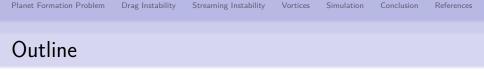


### Dynamics of Gas and Grains II

Christopher J. White

20 November 2012

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2 Drag Instability

Streaming Instability







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Vortices Simulation

Conclusion

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References

# How do we get planetesimals?

- Intermediate particles' orbits decay quickly
- Hurdle:  $1 \text{ cm}-1 \text{ km} (E_{\text{bind}} \ll E_{\text{kinetic}})$
- Solar wind time limit
- Growth must proceed quickly
- Difficult to accomplish with gravity

Conclusion References

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# Overview of Drag Instability

#### Environment

- Disk with particles and gas
- Settled particles
- Drag coupling slows particles coherently
- Mechanism
  - $\begin{array}{ccc} \bullet \ \rho & \uparrow \\ \bullet \ L & \uparrow \\ \bullet \ r & \downarrow \end{array}$
  - $\rho$   $\uparrow$

# Application of Drag

- Minimal drift, particle-dominated limit
- Growth rate pprox 0.49 $(\Sigma_{\rm p}/S_1)^{-2/5} 
  u_*^{-1/5}$ 
  - $S_1 \sim 0.28 \text{ g/cm}^2$
  - $u_* \sim 2.3 \times 10^{-5}$
- Exponential growth in inward-drifting frame

• Competitive with collisions

# Streaming Instability — Background

- $\bullet~\mbox{Gas}$  dragging particles  $\Rightarrow~\mbox{particles}$  affected by gas
- Two mixed fluids
- No vertical differentiation
- No self-gravity
- No turbulence
- Incompressible gas
- Stopping times

$$\tau_{\rm s} = \Omega_{\rm K} \times \begin{cases} \frac{\rho_{\rm s} a}{\rho_{\rm g} c_{\rm g}}, & a < \frac{4}{9} \lambda_{\rm mfp} \ {\rm (Epstein)} \\ \frac{\rho_{\rm s} a}{\rho_{\rm g} c_{\rm g}} \left( \frac{4a}{9 \lambda_{\rm mfp}} \right), & a < \frac{4}{9} \lambda_{\rm mfp} \ {\rm (Stokes)} \end{cases}$$

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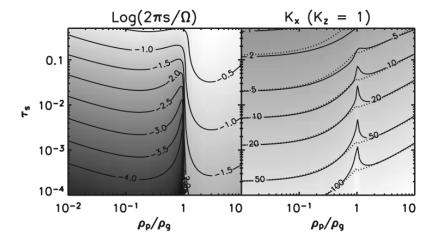
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# Features of Streaming

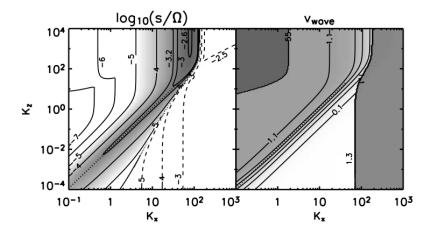
- 6-th order dispersion relation
  - 3 quickly decaying modes
  - 2 epicycles
  - 1 secular mode
- Growth rate dependence

• 
$$ho_{\rm p}/
ho_{\rm g} 
ightarrow 0, \infty \quad \downarrow$$
  
•  $ho_{\rm p}/
ho_{\rm g} 
ightarrow 1 \quad \downarrow$ 



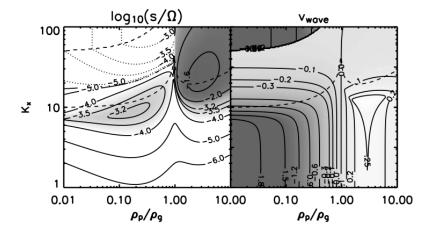
Best growth obtained varying radial wavenumber. (Youdin & Goodman 2005)

