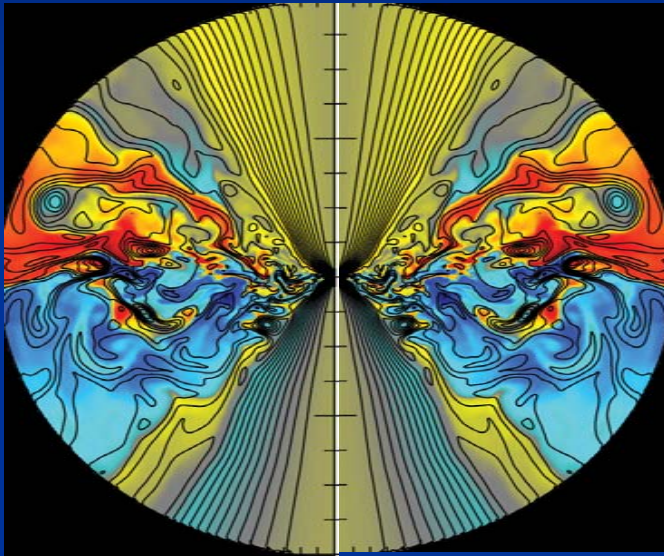


Hirose/McKinney '05



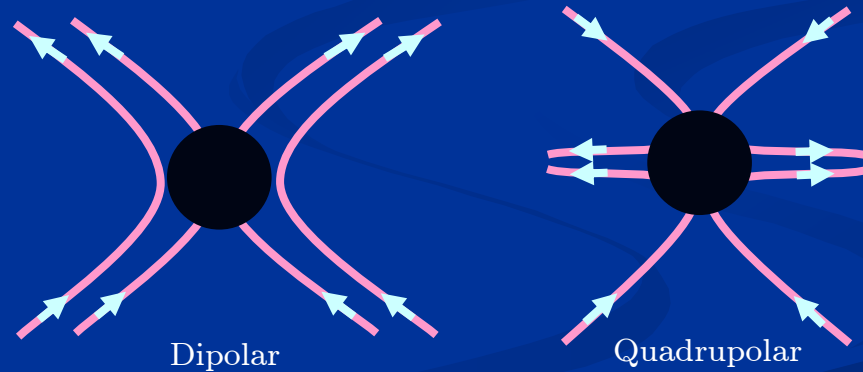
- Balbus & Hawley (MRI) [1]
- Gammie & Krolik [2,3]
- Effect of reconnections [4,5]
- Lovelace or Blandford-Payne [6,7]
- Konigl & Vlahakis [6,7,~9]
- Uzdensky, Matsumoto [8]
- Blandford & Znajek [9]

Jet Formation Stability



Issues:

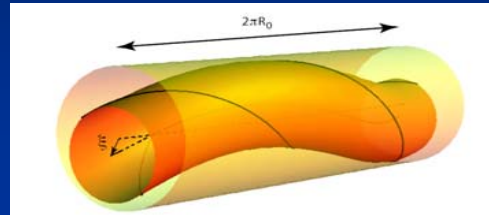
- Jet from Disk or BH?
- Unstable to Turbulence in Disk?
- Unstable to Accreting Disordered Field?