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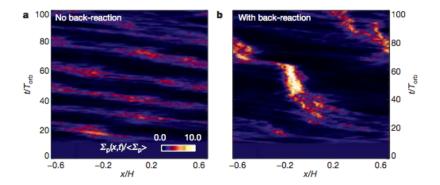


Growth and wave speed for  $\rho_{\rm p}/\rho_{\rm g}=$  0.2,  $\tau_{\rm s}=$  0.01. (Youdin & Goodman 2005)

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Growth and wave speed for  $\tau_s = 0.01$ ,  $K_z = 1$ . (Youdin & Goodman 2005)



(Johansen et al. 2007)

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# Streaming Summary

- Growth faster than diffusion:  $K \lesssim 2\pi \sqrt{s\eta} \Omega \alpha$
- Maximum growth:  $(k/2)(V_{max} V_{min})$
- $\bullet$  Typical conditions:  $\sim 10^{18} \text{--} 10^{20}~\mathrm{g}$

#### Formation of Vortices

- Need turbulence
- Coriolis force at large scales
- Instability or mergers
- Shear  $\Rightarrow$  anticyclonic
- Grow to size of disk thickness and velocity of sound speed

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Separated radially

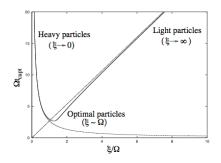
### Particle Capture

- Balance of Coriolis, centrifugal, and friction forces
- Elliptical vortex (aligned with shear), q = a/b
- Light particles:

$$t_{ ext{capt}} = rac{4\xi}{3\Omega^2} rac{q(q-1)^2}{(q-2)(2q+1)}$$

Heavy particles:

$$t_{ ext{capt}} = rac{1}{\xi} rac{2q(q-1)}{(q-3)(2q+1)}$$



Capture times. (Chavanis 2000)

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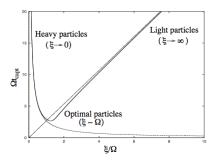
#### Particle Capture

• Optimal trapping:

$$\begin{split} \xi_{\text{opt}} &\approx \left(\frac{3(q-2)}{2(q-1)(q-3)}\right)^{1/2} \Omega \\ t_{\text{capt}}^{\text{opt}} &\approx \left(\frac{8(q-1)^3 q^2}{3(q-3)(2q+1)^2(q-2)}\right)^{1/2} \frac{1}{\Omega} \end{split}$$

• Condition for any trapping

$$q>3$$
  
 $-rac{5}{2}\Omega<\omega<-rac{3}{2}\Omega$ 



Capture times. (Chavanis 2000)

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### Mass Accumulation

- Particles concentrated by turbulence
- Particles brought in by shear
  - Vortex mass

$$\dot{M} = rac{3}{2} \Sigma_{p} \Omega R^{2} (f(\xi))^{2}$$
 $M 
ightarrow rac{3}{2} (\Omega t_{ ext{life}}) \Sigma_{p} R^{2} (f(\xi))^{2}$ 

• Heavy vs. light

$$f(\xi)pprox egin{cases} \left\{ inom{inom{\Omega}}{\xi} 
ight\}^{1/2}, & {
m light} \ rac{\xi}{\Omega}, & {
m heavy} \end{cases}$$

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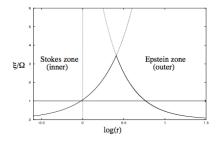
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References

# Application to Solar System

- Minimum mass nebula
- Stokes vs. Epstein

$$\frac{\xi}{\Omega} = \begin{cases} \frac{1913}{a^2 \rho_s} r^{5/4}, & r < r_c \\ \frac{850}{a \rho_s} r^{-3/2}, & r > r_c \end{cases}$$
$$r_c = \left(\frac{4}{9} \frac{a}{1 \text{ cm}}\right)^{4/11} \text{ AU}$$



 $a=30~{
m cm},~
ho_{
m s}=2~{
m g/cm^3}.$  (Chavanis 2000)

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Vortices

References

## Application to Solar System

Transition

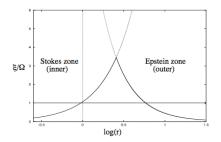
 $1.7 \text{ AU} < r_c < 3.9 \text{ AU}$ 

 Interior optimum near  $1 \, \mathrm{AU}$ 

$$r_{\rm in} = \left(\frac{a^2 \rho_{\rm s}}{1913}\right)^{4/5}$$

• Exterior optimum near 6 AU

$$r_{\rm out} = \left(\frac{850}{a\rho_{\rm s}}\right)^{2/3}$$



 $a = 30 \text{ cm}, \rho_s = 2 \text{ g/cm}^3$ . (Chavanis 2000)

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Vortices

References

# Application to Solar System

Transition

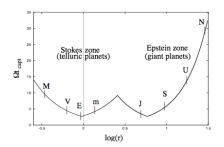
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 $a = 30 \text{ cm}, \ \rho_{s} = 2 \text{ g/cm}^{3}.$  (Chavanis 2000)

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Streaming Instability

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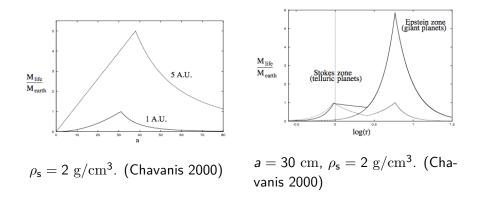
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### Application to Solar System

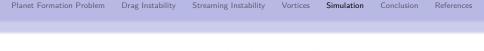


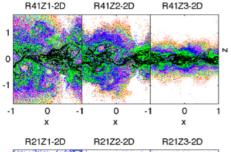
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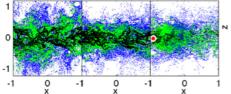
References

# Simulation Considerations

- Hydrodynamic
- Grain sizes
- Disk structure
- 2D vs. 3D
- Pure hydro vs. MHD







Short stopping times. (Bai and Stone 2010)

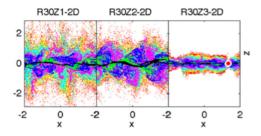
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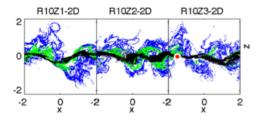
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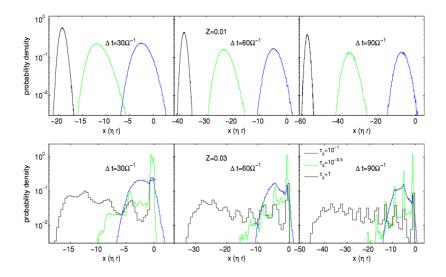


Long stopping times. (Bai and Stone 2010)

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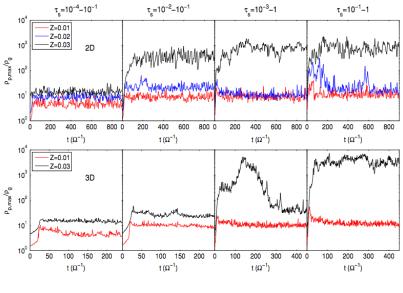
# Simulation Results

- Streaming instability
  - $\tau_{\rm s}\gtrsim 10^{-2}$
  - Prevents Kelvin-Helmholtz
  - Large, abundant particles
- Range of particle sizes
  - Clumping counters radial drift
  - Reduced collisional velocity
- Formation in dead zone



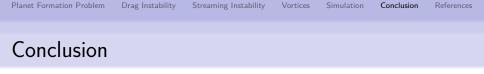
Radial drift. (Bai and Stone 2010)

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Particle density. (Bai and Stone 2010)

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- $\bullet\,$  Need mechanism to grow to  $\rm km$  sizes
- Drag hurts and helps
- Streaming instability critical
- Vortical structure may play role
- Seen in simulations
- Generality?

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