Jet Propagation Stability: Kink

- $|m|=1$ most dangerous: Center-of-mass shifted

$$e^{i(kz + lR + m\phi - \omega t)}$$



- Kruskal-Shafranov non-rel. criterion

$$\frac{|B^\phi|}{|B^z|} > \frac{2\pi R}{L}$$

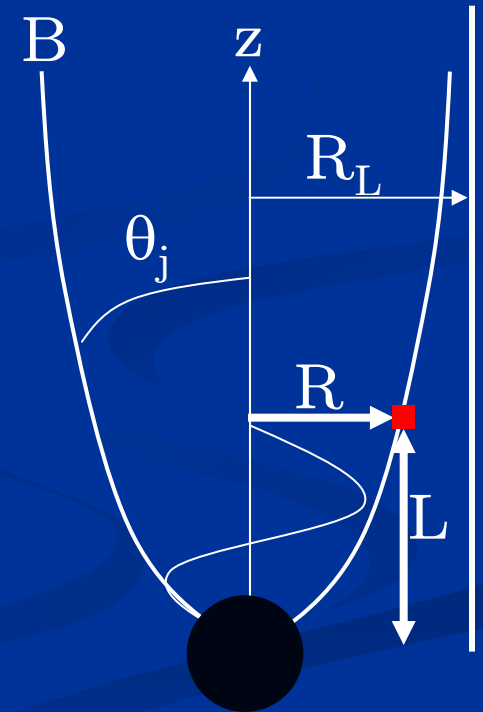
- Tomimatsu (2001) \sim rel. criterion

$$\frac{|B^\phi|}{|B^z|} > \frac{R|\Omega_F|}{c} = \frac{R}{R_L}$$

- Narayan et al. (2009) rel. criterion

$$\frac{|B_{co}^\phi|}{|B_{co}^z|} > \text{few}$$

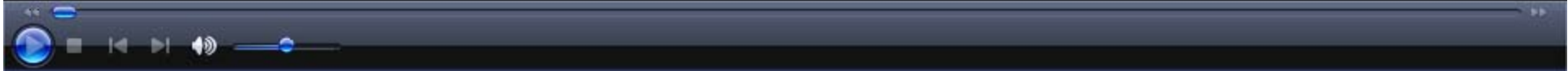
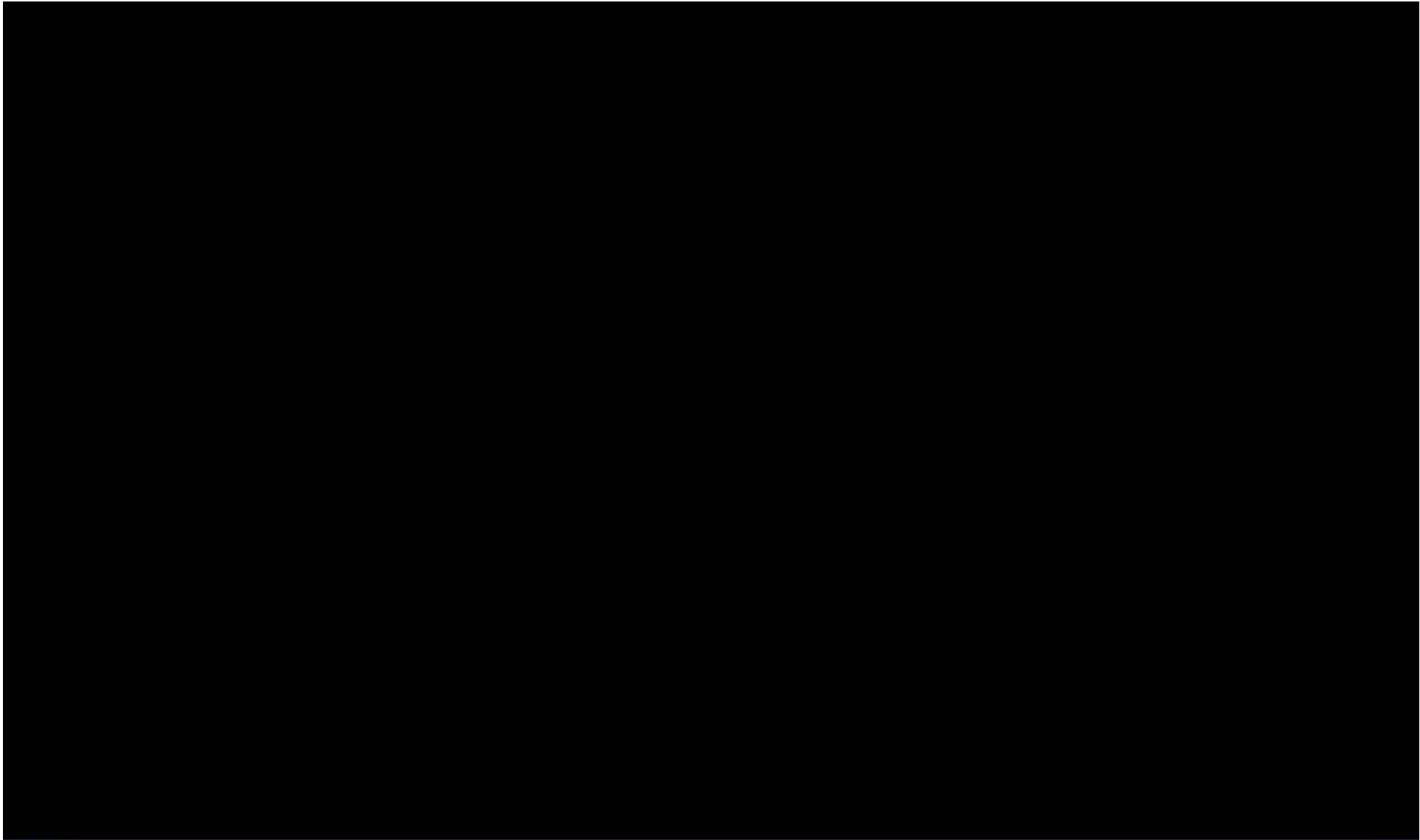
- Expansion & Finite Mass-loading: Jet goes out of causal contact



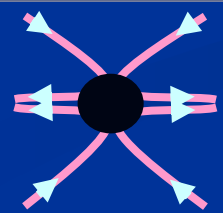
McKinney (2006)
Narayan et al. (2009)

Fully 3D GRMHD Sims

- Initial HD Eq. Thick Torus, $a=0.93$
- Field: Dipolar loop, Quadrupole loops, and large-scale versions
- No symmetries (in θ or ϕ)
 - Required to resolve the dangerous $m=1$ mode
- Conservative HARM 3D w/ Staggered Field
- PPM's base interpolation: 3rd order polynomial fit around flux positions attempted
- $128 \times 64 \times 32$, $64 \times 128 \times 64$, $256 \times 128 \times 32$, $512 \times 256 \times 64$
- Grid resolves disk near BH and jet far from BH



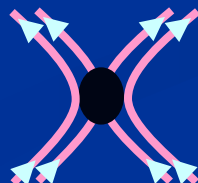
Fully 3D GRMHD Jet Simulations
McKinney & Blandford (2008)



- Quadrupolar Field Jet Fails
- Magnetic field Crucial to explain x-ray binary states: i.e. for Jet or NOT



Fully 3D GRMHD Jet Simulations
McKinney & Blandford (2008)



- Dipolar Field Jet Succeeds: Relativistic Rotation, Expansion, Non-linear Saturation

Review:

- Magnetosphere of BH vs. NS
 - No surface, so flux can be pinched and slip around
 - Stagnation Point: Inflow and Outflow, particle creation
- Black Hole Driven Jet can become Relativistic
 - Requires Organized [mostly dipolar] Field
- Disk Driven Wind-Jet Weakly Relativistic
 - Mass-Loaded by Disk Turbulence
- Stability Maintained by ...
 - Relativistic Rotation of Field Lines
 - Expansion of Jet [and so Causal Disconnection]
 - Finite Mass-Loading [and so Causal Disconnection]
 - Non-linear Saturation [even Non-Rel. Jets can avoid Diss.]